

Structure and Dynamics of High-Z Ions Studied in Relativistic Atomic Collisions at the ESR Storage Ring

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- **Introduction**

- **Experiments at the jet-target**
Alignment studies: population of magnetic sublevels

=> Interference between E1 and M2 transition amplitudes

- **Experiments at the electron cooler**

Two-electron contribution to the K-shell ionization potential in He-like uranium

- **Summary and Outlook**

Collaboration

Experiment

H.F. Beyer, K. Beckert, G. Bednarz, F. Bosch, R.W. Dunford, B. Franzke, A. Gumberidze, S. Hagmann E. Kanter, O. Klepper, A. Krämer, C. Kozhuharov, D. Liesen, T. Ludziejewski, P.H. Mokler, X. Ma, A. Muthig, M. Steck, Z. Stachura, Th. Stöhlker, S. Toleikis, A. Warczak

Atomic Physics Group, GSI-Darmstadt, Germany

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University of Cracow, Poland

University of Frankfurt, Germany

Institute for Nuclear Studies, Swierk, Poland

Argonne National Laboratory, Argonne, USA

Kansas State University, Kansas, USA

IMP, Lanzhou, China

Theory

J. Eichler, S. Fritzsch, A. Ichihara, D.C Ionescu, T. Shirai, A. Surzhykov

Theoretische Physik, HMI-Berlin, Germany

JAERI, Japan

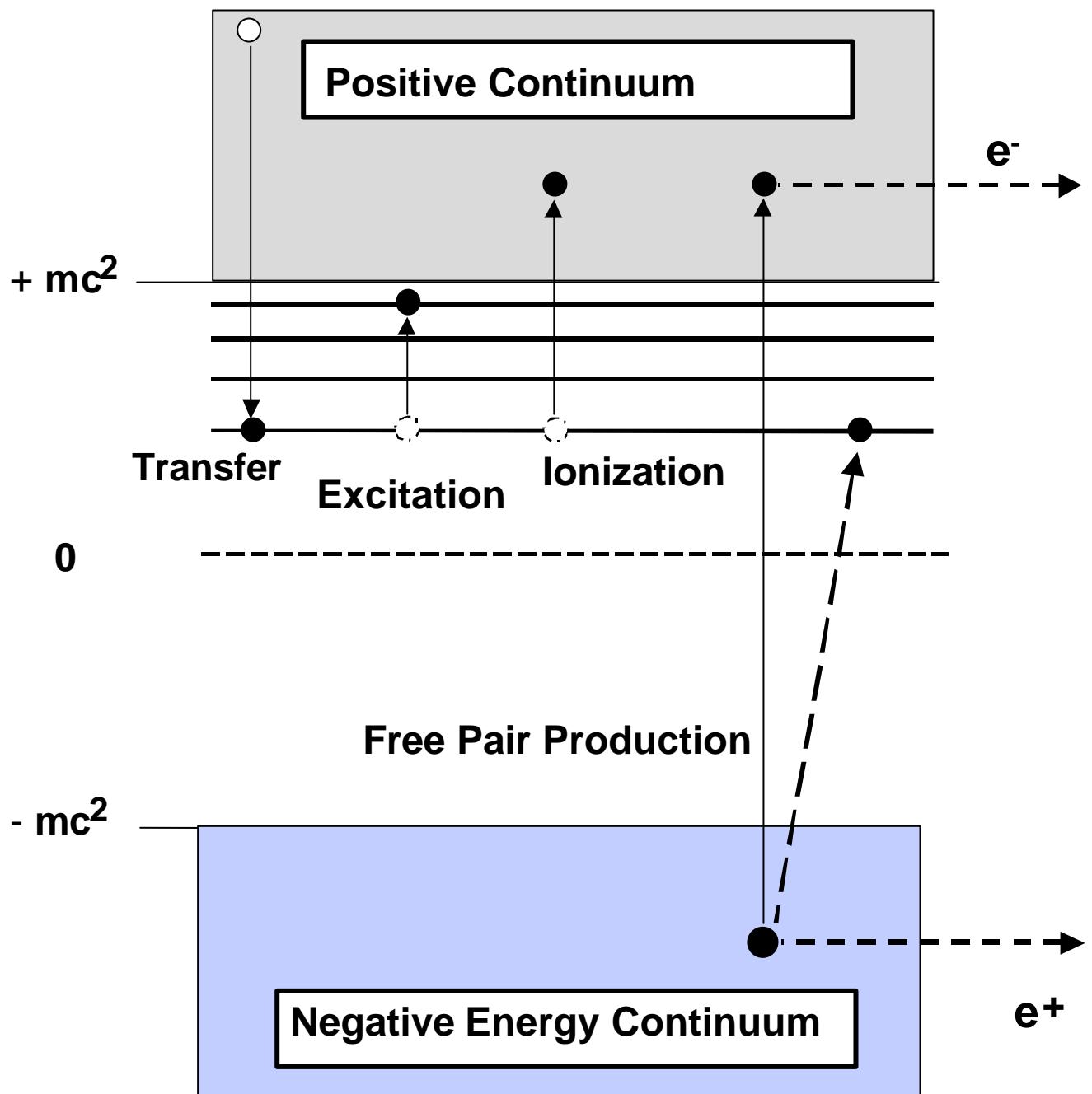
TU-Dresden, Germany

GSI-Darmstadt, Germany

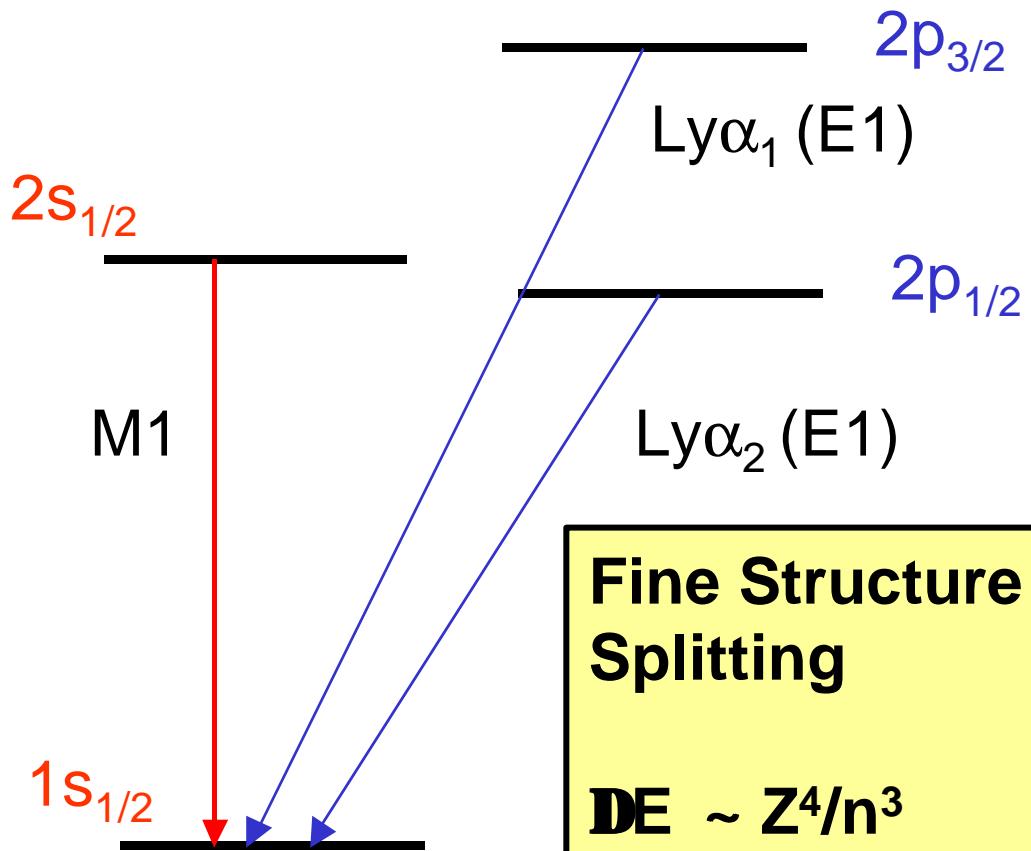
University of Kassel, Germany

Collision Dynamics of Relativistic Heavy Ions

Collision times in the sub-attosecond regime
 $(10^{-19} \text{ s} < t < 10^{-21} \text{ s})$



The Structure of One-Electron Systems



Fine Structure Splitting

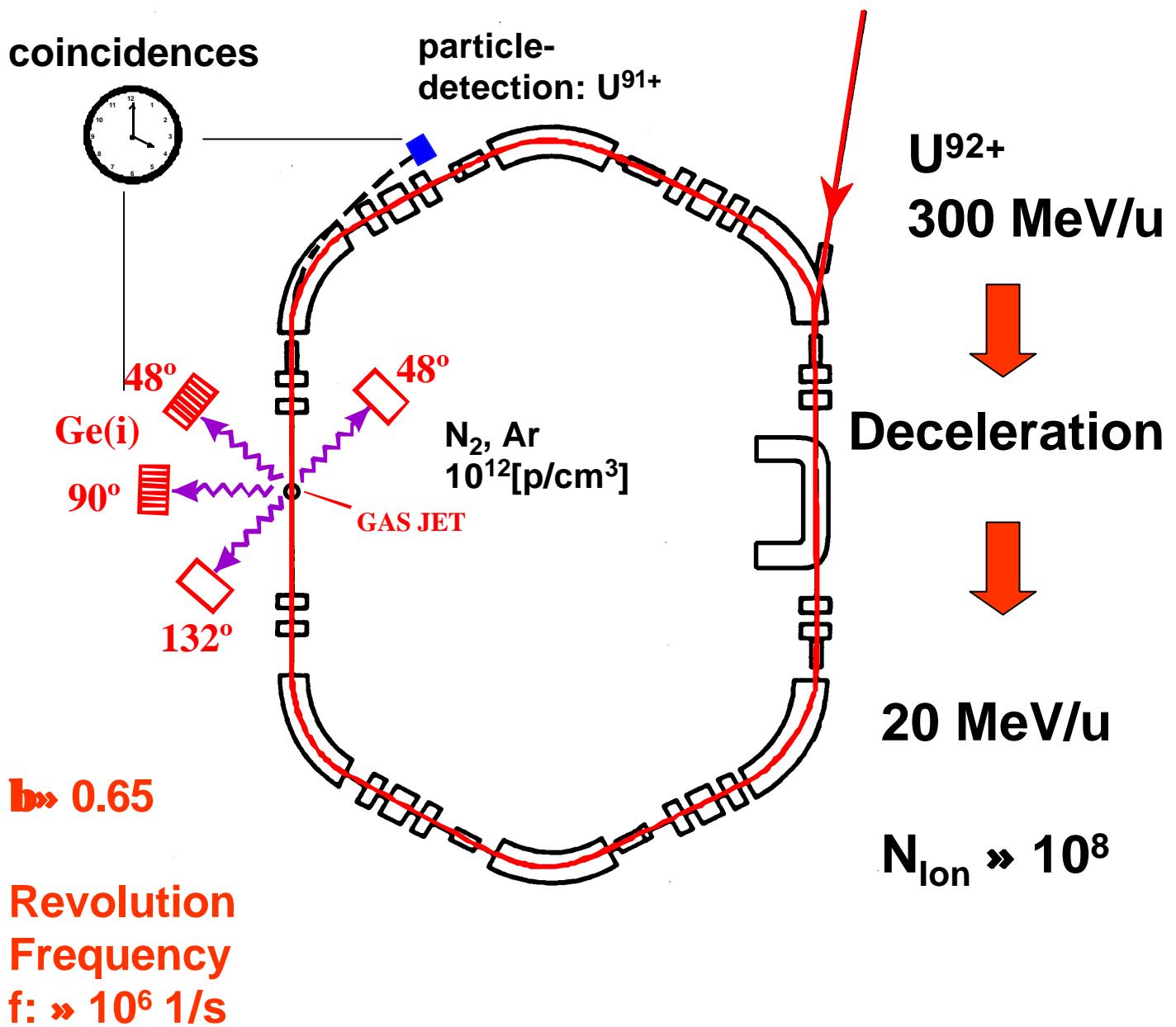
$$\Delta E \sim Z^4/n^3$$

Z: nuclear charge number
n: principal quantum number

large shell and sub-shell splitting in high-Z systems

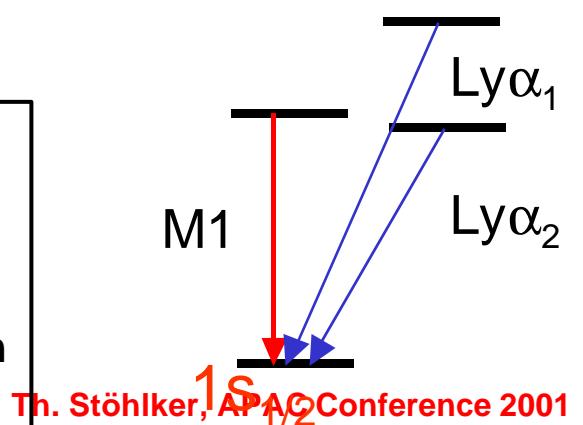
- *photon emission enables to study dynamic atomic processes state selectively*
- *Via angular distributions, even information about magnetic substates get accessible*

