

heavy ion collisions and hydrodynamics

for $T > 200$ MeV in 2-flavor QGP $n_{\text{parton}} > 4/\text{fm}^3$ and with typical perturbative cross sections $\lambda < 0.8$ fm
rescattering between particles formed in primary collisions may lead to local thermal equilibrium rapidly
treat system as particle fluid using language and tools of hydrodynamics

$\partial_{\mu} T^{\mu\nu} = 0$ $\partial_{\mu} j^{\mu} = 0$ with energy-mom tensor $T^{\mu\nu}$ and 4-current of cons. charge j^{μ}

for ideal fluid: $T^{\mu\nu} = (\varepsilon + p)u^{\mu}u^{\nu} - pg$ and $j^{\mu} = n_i u^{\mu}$

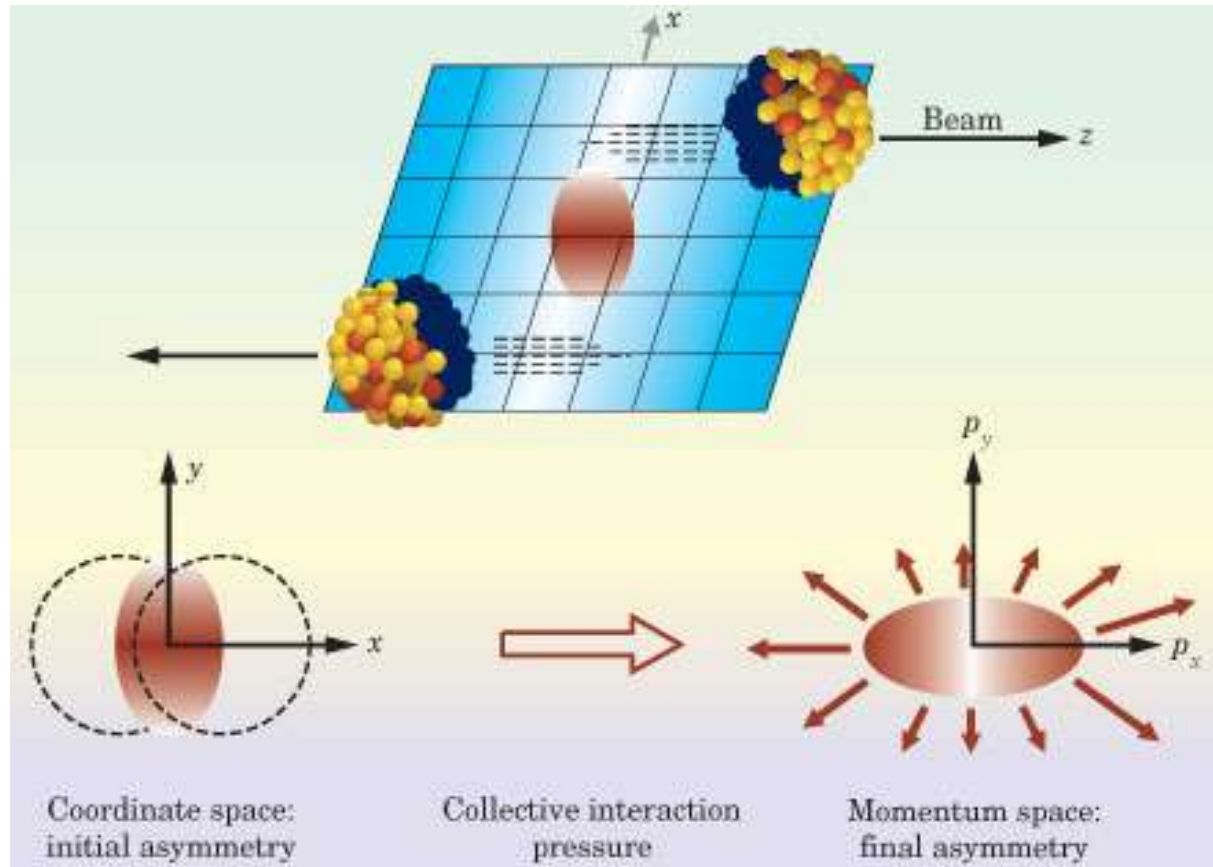
ε : energy density p : pressure u^{μ} : flow 4 velocity
generally all fields functions of x

generally only **EoS** and **initial condition** needed to calculate evolution

EoS: $p = p(\varepsilon, n_1, \dots, n_n)$ connection pressure – densities

initial cond.: in ideal fluid expansion isentropic, final state multiplicity gives initial entropy, pick volume \rightarrow system completely determined

Azimuthal Anisotropy Parameter v_2



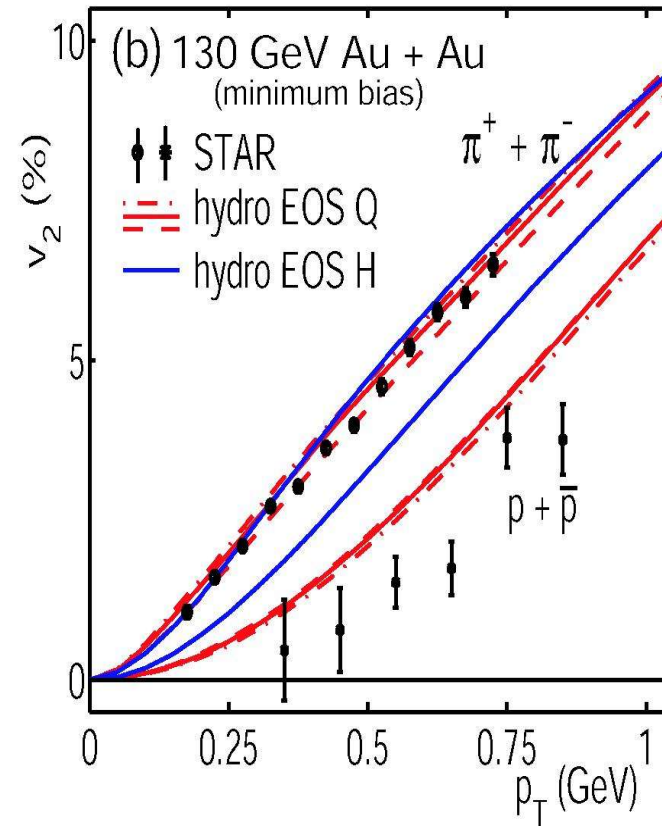
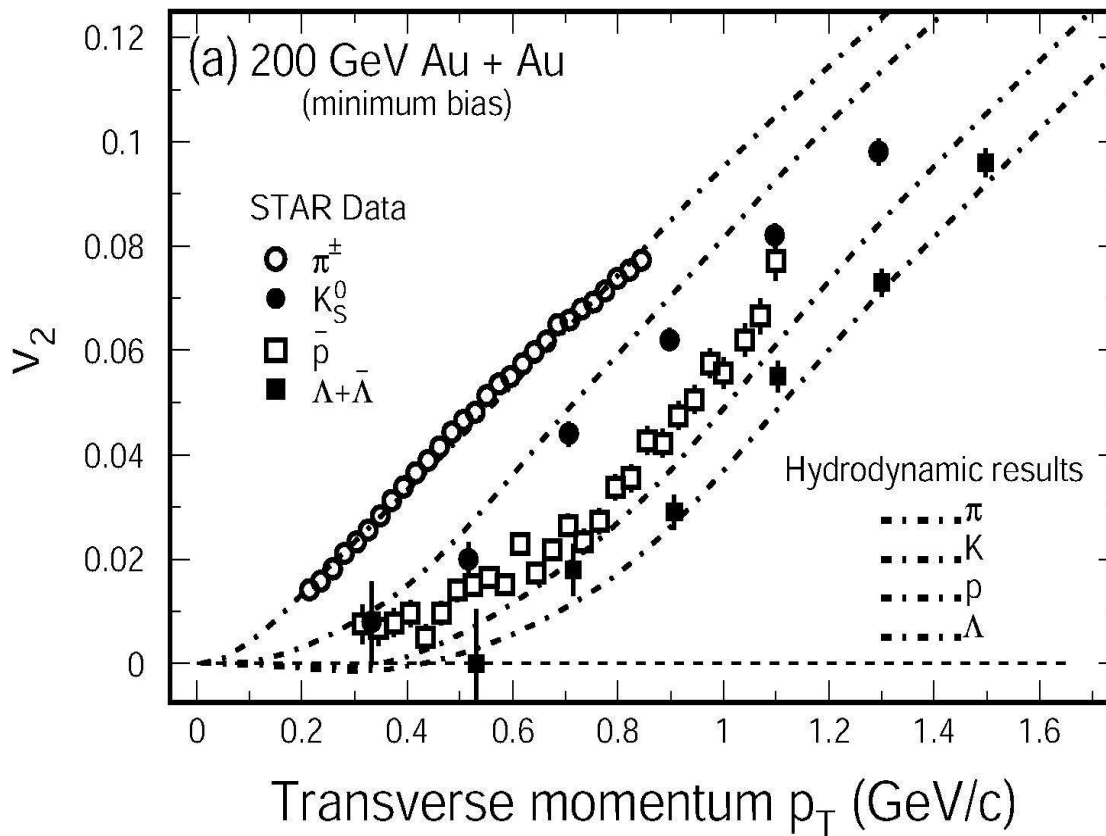
Fourier decomposition of momentum distributions rel to reaction plane:

$$\frac{dN}{dp_t dy d\phi} = N_0 \cdot \left[1 + \sum_{i=1} 2 v_i(y, p_t) \cos(i\phi) \right]$$

quadrupole component v_2
“elliptic flow”

elliptic flow for different particle species and p_t

$v_2 = \langle \cos(2\phi) \rangle$ where angle ϕ is relative to reaction plane



QGP
hadron gas

data: STAR PRL92 (2004) 052302

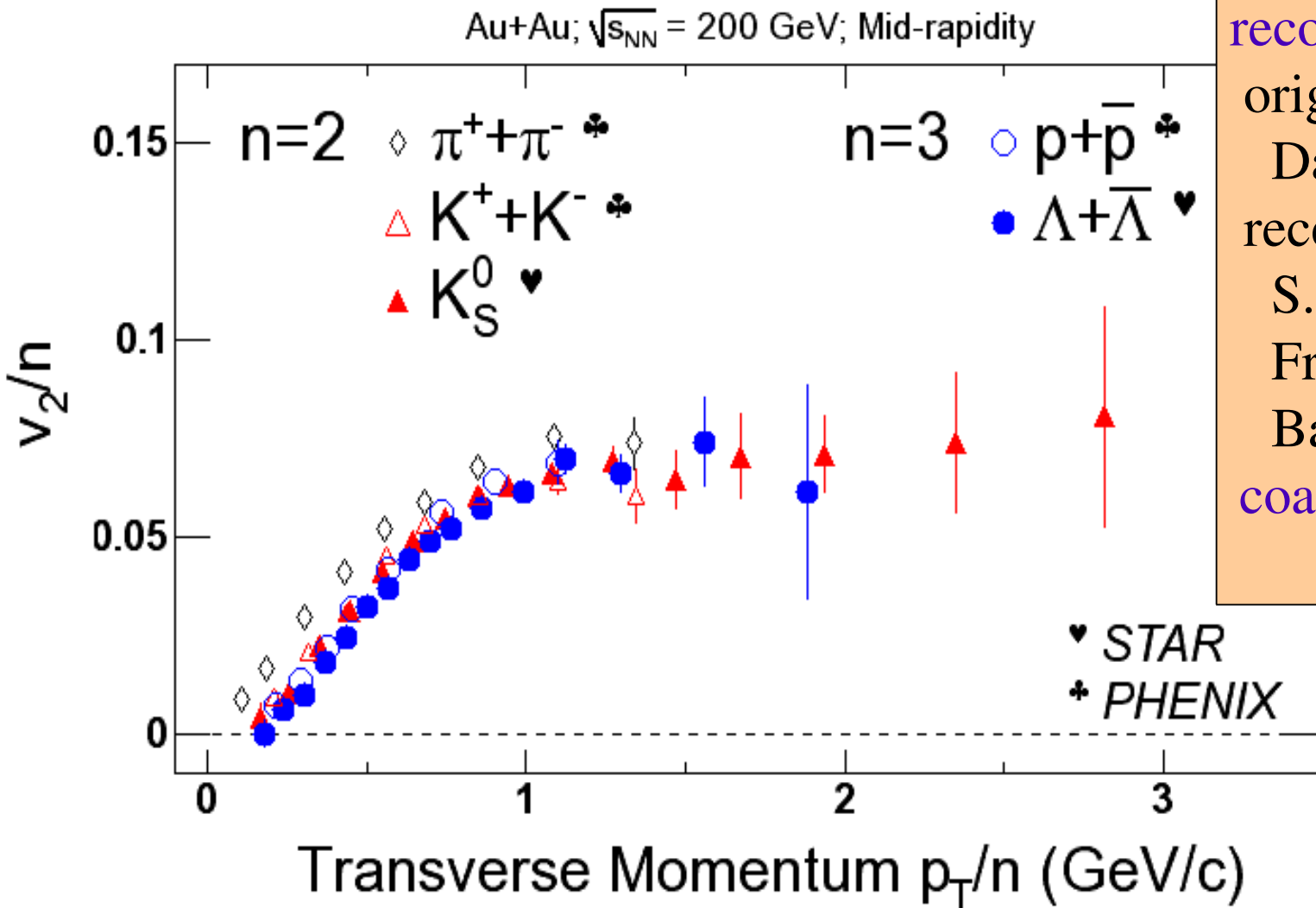
and PRL 87 (2001) 182301

hydro: P. Huovinen et al. PL B503 (2001) 58

and priv. comm.

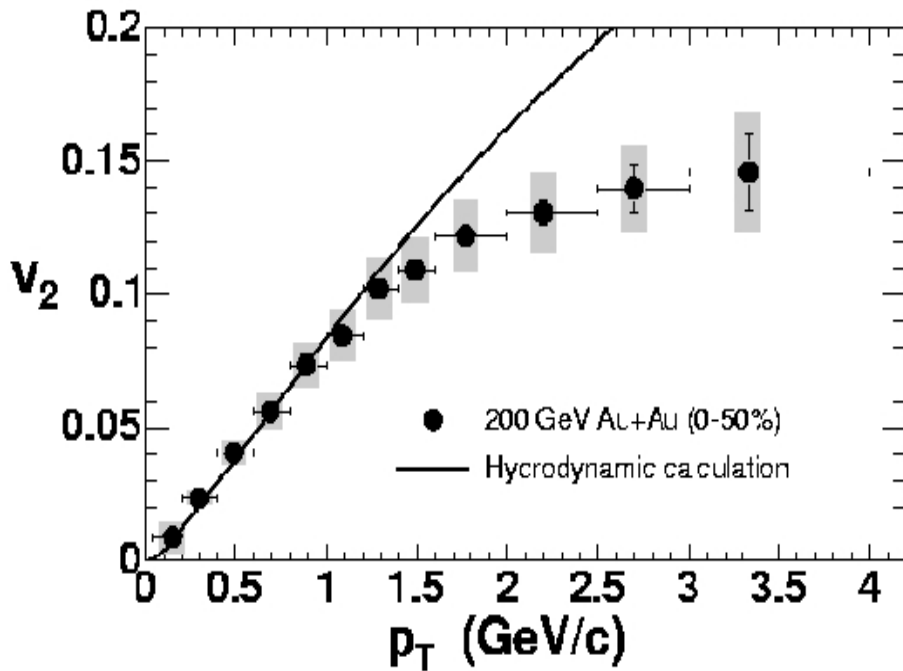
well described by hydrodynamics
sensitivity to EOS

valence quark scaling of elliptic flow



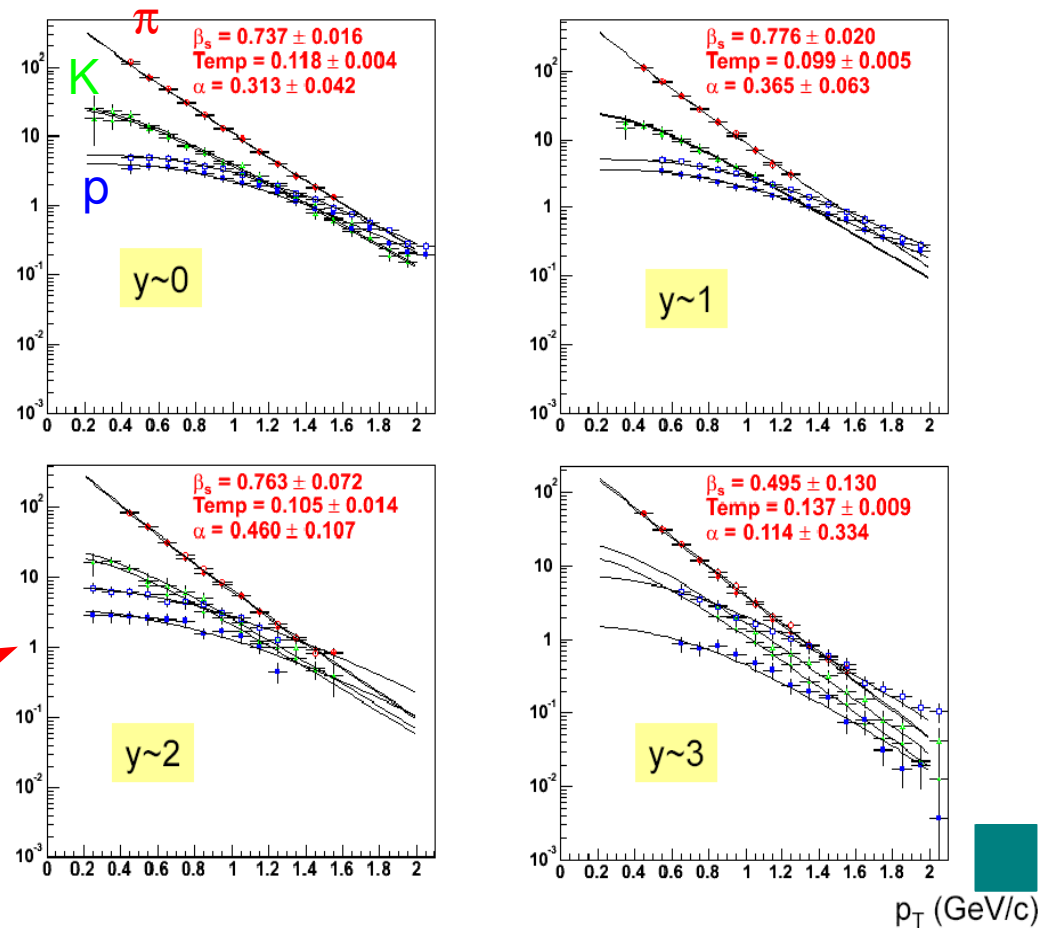
recombination of quarks:
 original idea:
 Das & Hwa 1977
 recent connection to data:
 S. Voloshin 2003
 Fries, Müller, Nonaka,
 Bass 2003
 coalescence model works
 for $p_t > 2-3$ GeV/c