X-Ray Spectroscopy at the Jetttarget



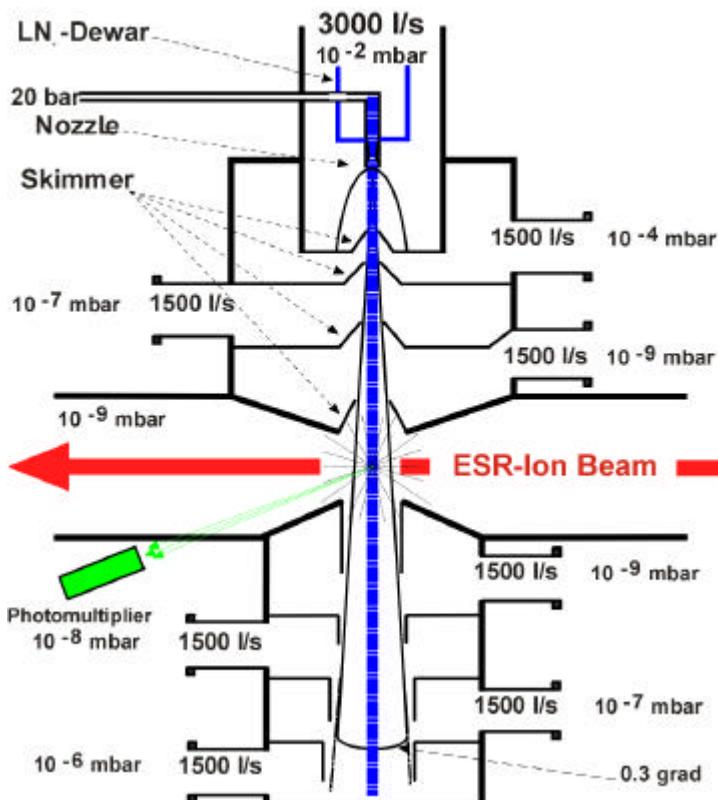
$$E_{\text{lab}} = \frac{E_{\text{proj}}}{\tilde{a} \cdot (1 - \hat{a} \cdot \cos \hat{\epsilon}_{\text{lab}})}$$

E_{lab} : Photon energy in laboratory system
 E_{proj} : Photon energy in emitter system



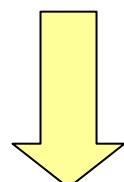
The Jet-Target

Supersonic jet, operates in ultra high vacuum environment (10^{-11} mbar)

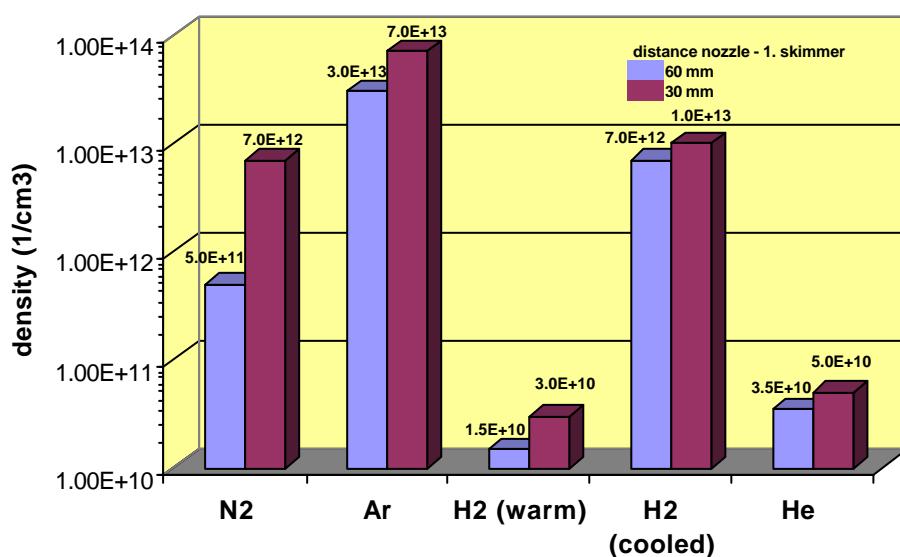


Target densities

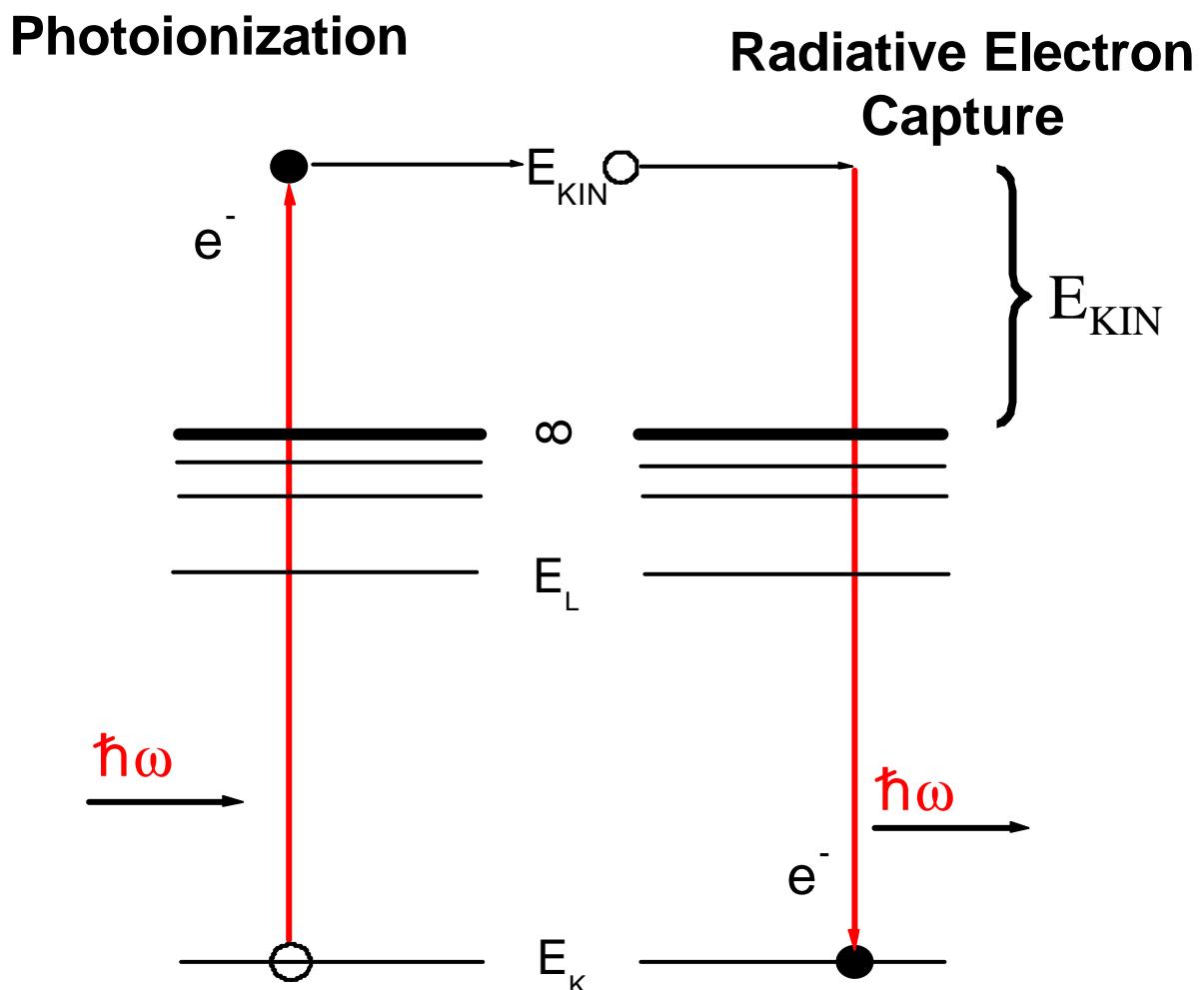
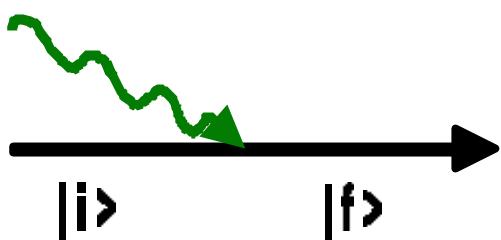
$10^{12} - 10^{14}$ p/cm³



Single collision conditions



Photon-Matter Interaction in the Relativistic Regime: *Study of Photoionization of high-Z Ions*

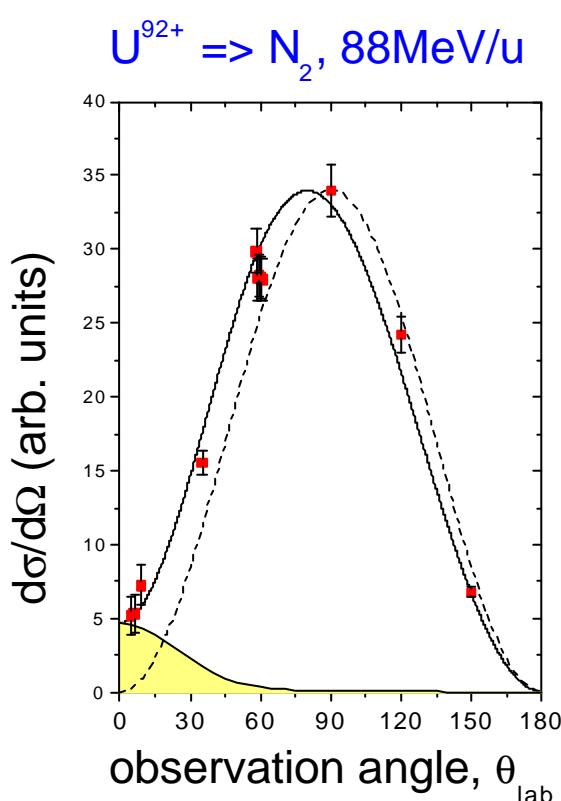


Angular Distribution

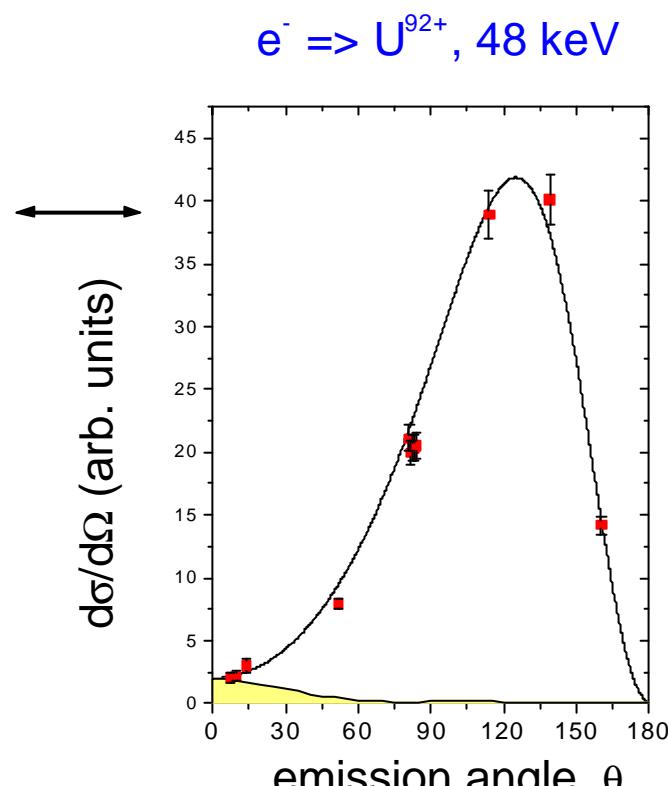
radiative capture into U^{92+} and photoionization of U^{91+}