typical data: p_t spectra and elliptic flow



PHOBOS, nucl-ex/0410022

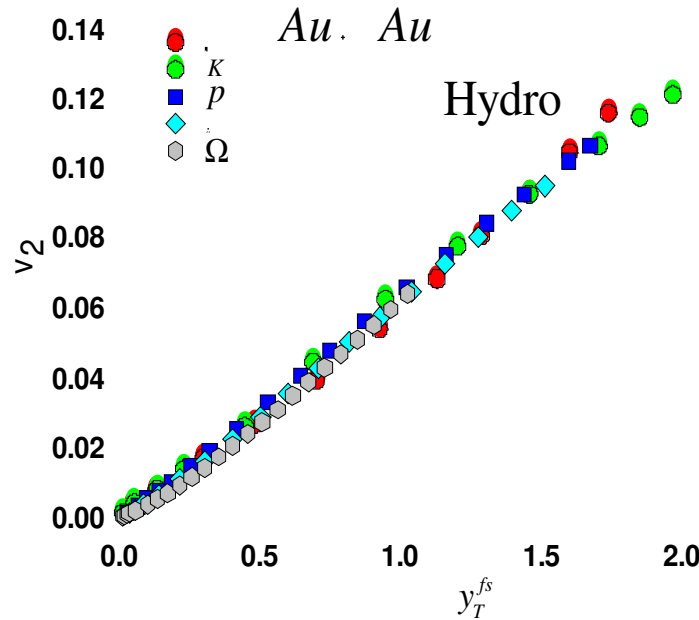
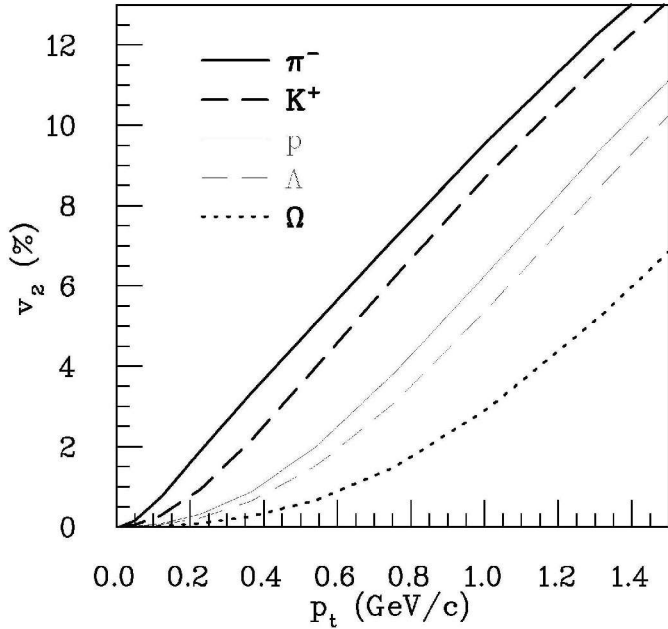
inv. slope constant grows with mass – typical for expanding fireball (Hubble expansion)



R. Debbé, BRAHMS, Proc. ICPAQGP 2005, Kolkata

particular scaling of hydrodynamics with mass – observed in hydrodynamic models as well as data

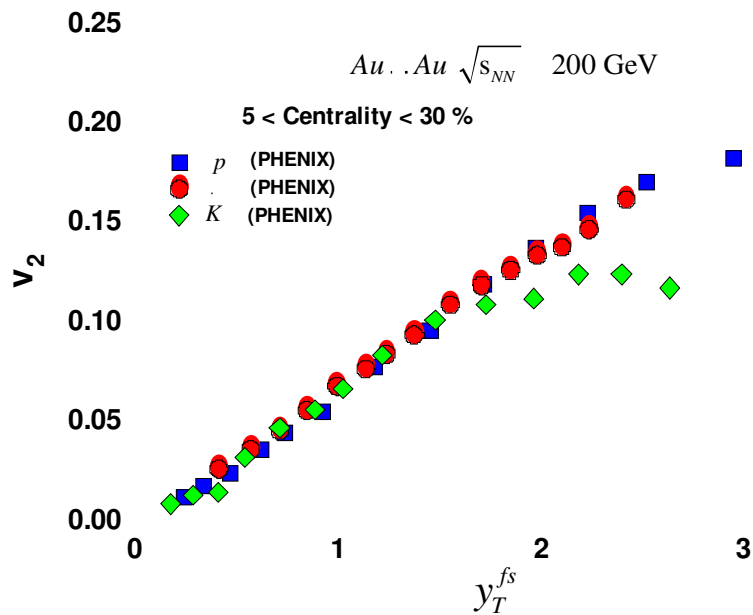
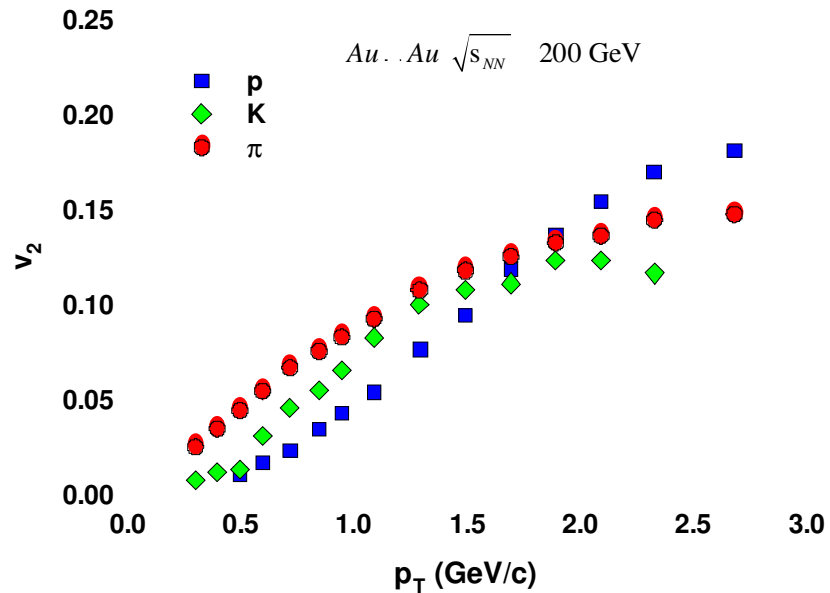
P. Huovinen, P.F. Kolb .. Phys. Lett. B503:58 (2001)



hydro scaling realized:
Csörgö et al., PRC67
(2003)034904

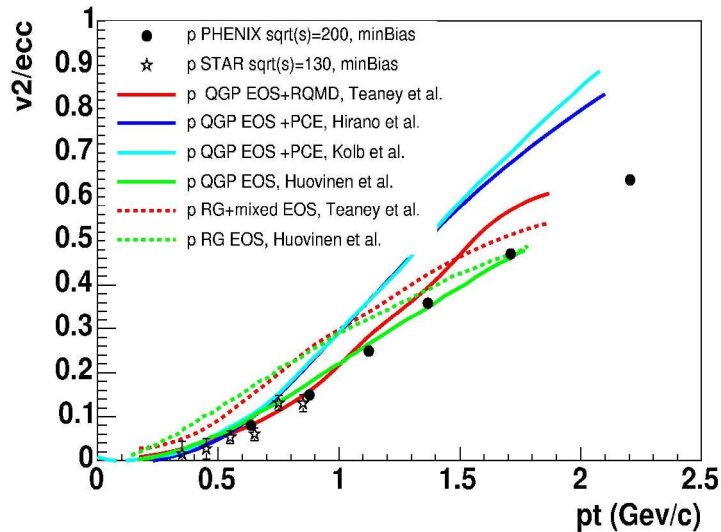
$$p_T \rightarrow y_T = \sinh^{-1}(p_T/m)$$