laboratory frame



emitter frame



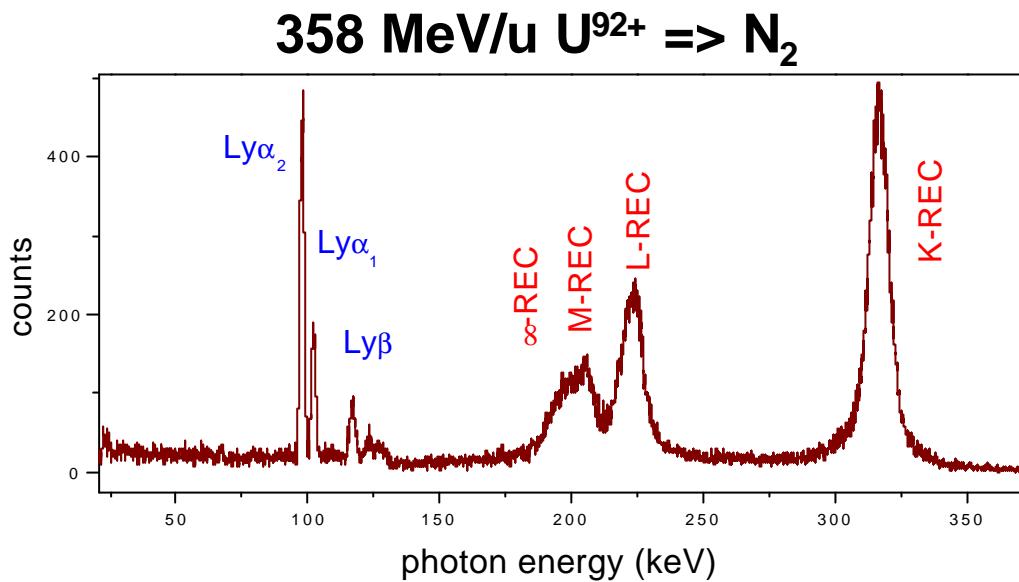
Theory: J. Eichler 1999

spin-flip contribution: 2%

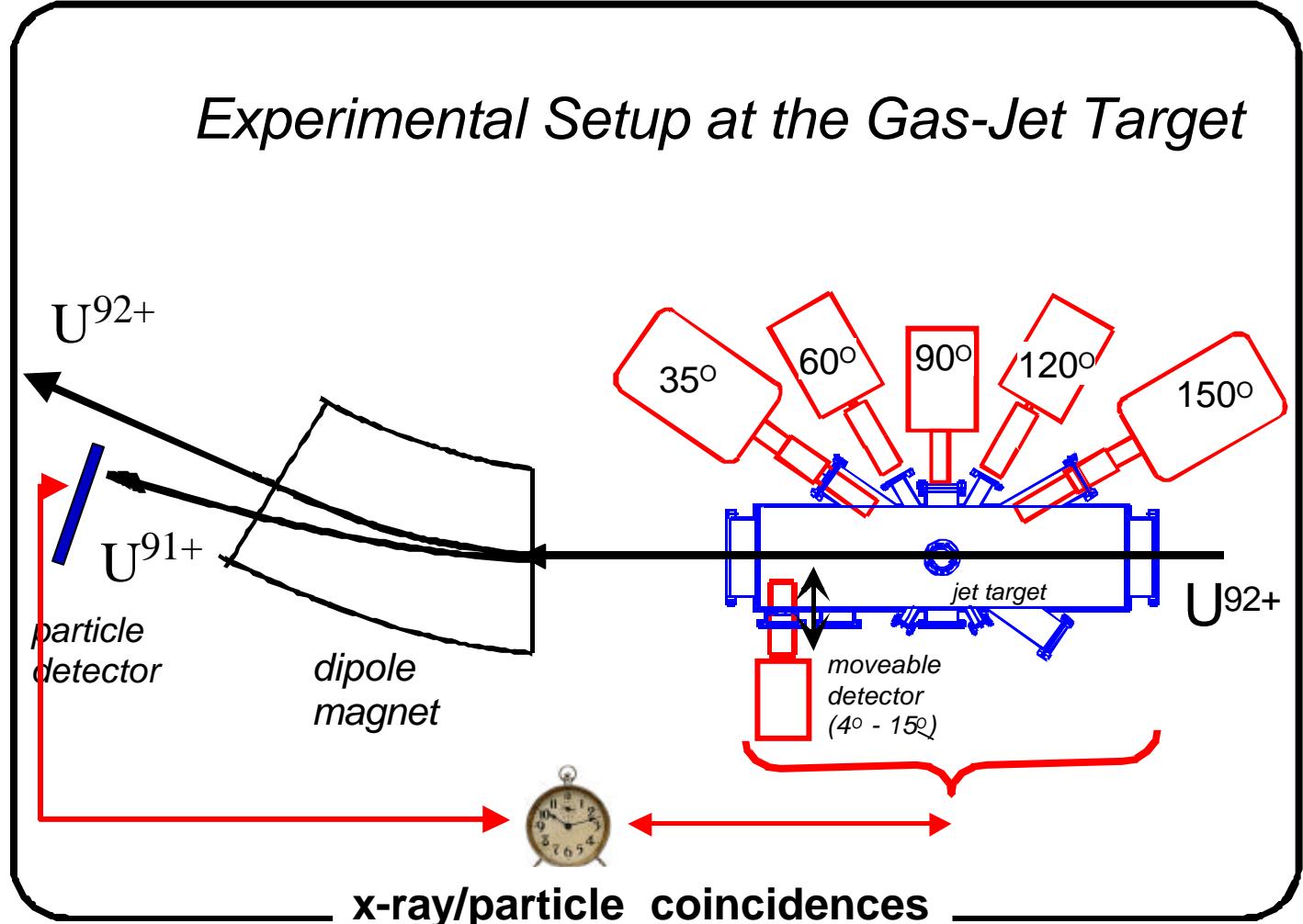
Th. Stöhlker et al., PRL 82, 3232 (1999)
PRL 86, 983 (2001)

Th. Stöhlker, APAC Conference 2001

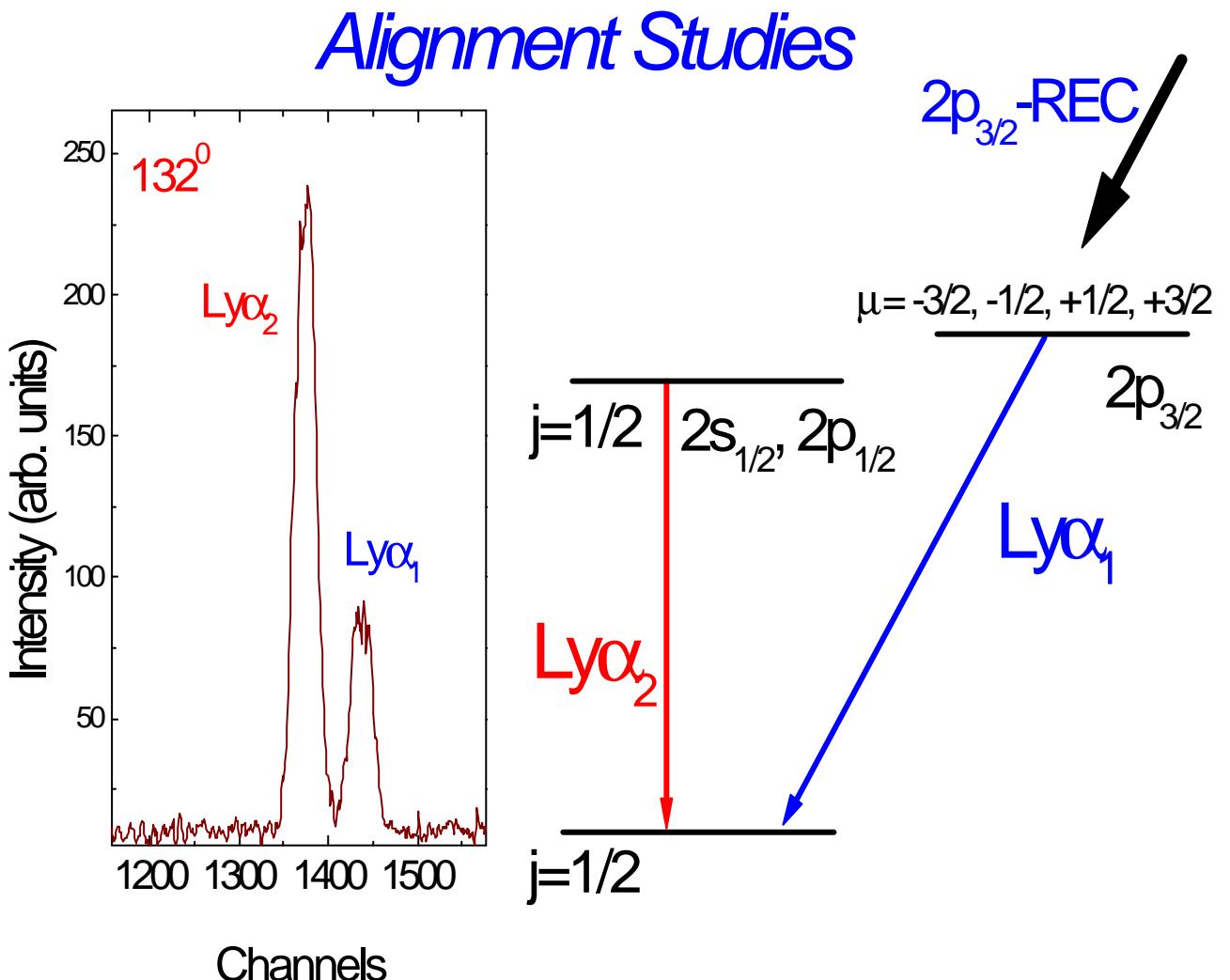
Angular Distribution Studies for the Time-Reversed Photoionization Process



Experimental Setup at the Gas-Jet Target



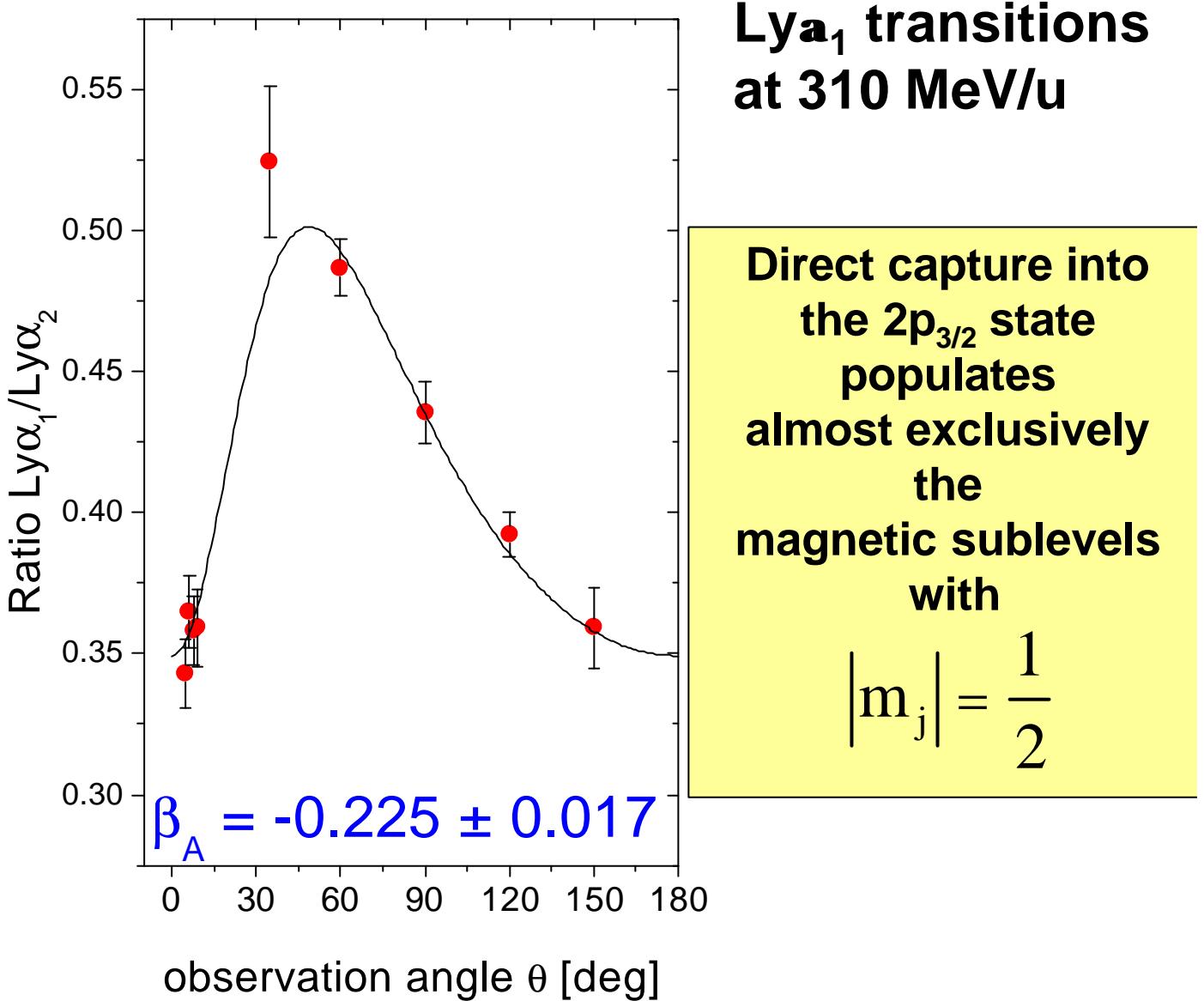
2p_{3/2} transitions in high-Z ions produced by REC: a source of polarized high energy photons