application to model
and data: R.Lacey
and A.Taranenko
nucl-ex/0506019

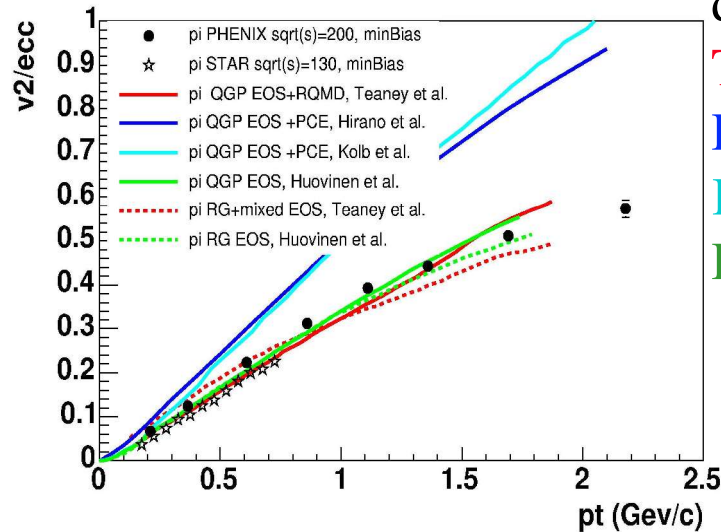


ideal hydrodynamics describes spectra and elliptic flow

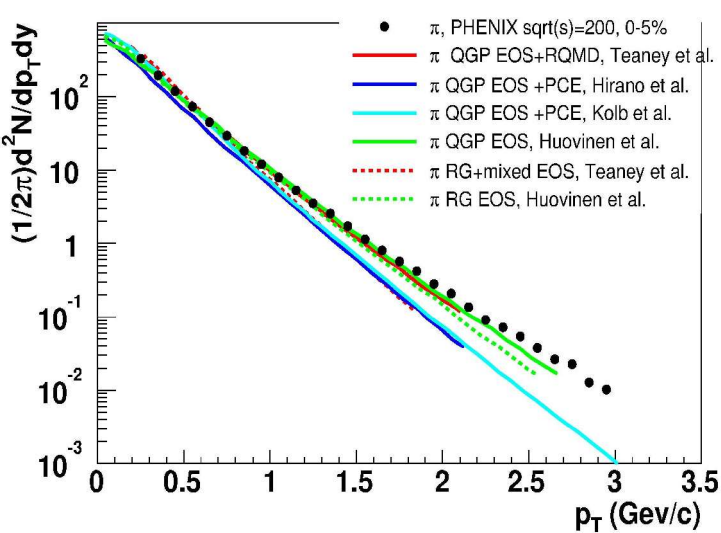
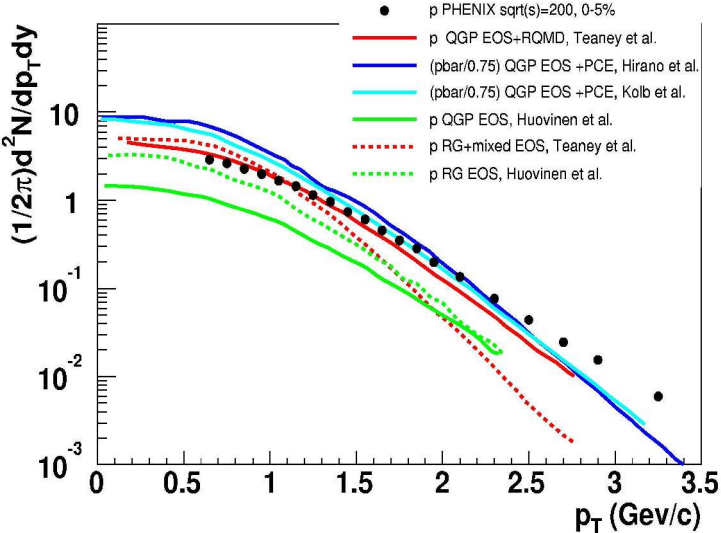
proton



pion



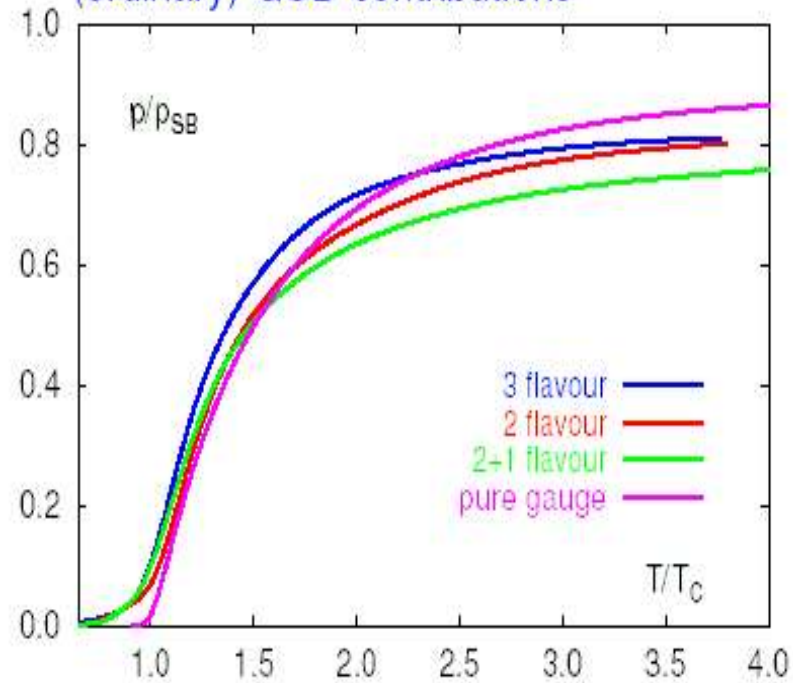
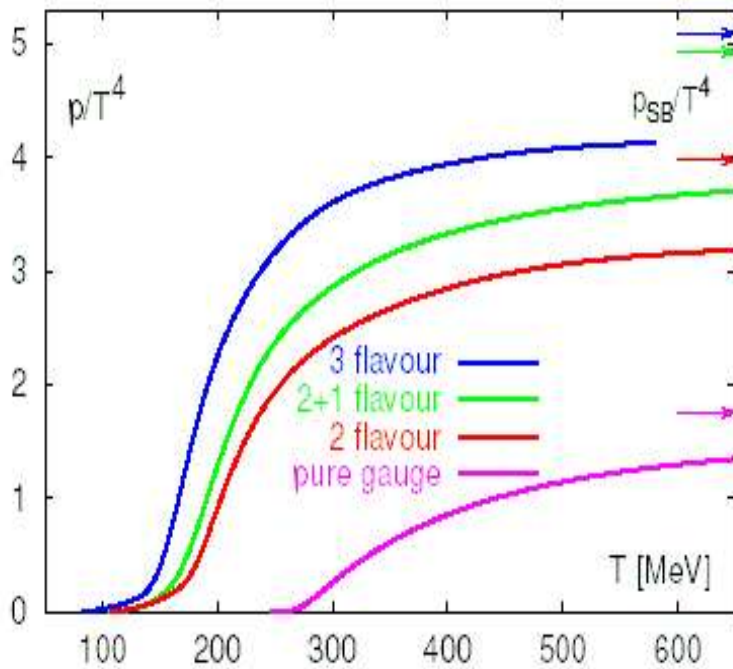
different hydrodyn. models:
 Teaney (*w/ & w/o RQMD*)
 Hirano (*3d*)
 Kolb
 Huovinen (*w/ & w/o QGP*)



works up to
 about 2 GeV/c
 „requires
 strong interactions at
 short times“
 (role of CGC?)

the pressure indicates non-ideal behavior

$T \gtrsim (2-3)T_c$: deviations from ideal gas understood in terms of non-perturbative (ordinary) QCD contributions

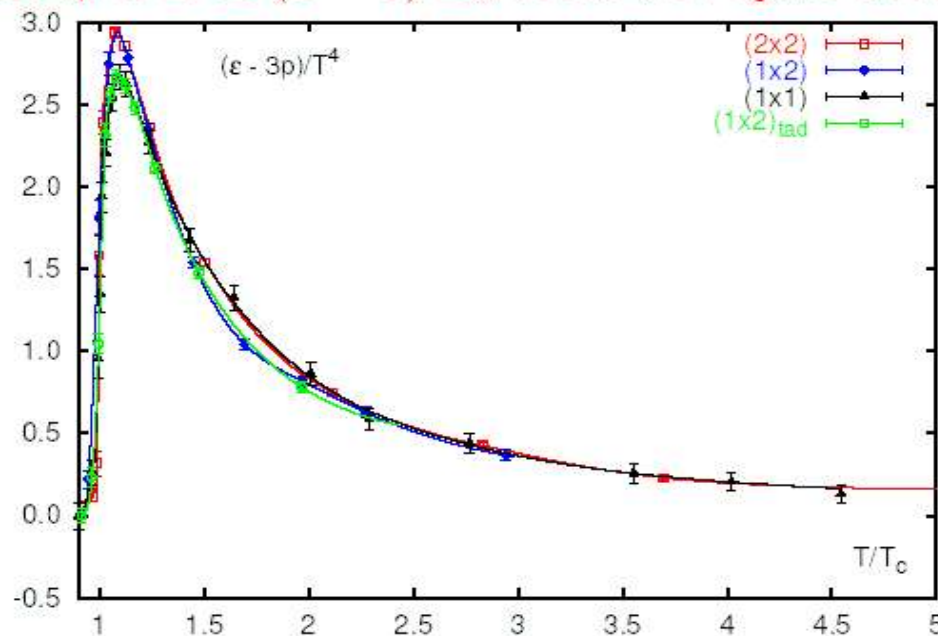


FK, E. Laermann and A. Peikert, Phys. Lett. B478 (2000) 447.

strong interactions near T_c

Deviation from ideal gas (pure gauge) – suggests strong interactions –

evidence for significant deviations from ideal gas behavior;
at least up to $T \simeq (2 - 3) T_c$; i.e. in the regime accessible to RHIC



colored bound states above T_c ??