Theory: J. Eichler Nucl. Phys. A572, 147 (1994)
Inverse of Two-Photon One-Electron Ionization

1st photon: resonant excitation to the
2p_{3/2} state
+
2nd photon: ionization

Dt » 10⁻¹⁷ s



Alignment

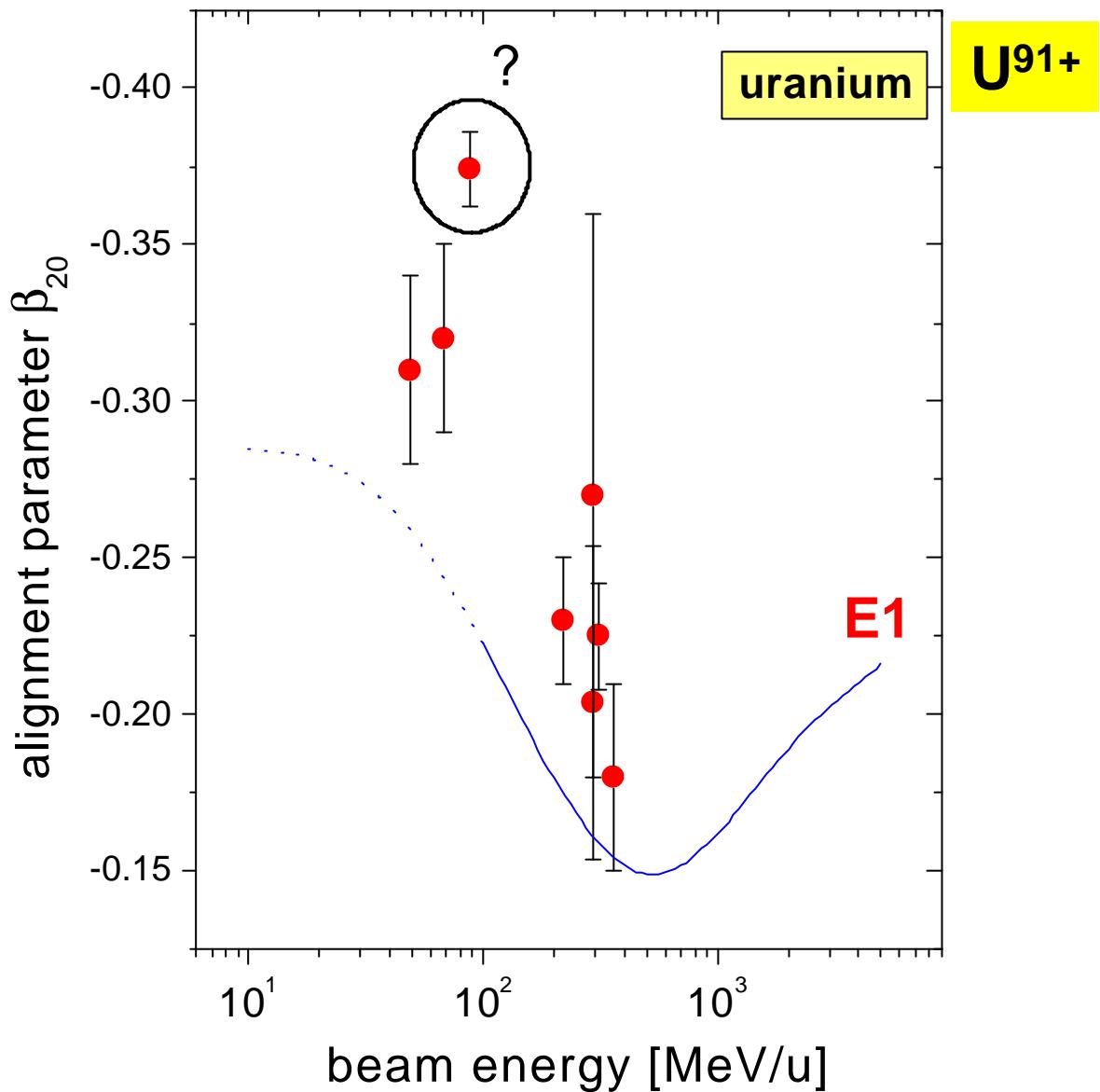


Alignment Parameter

$$W(\theta) \propto 1 + \beta_A \left[1 - \frac{3}{2} \sin^2 \theta \right]$$

$$\beta_A = \frac{1}{2} \frac{\sigma\left(\frac{3}{2}\frac{3}{2}\right) - \sigma\left(\frac{3}{2}\frac{1}{2}\right)}{\sigma\left(\frac{3}{2}\frac{3}{2}\right) + \sigma\left(\frac{3}{2}\frac{1}{2}\right)}$$

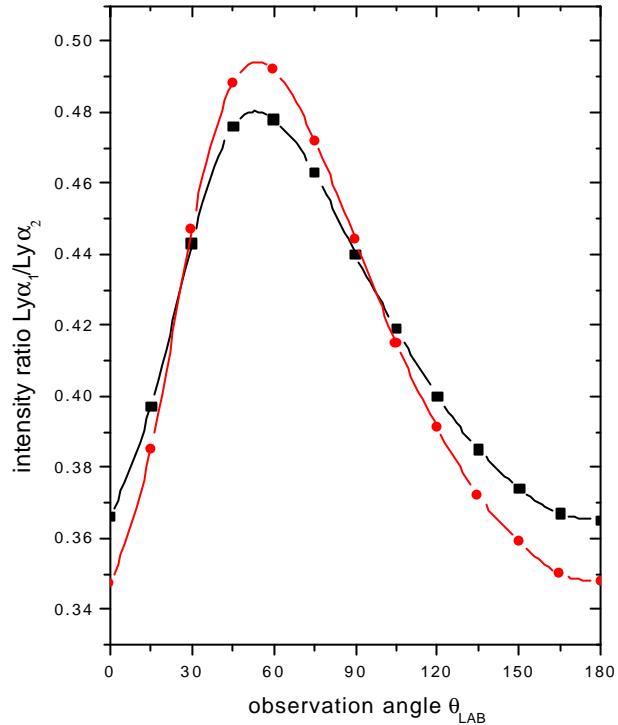
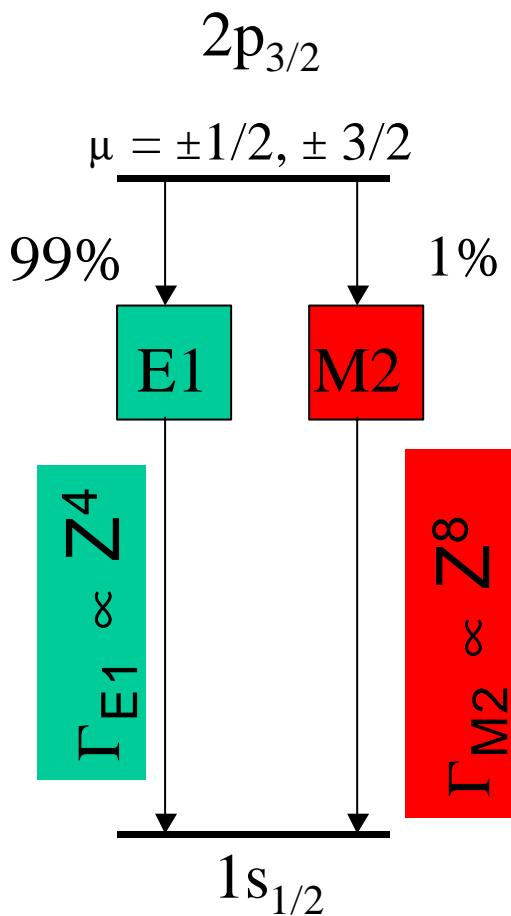
Th. Stöhlker, APAC Conference 2001



***Strong alignment observed for REC into the $2p_{3/2}$ state
but data deviate markedly from theory***

(Th. Stöhlker et al., PRL 79, 3270 (1997))

Interference between the E1 and M2 transition amplitudes leads to a disturbed angular distribution for the $2p_{3/2}$ decay



Photon angular distribution of the decay of the $2p_{3/2}$ state in H-like ions

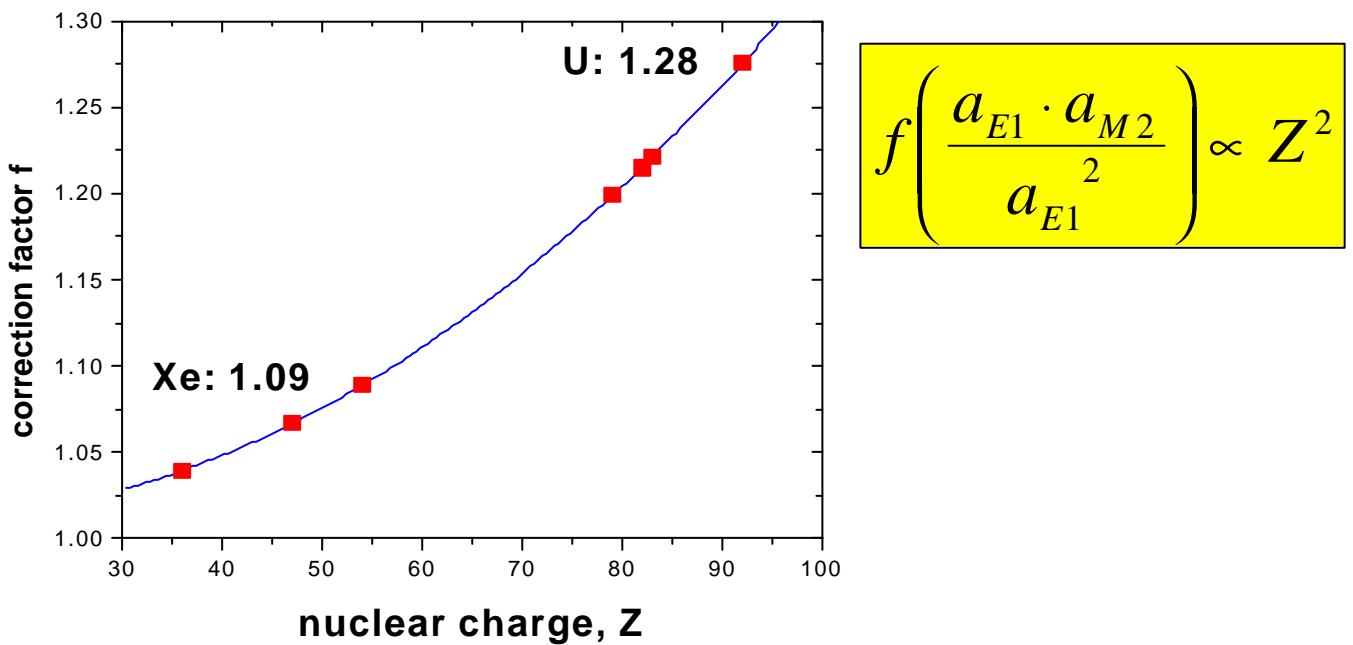
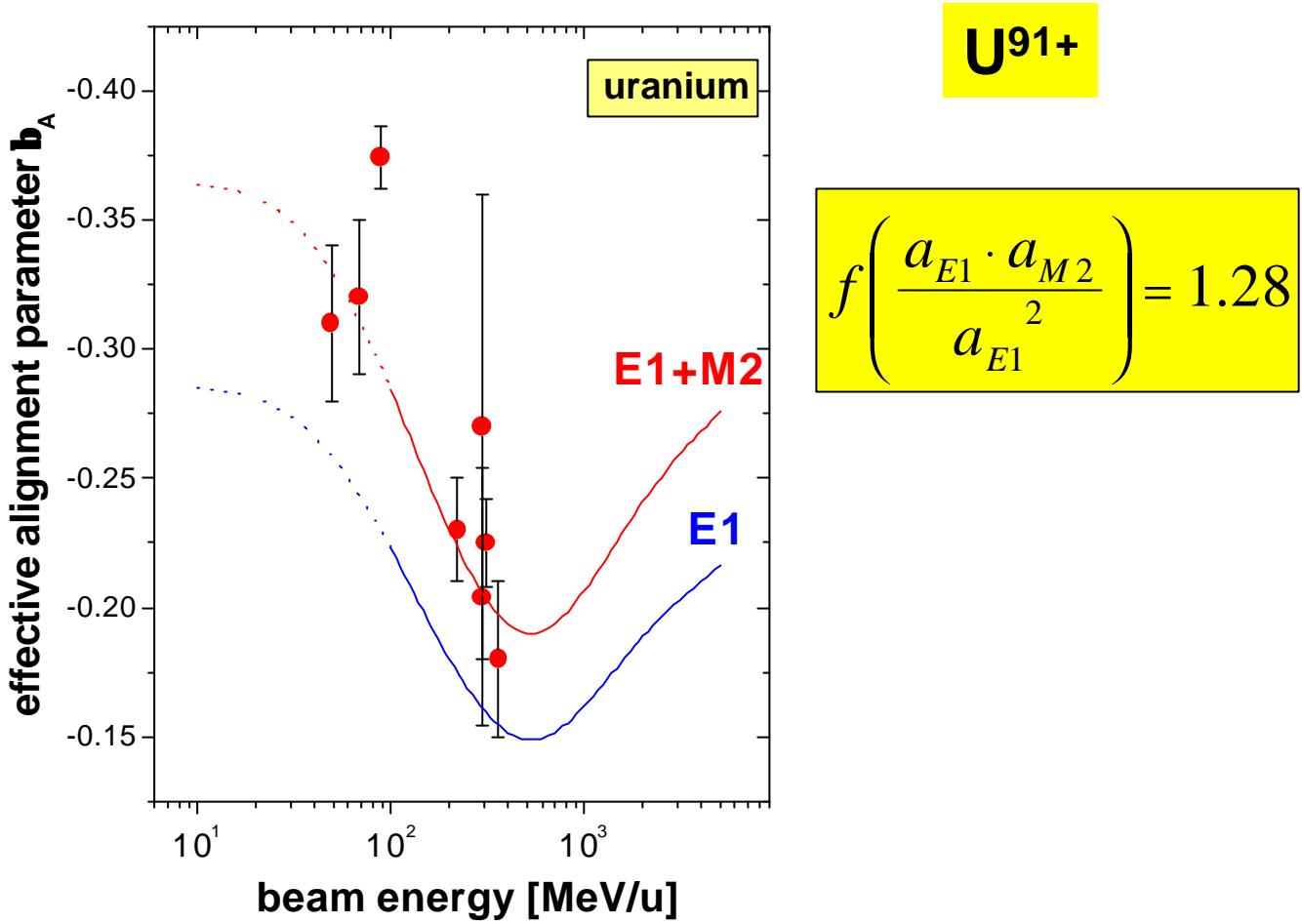
$$W(\theta) \propto 1 + f \left(\frac{a_{E1} \cdot a_{M2}}{a_{E1}^2} \right) \bullet \beta_A \bullet \left[1 - \frac{3}{2} \sin^2 \theta \right]$$

for Z=92

$$f \left(\frac{a_{E1} \cdot a_{M2}}{a_{E1}^2} \right) = 1.28$$

Prediction by S. Fritzsche und A. Surzhykov, 2001

TM St. Petersburg, XPA Conference 2001



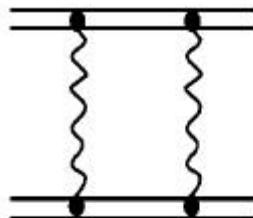
Electron-Electron Interaction in Strong Fields

Measurement of the 2eQED for Uranium at the ESR

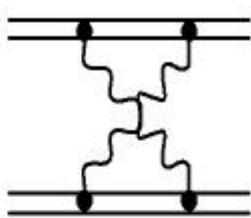
Accuracy:
2 eV; $\Delta E/E \approx 0.1\%$

Z=92

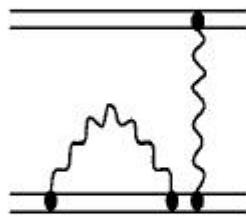
Two-Electron Contribution: **2246.0 eV**



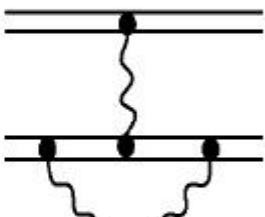
a)



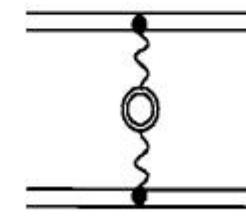
b)



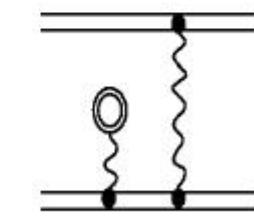
c)



d)



e)



f)

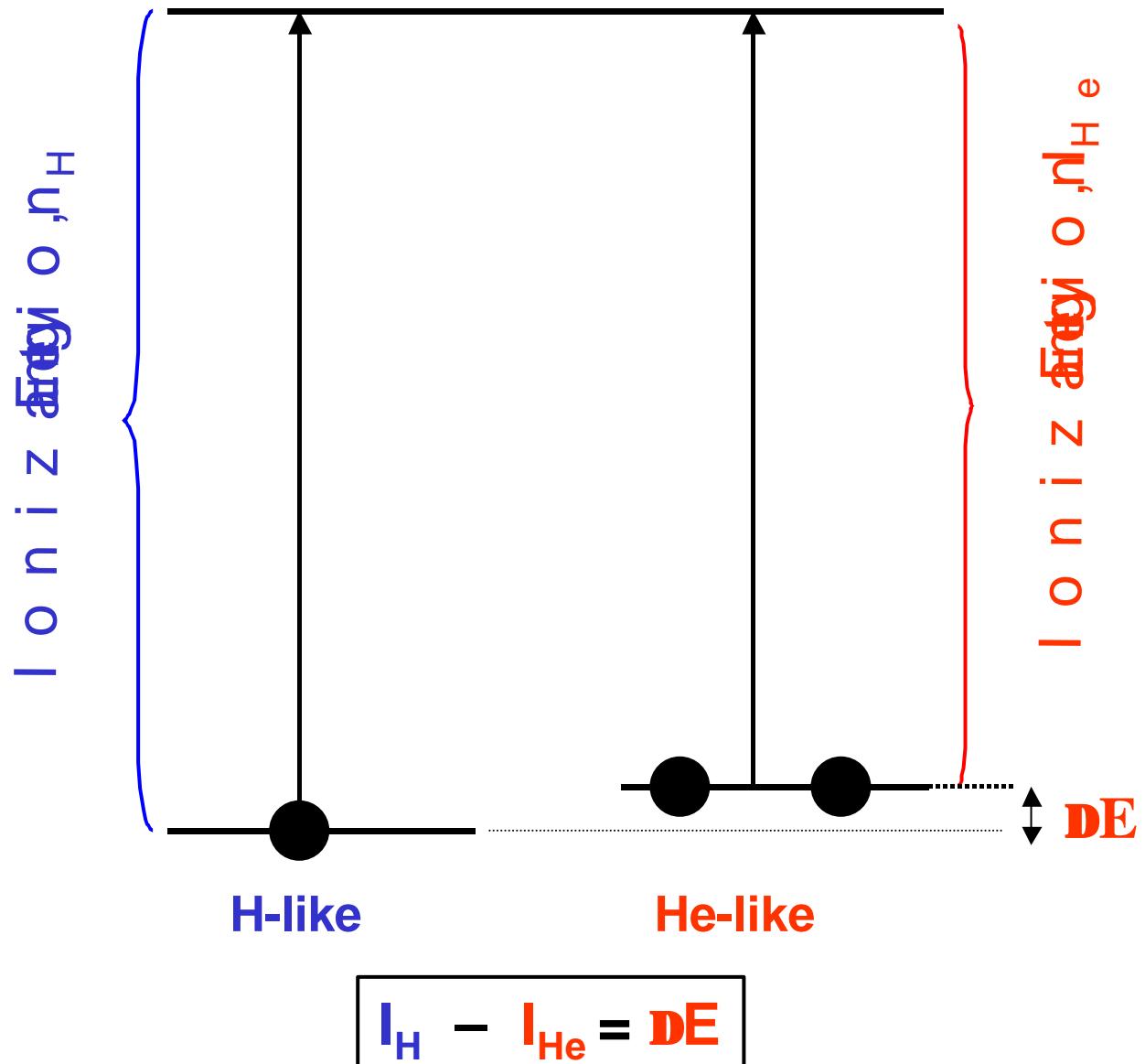
a,b) Non-Radiative QED
+1.3 eV [U⁹⁰⁺]
0.06%

c,d) Two-Electron Self Energy
-9.7 eV [U⁹⁰⁺]
0.4%

e,f) Two-Electron Vacuum Polarization
+2.6 eV [U⁹⁰⁺]
0.1%

Electron-Electron Interaction in Strong Fields

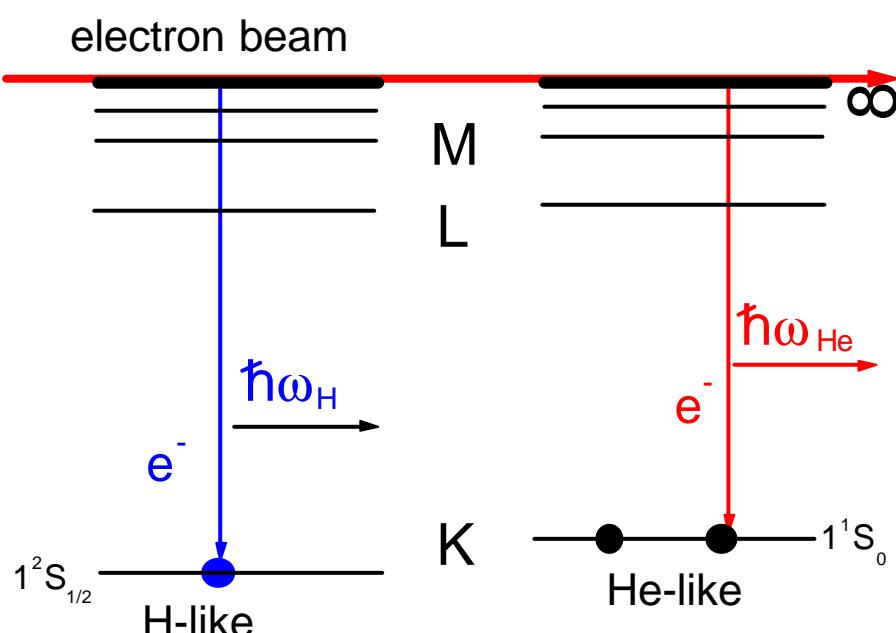
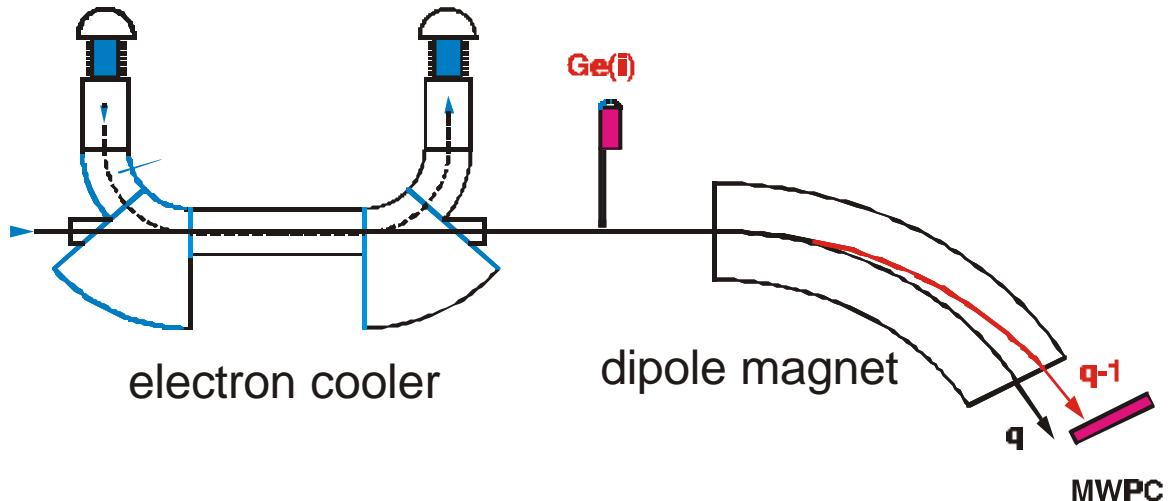
Continuum