- success of hydrodynamic description – short mfp
- from lattice: QGP is strongly coupled near T_c
- but: ideal liquid interpretation is a guess
- no computation of viscosity from LQCD yet
- no computation of viscosity in hadron gas
- no systematic study of viscous hydrodynamics for RHIC data

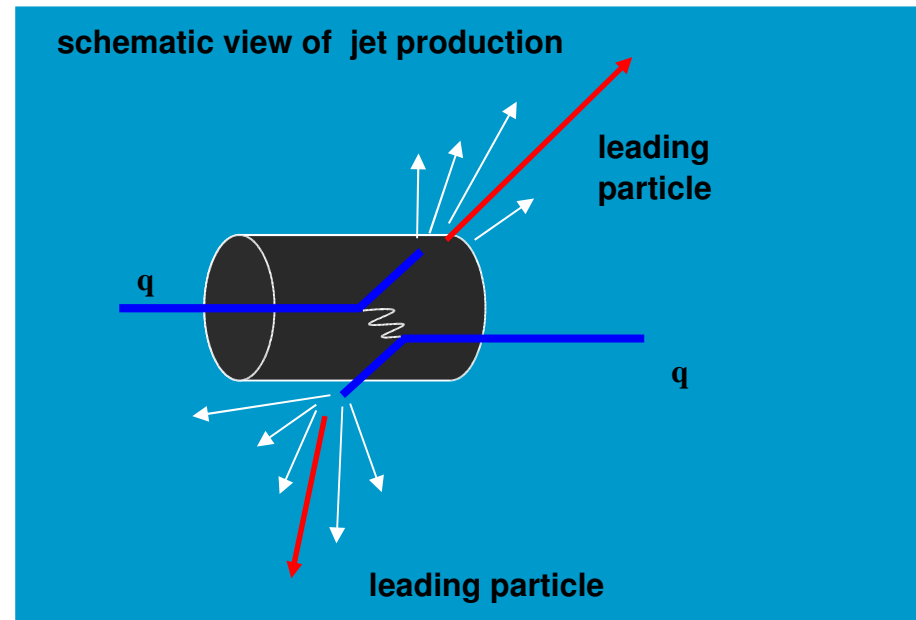
many open questions

Jet quenching

- suppression of high p_t particles in AA relative to pp collisions
- disappearance of jet-like correlations
- connected to large gluon density in hot (QGP) fireball

Jets of hard partons as probe of the hot medium

- **Hard parton scattering observed via leading particles**
- **Expect strong $\Delta\phi=\pi$ azimuthal correlations**



However, the scattered partons may lose energy (~ several GeV/fm) in the colored medium

- **momentum reduction (fewer high p_T particles in jet)**
- **no jet partner on other side**

Jet Quenching

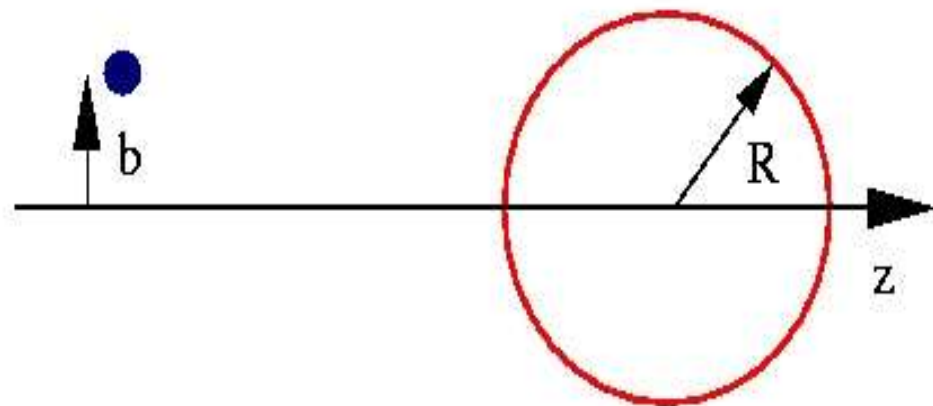
From pp to AA collisions

all valid for hard scattering only

start with normalized nuclear density, $\rho(b, z)$

$$\int \rho(b, z) dz db \cdot 2\pi b = A$$

$$R = r_0 A^{1/3}$$



Thickness function:

- pA collisions:

$$T_A(b) = \int_{-\infty}^{+\infty} dz \rho(b, z)$$

with $\int d^2b T_A(b) = \int_0^\infty 2\pi b db T_A(b) = A$

Note: $\sigma_{NN} \cdot T_A(b) = \langle N_{coll}^{pA}(b) \rangle$

sharp spheres: $T_A(0) = \frac{3}{2\pi R_A^2} A$

- AB collisions:

$$T_{AB}(b) = \int d^2b_1 d^2b_2 \delta^{(2)}(\vec{b} - \vec{b}_1 - \vec{b}_2) T_A(b_1) T_B(b_2)$$

with $\int d^2b T_A(b) = A \cdot B$

$$\sigma_{NN} \cdot T_{AB}(b) = \langle N_{coll}^{AB}(b) \rangle$$

sharp spheres: $T_{AB}(0) = \frac{9}{8} \frac{A^2}{\pi R_A^2}$

Thickness function continued

Implications:

$$\frac{d\sigma^{AB}}{dy}(\Delta b = b_2 - b_1) = \frac{d\sigma^{NN}}{dy} \int_{b_1}^{b_2} d^2b T_{AB}(b)$$

for minimum bias ($b = [0, R_A + R_B]$):

$$\frac{d\sigma^{AB}}{dy} = \frac{d\sigma^{NN}}{dy} \cdot A \cdot B$$

also:

$$\frac{dN^{AB}}{dy} = \frac{dN^{NN}}{dy} \cdot \sigma_{inelastic}^{NN} \cdot \frac{\int_{b_1}^{b_2} d^2b T_{AB}(b)}{\pi(b_2^2 - b_1^2)}$$

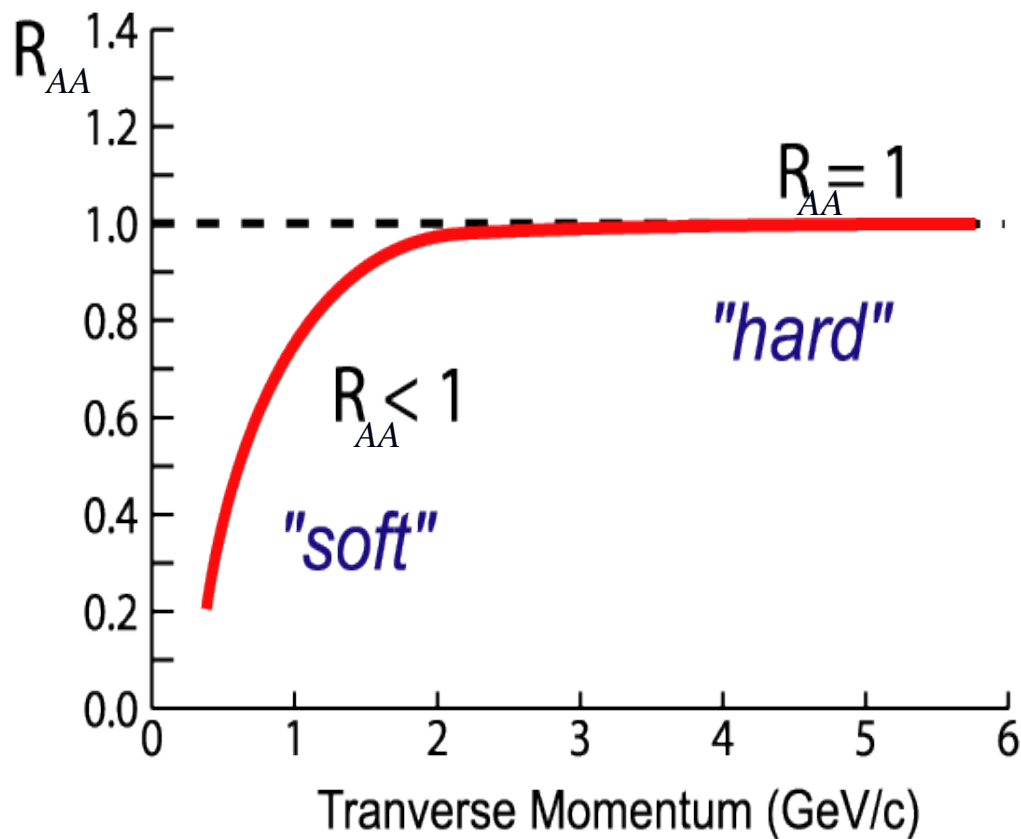
The R_{AA} function:

$$R_{AA}(b) = \frac{\frac{d^2 N^{AA}}{dp_t^2 dy}}{N_{coll}^{AA}(b) \cdot \frac{d^2 N^{NN}}{dp_t^2 dy}}$$

if hard scattering only:

$$R_{AA}(b) = 1$$

Qualitative expectations



no medium effects:

$R_{AA} < 1$ in regime of soft physics

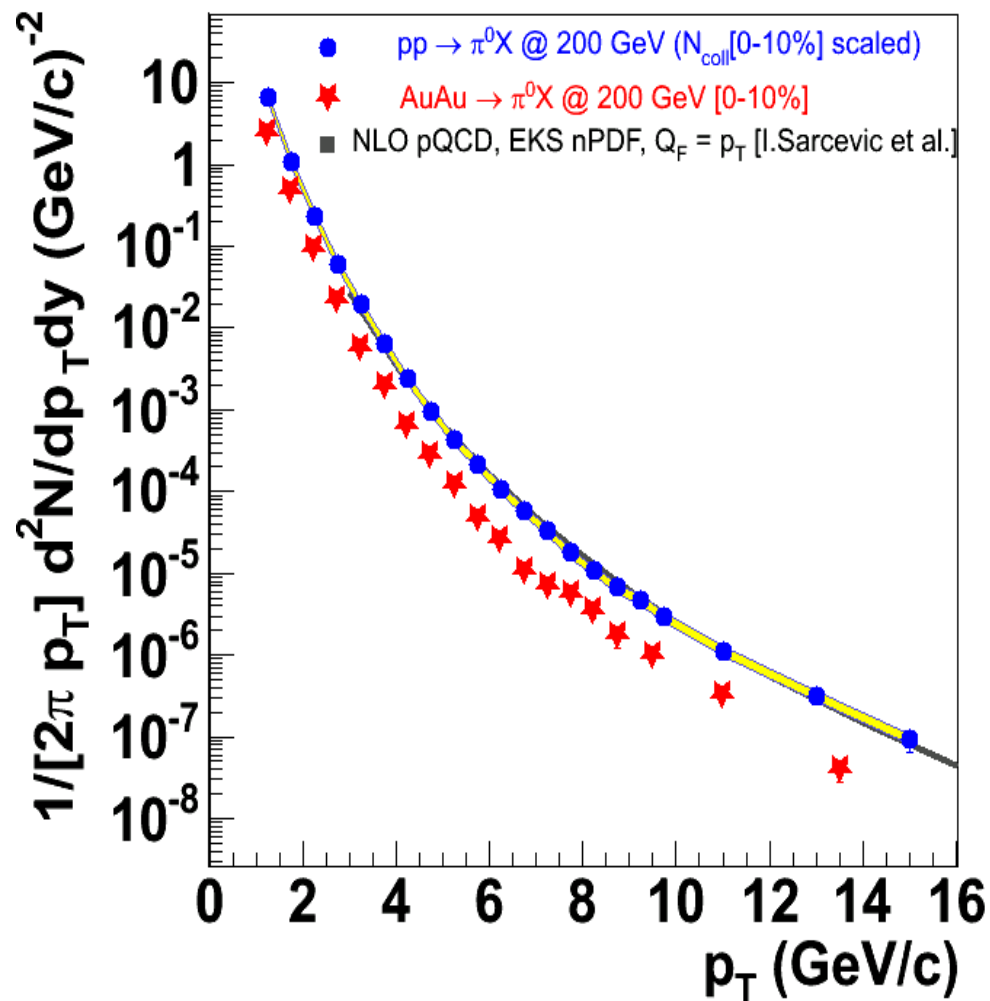
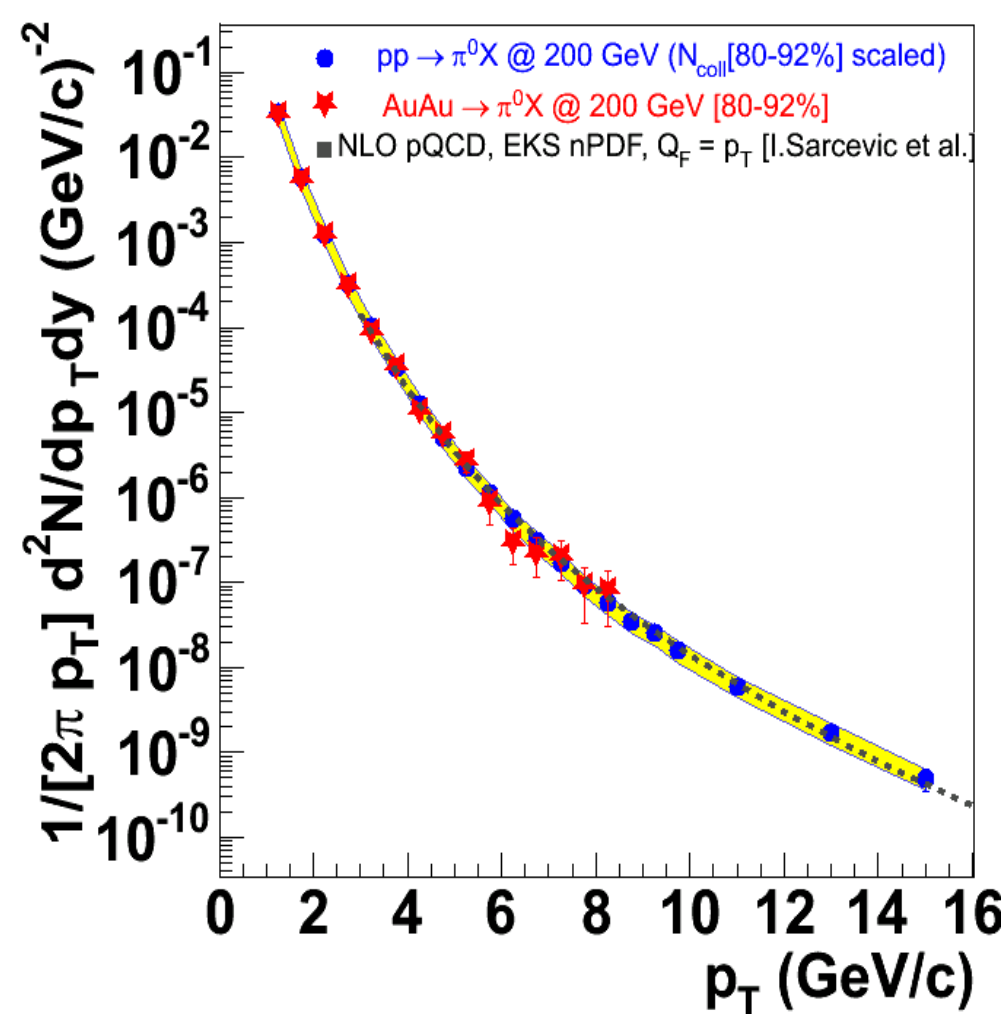
$R_{AA} = 1$ at high- p_T where hard scattering dominates

Suppression:

$R_{AA} \ll 1$ at high- p_T

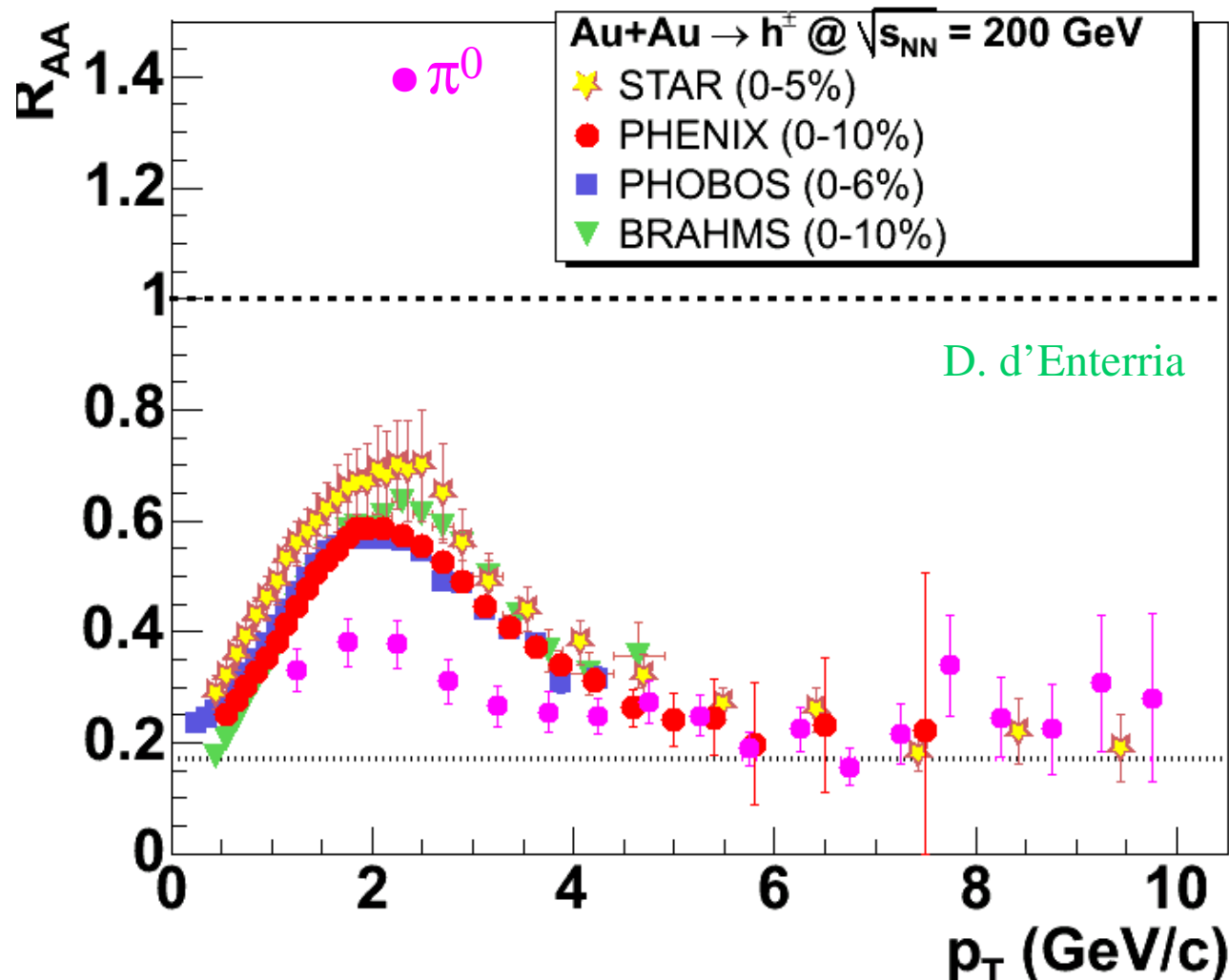
spectra suppressed at high p_t in AuAu relative to pp

proton data scaled to AuAu with appropriate number of binary collisions



high p_t suppression seen by all experiments

$$R_{AA} = \text{yield}(\text{AuAu}) / N_{\text{coll}} \text{ yield}(\text{pp})$$



- ★ all expts. see large suppression in AuAu
- ★ π^0 lower than h^\pm
- ★ no suppression in dAu rather

Cronin enhancement
→ medium effect, not incoming partons

→ reasonable agreement between 4 experiments

Suppression predicted due to energy loss of partons in hot matter “jet quenching”

H. Baier, Y.L. Dokshitzer, A.H. Mueller,
S. Peigne, D. Schiff, Nucl. Phys. B483
(1997) 291 and 484 (1997) 265

energy loss of high energy parton
traversing color charged medium ->

medium induced gluon radiation
in high energy limit

$$\Delta E \approx \alpha_s \mu^2 L^2 / \lambda (1 + O(1/N))$$

implemented in models in different ways:

high initial densities $dN_g/dy=1100$ (Vitev/Gyulassy)

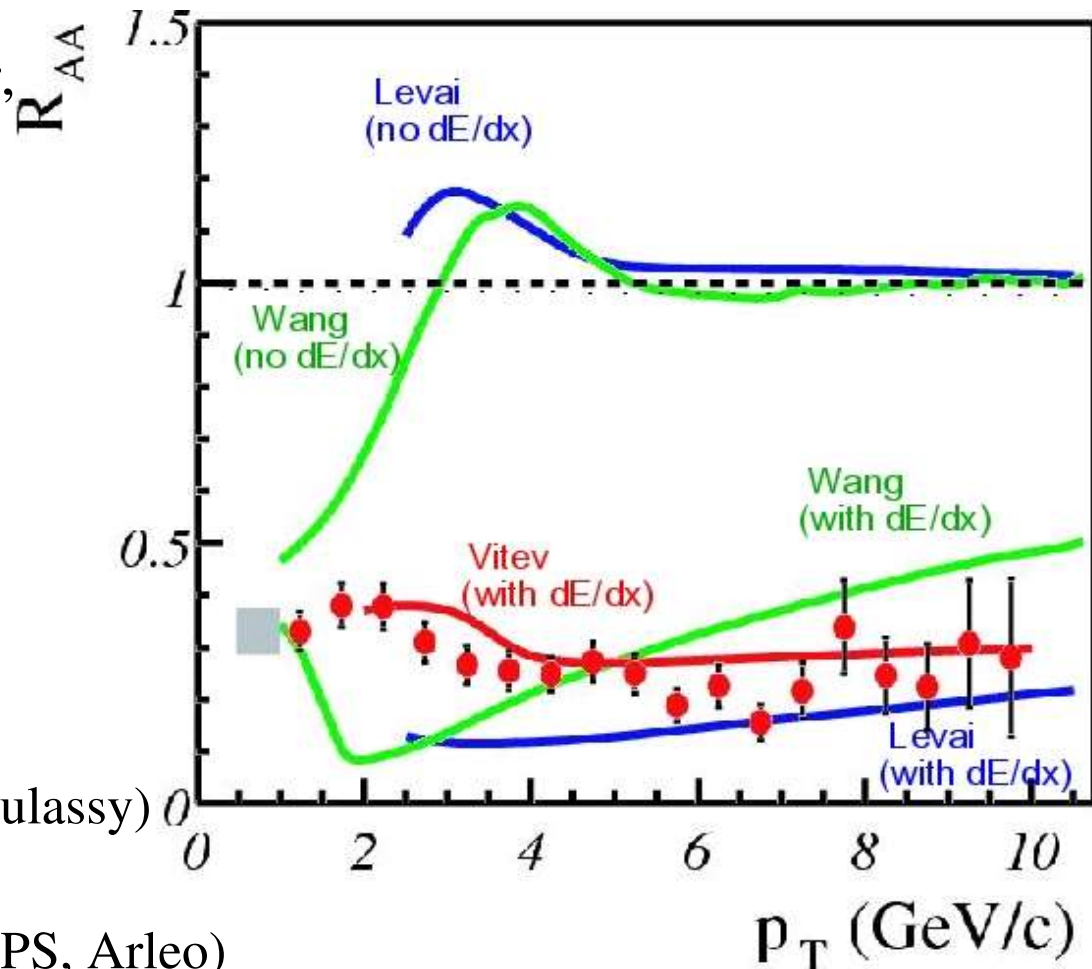
large opacities $\langle n \rangle = L/\lambda \approx 3-4$ (Levai et al.)

transport coefficients $q_0=3.5 \text{ GeV}/\text{fm}^2$ (BDMPS, Arleo)

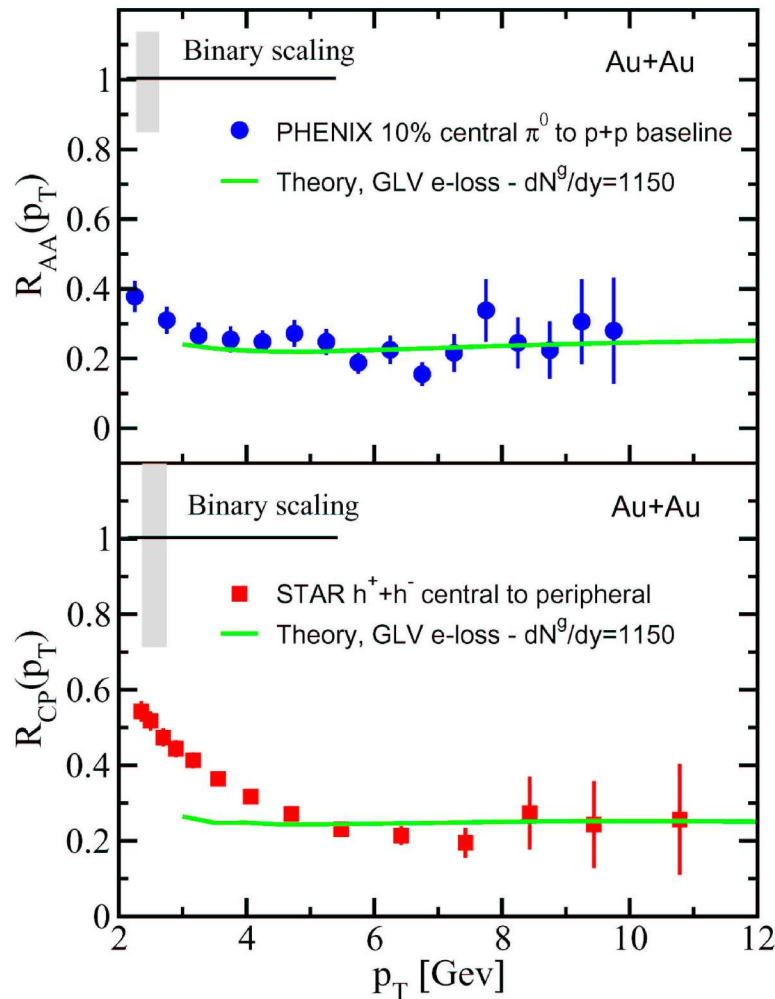
plasma temperature $T = 400 \text{ MeV}$ (G. Moore)