ΔE : Two-Electron Contribution to the Binding Energy in the He-like System

Goal of the experiment

To probe the first time higher order QED (in a)
corrections in the domain of high-Z systems



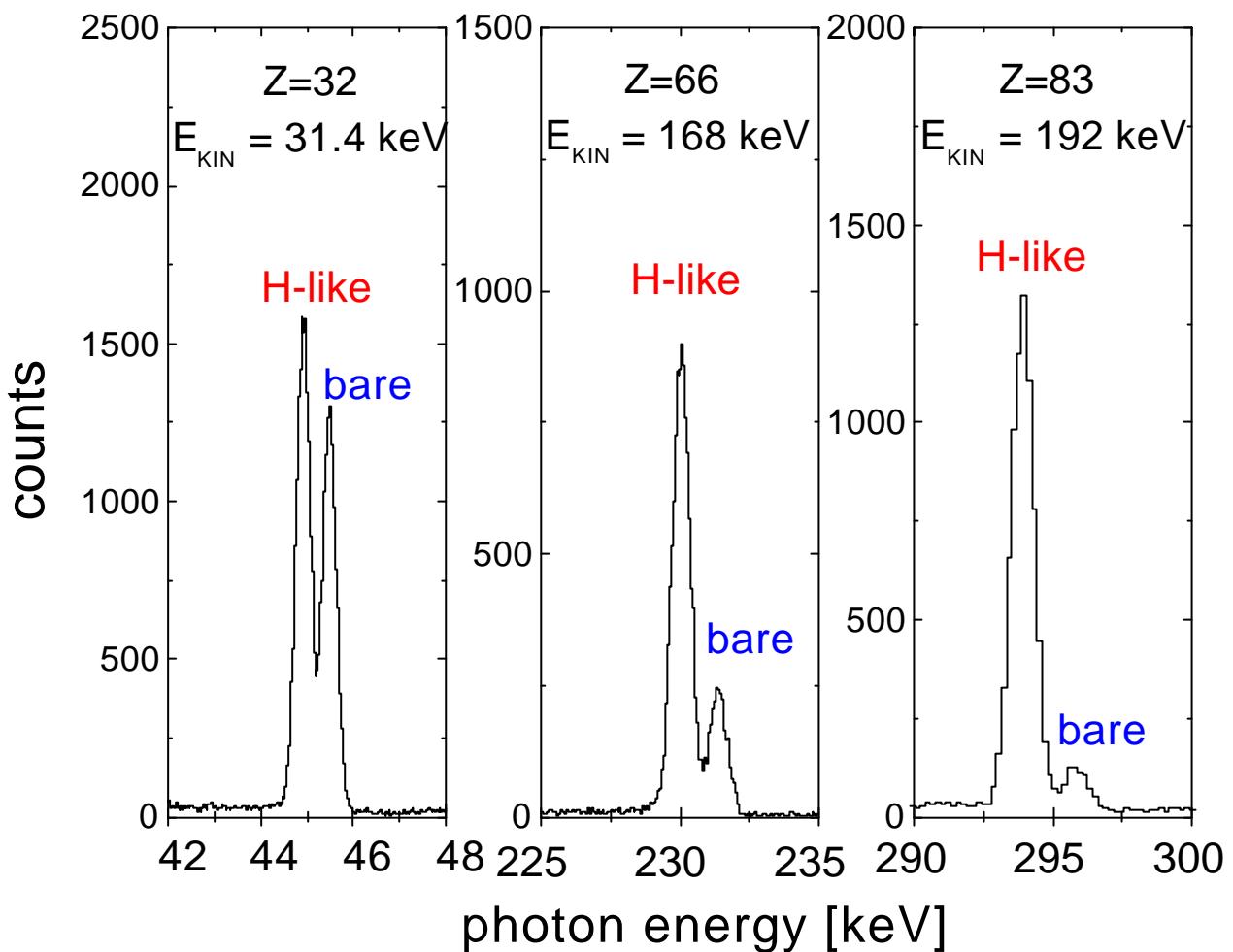
$$E_{KIN} + I_H = \hbar\omega$$

$$E_{KIN} + I_{He} = \hbar\omega$$

$$\Delta E(\hbar\omega - \hbar\omega) = I_H - I_{He}$$

Th. Stöhlker, APAC Conference 2001

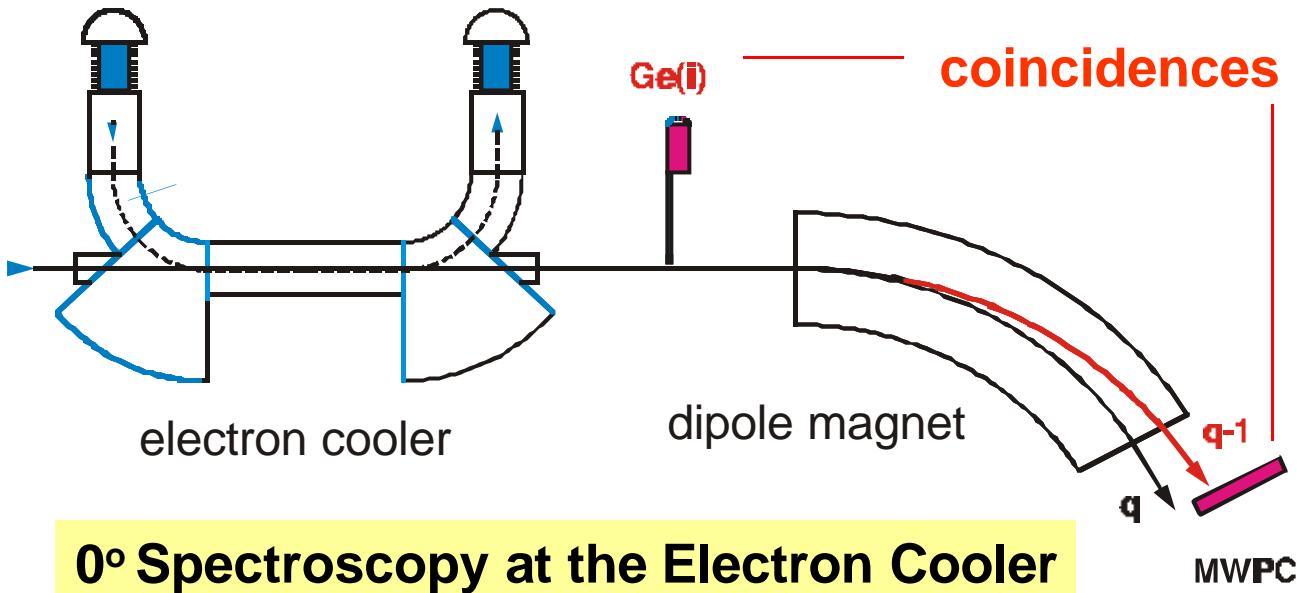
result



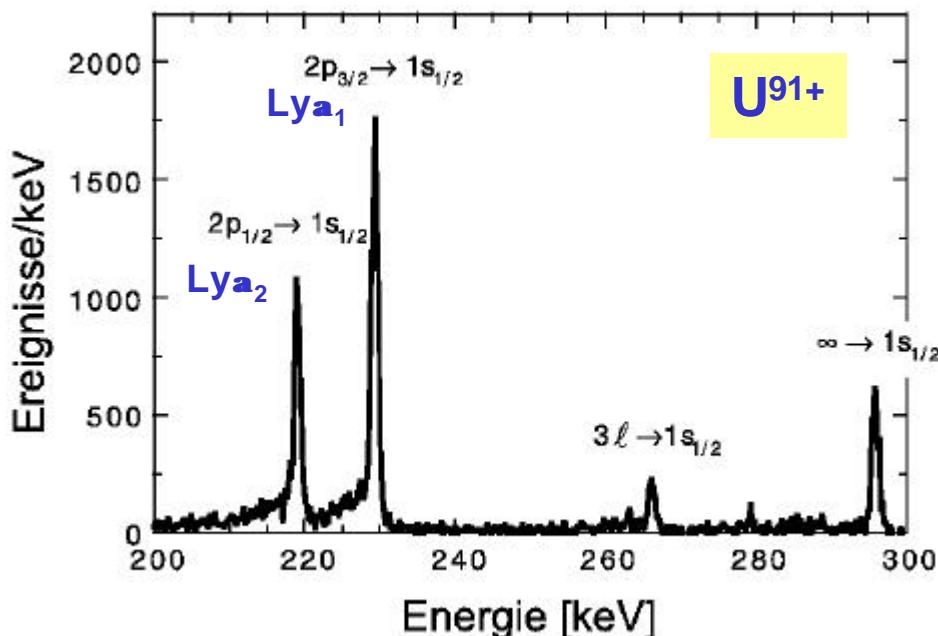
Z	32	54	66	74	76	83
Exp. [eV]	562.6 ± 1.6	1027.2 ± 3.5	1341.6 ± 4.3	1568.9 ± 15	1608 ± 20	1876 ± 14
$\Delta E/E$	2.8×10^{-3}	3.4×10^{-3}	3.0×10^{-3}	1×10^{-2}	1.2×10^{-2}	0.7×10^{-3}

Results are only limited by counting statistics

Th. Stöhlker, APAC Conference 2001

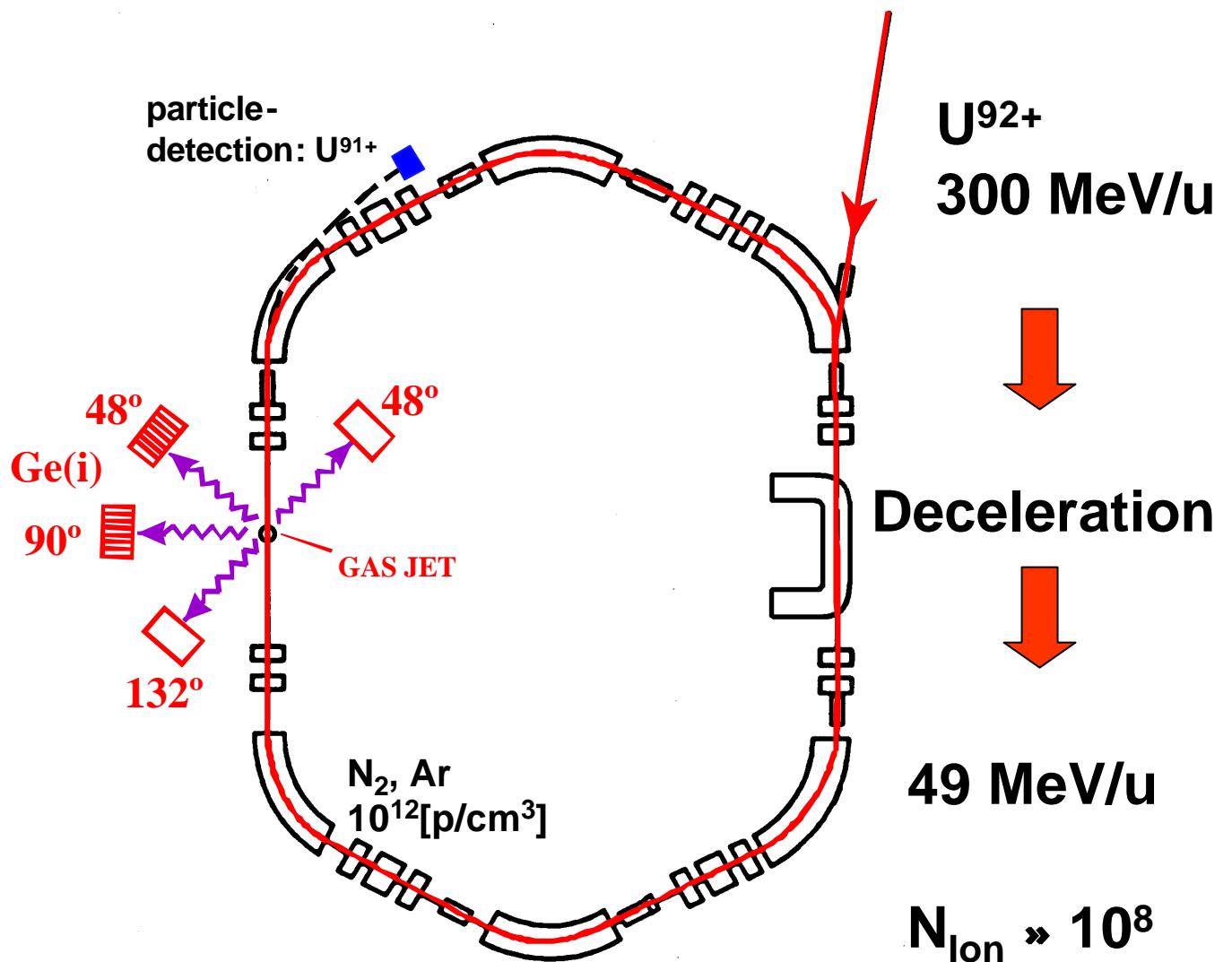


- strongest blue shift
 $b \gg 0.65 \cdot p_{\text{lab}} \gg 2.2 \cdot E_{\text{proj}}$
- Δq_{LAB} not critical, Doppler width has its minimum
- Uncertainty is caused by Δb



H.F. Beyer et al., 1995
 Th. Stöhlker, APAC Conference 2001

X-Ray Spectroscopy at the Jetttarget

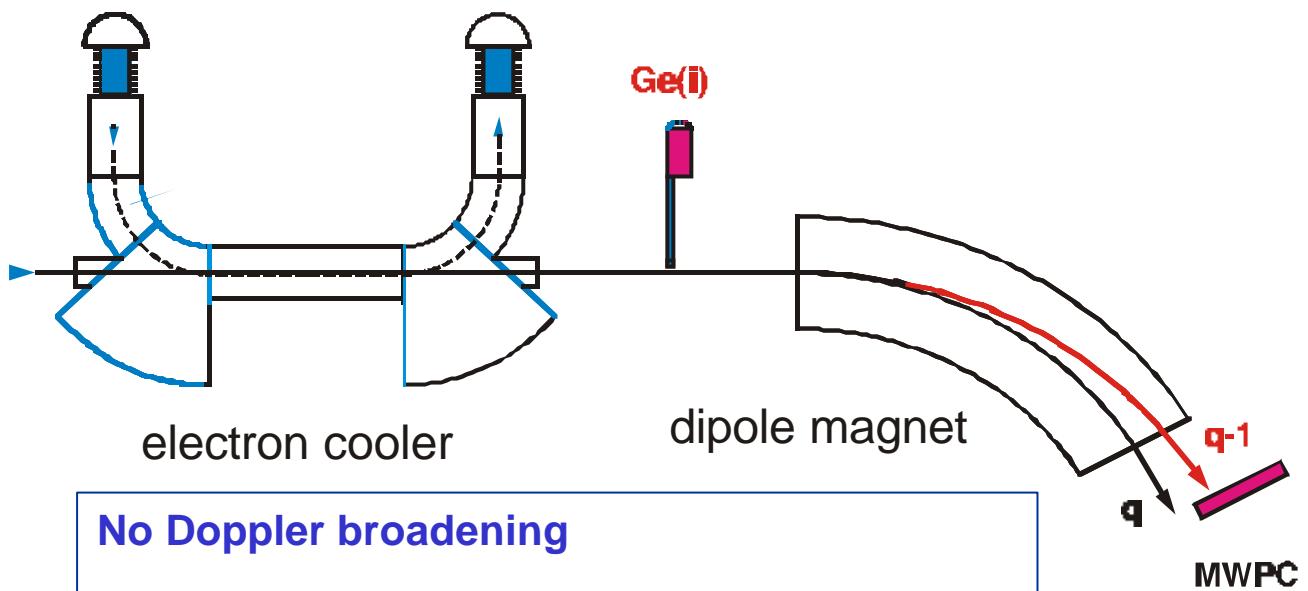


$$E_{\text{lab}} = \frac{E_{\text{proj}}}{\hat{a} \cdot (1 - \hat{a} \cdot \cos \hat{\epsilon}_{\text{lab}})}$$

E_{lab} : Photon energy in Laboratory system

E_{proj} : Photon energy in emitter system

0 deg spectroscopy at the electron cooler



No Doppler broadening

practically, no uncertainties introduced by the geometry of the set-up

X-ray projectile coincidences

The relative measurement can be performed in sequential steps by changing frequently between bare and H-like uranium

Application of the deceleration mode of the ESR

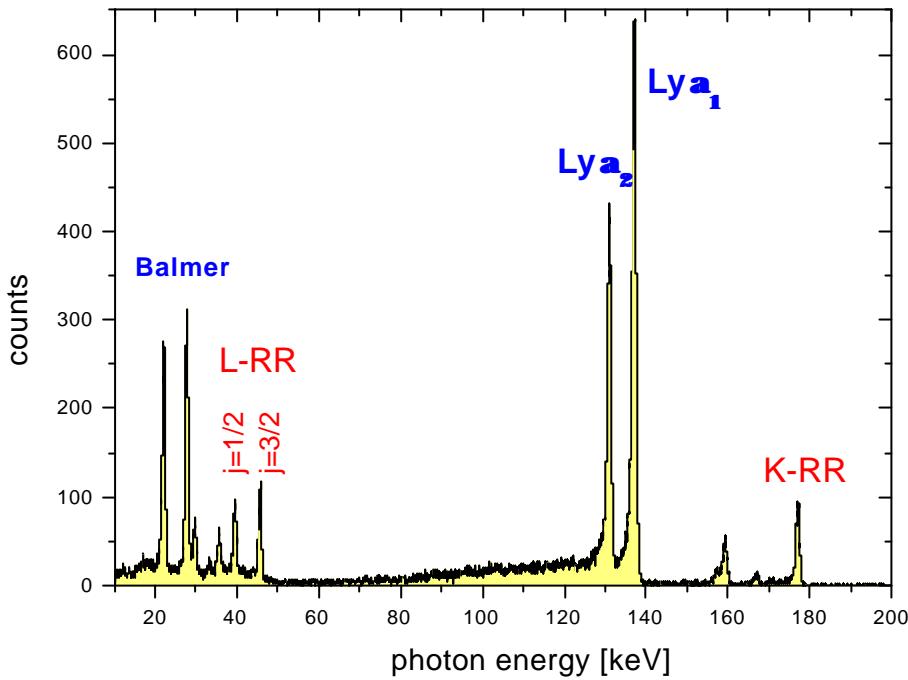
→ strongly reduced bremsstrahlung uncertainty in b not critical

To probe the QED corrections for the ground state in helium like uranium, a relative accuracy of only $\Delta E/E \approx 10^{-3}$ is required

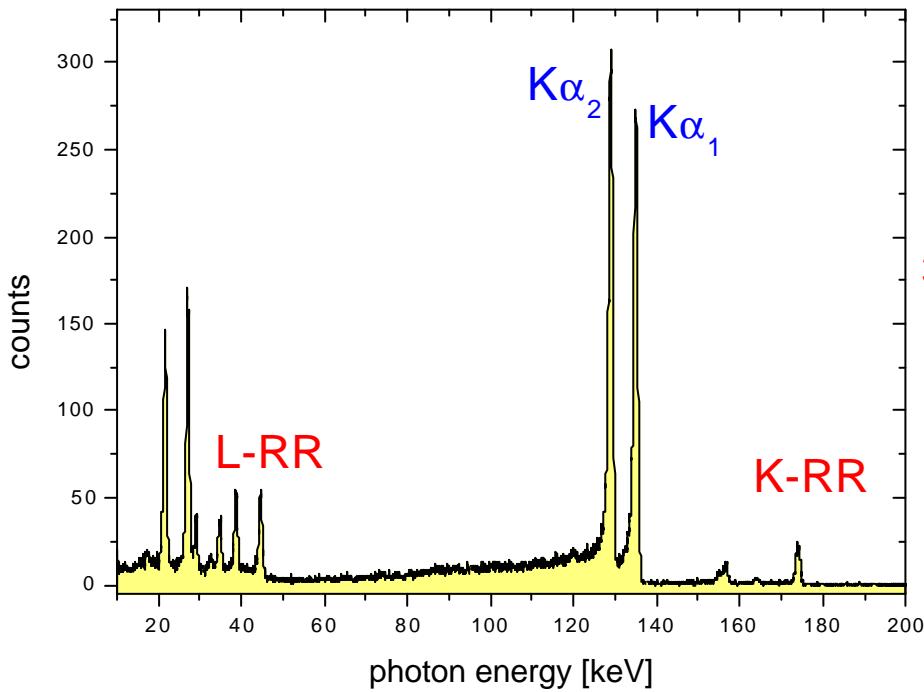
For uranium, we like to measure the 2.2 keV energy splitting with an accuracy of better 5 eV ($\Delta E/E \gg 2 \times 10^{-3}$)

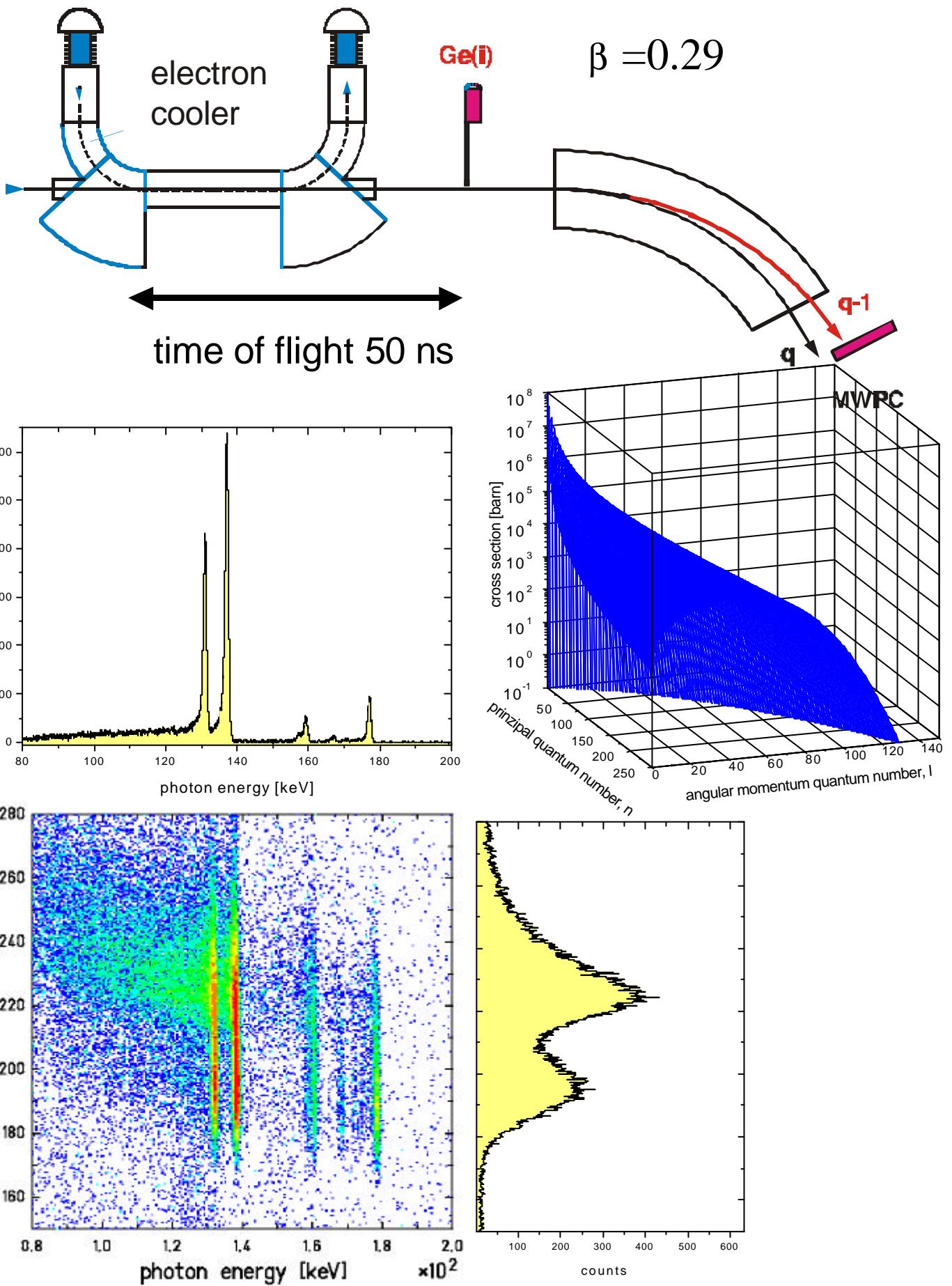
Th. Stöhlker, APAC Conference 2001

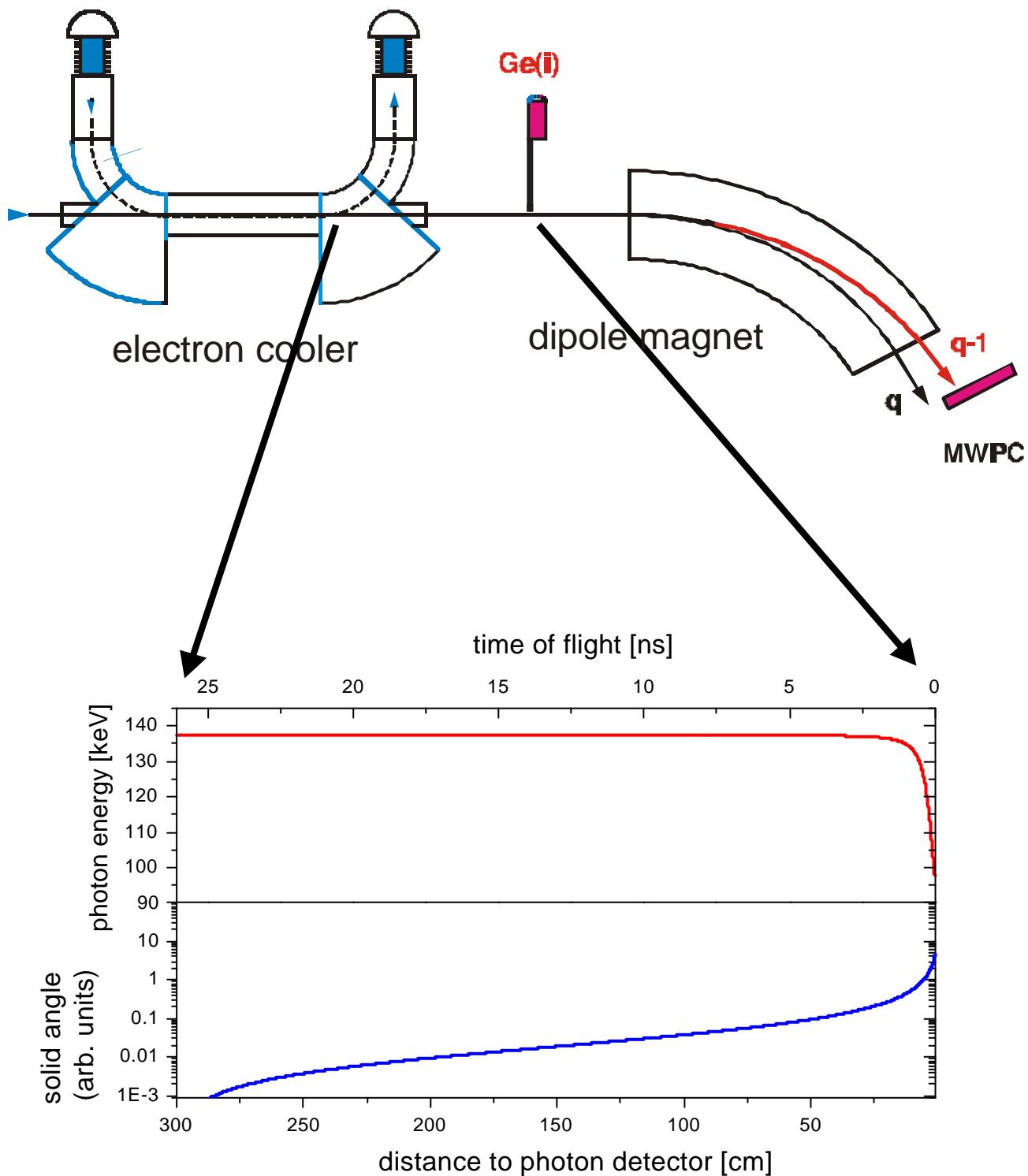
H-like uranium

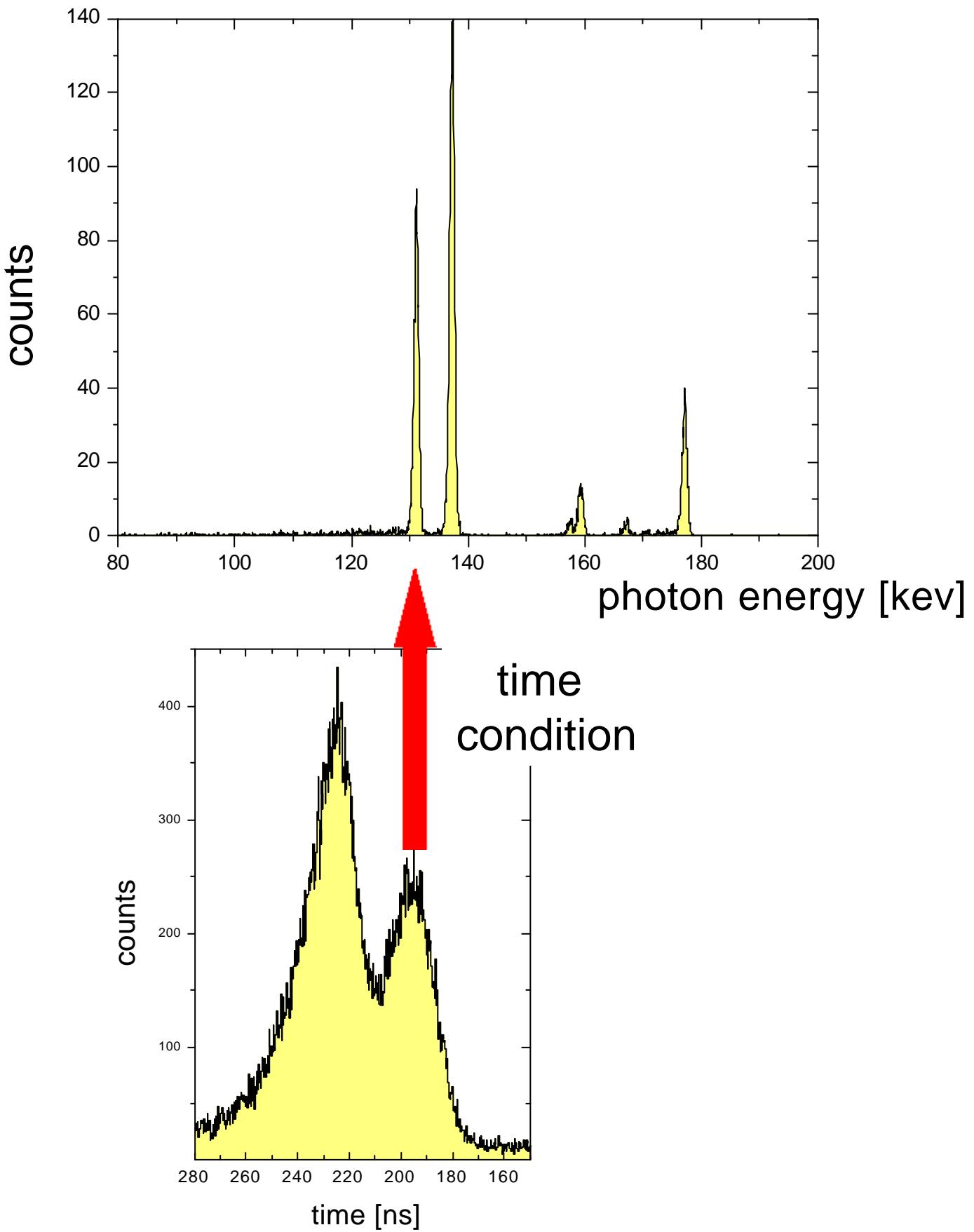


He-like uranium

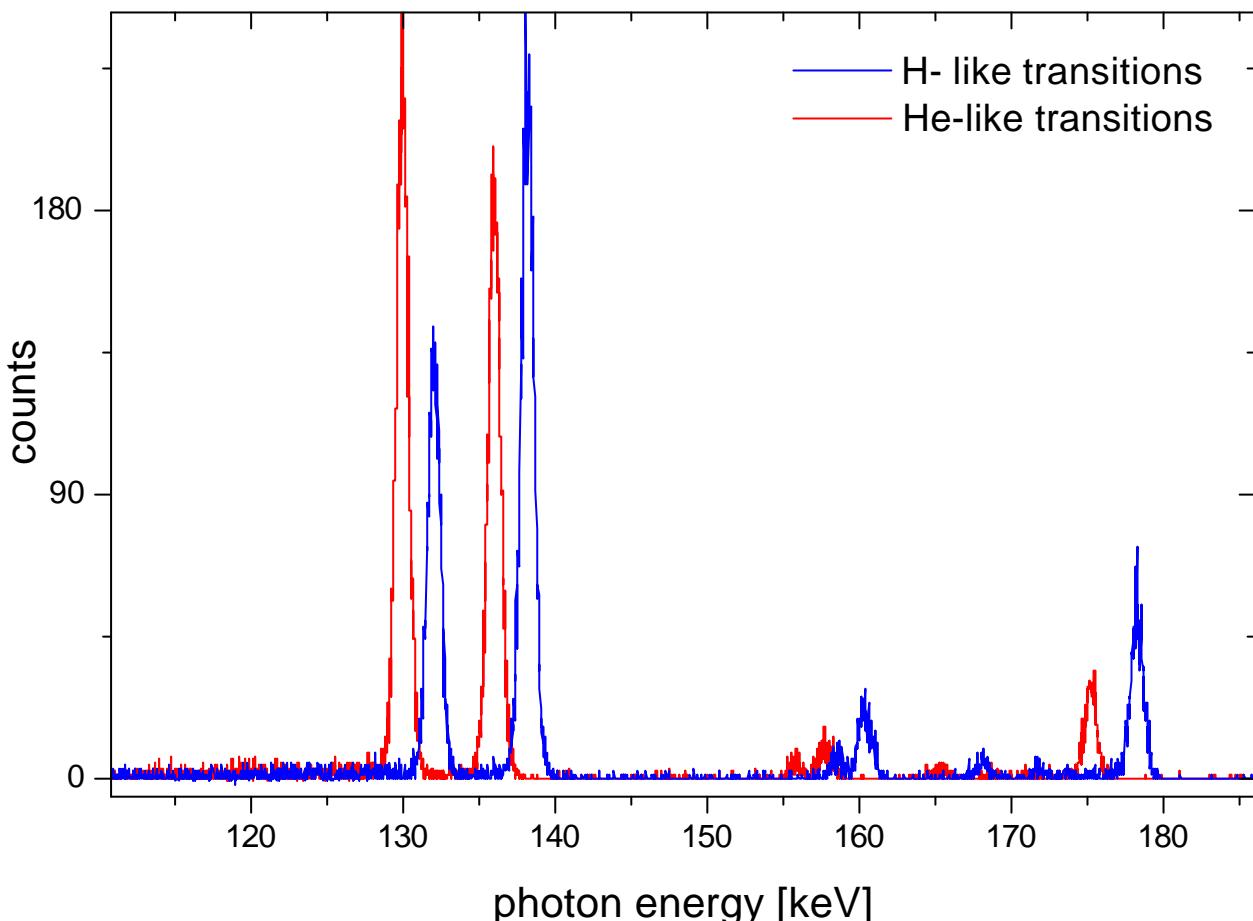








Relative measurement at electron cooler for decelerated ions



Estimated statistical uncertainty
for RR into H- and He-like uranium: ≈ 9 eV

Uncertainty caused by doppler shift:

Additional systematic errors: ?

Summary

- The ESR provides ideal conditions for precise studies of the atomic structure and collision of high-Z ions (at the jet target as well as at the electron cooler)
 - *determination of relative ionization potentials*
 - *Lamb shift experiments*
- The study of elementary atomic processes for highly-charged heavy ions via their time-reversal
 - *identification of spin-flip transitions for photoionization*
 - *interference between E1 and M2 transition amplitudes*

Outlook

- Decelerated ions
- photon correlation spectroscopy
- polarisation studies
- Recoil momentum spectroscopy and electron spectroscopy
- Combination of electron and photon spectroscopy
- Laser assisted experiments