medium induced radiative energy loss

$dE/dx(\text{expanding})=0.25 \text{ GeV}/\text{fm}$ or $dE/dx(\text{static source})=14 \text{ GeV}/\text{fm}$ (S.N.Wang)



suppression of hadron yields at high p_t in central AuAu collisions

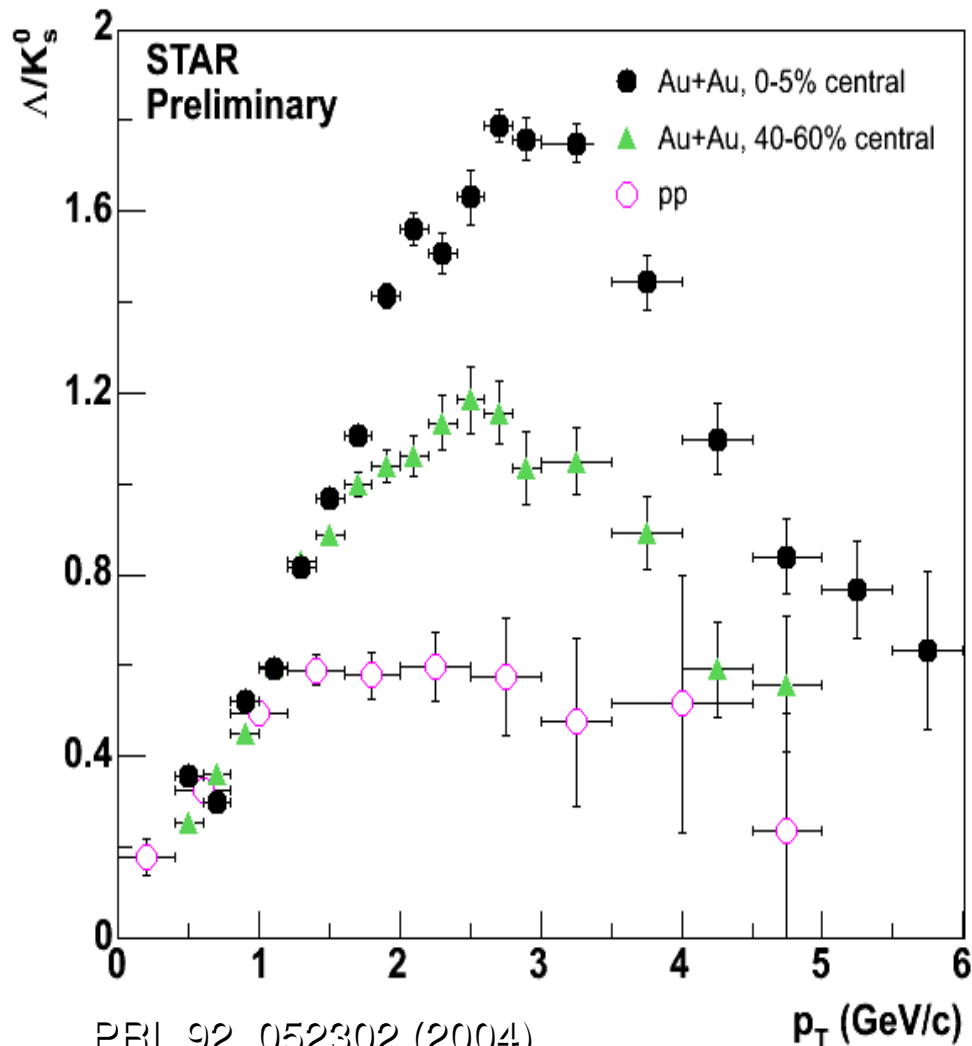


AuAu compared to pp scaled
with number of binary collisions

AuAu central collisions compared to
peripheral collisions scaled with
number of binary collisions

in central collisions hadron yields suppressed
indicative of jet quenching due to parton
energy loss due to high gluon density

suppression pattern depends on hadronic species



PRL 92, 052302 (2004)
J.Phys. G30, S963 (2004)

observed early on: protons exceed
number of pions at 2 GeV/c

general feature for all baryons?

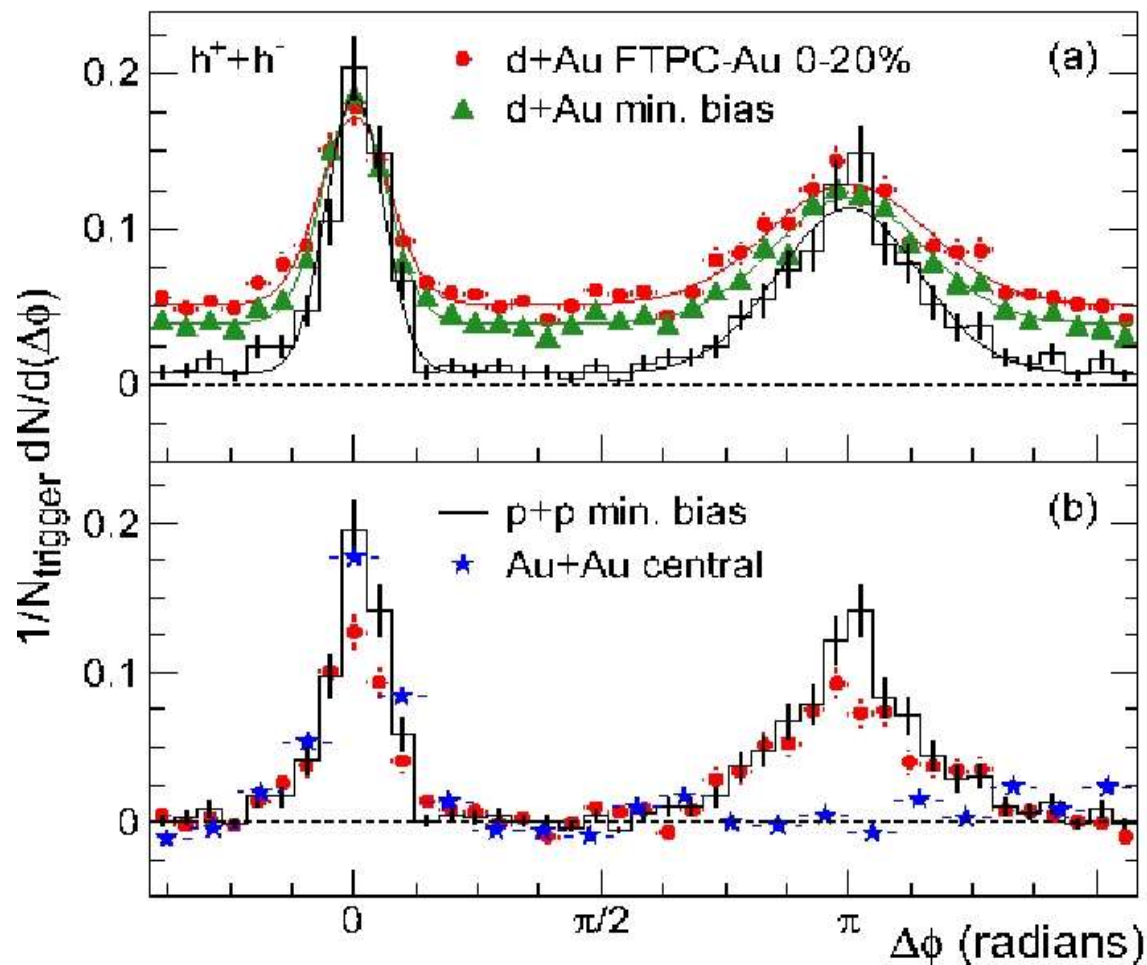
very cleanly shown recently for
 Λ as compared to K_s^0

while pp shows normal
fragmentation

coalescence of valence quarks? 3:2

Duke model...

Azimuthal correlations of high p_t particles



trigger particle: 4-6 GeV/c
correlated with all others
with $p_t=2-4$ GeV/c

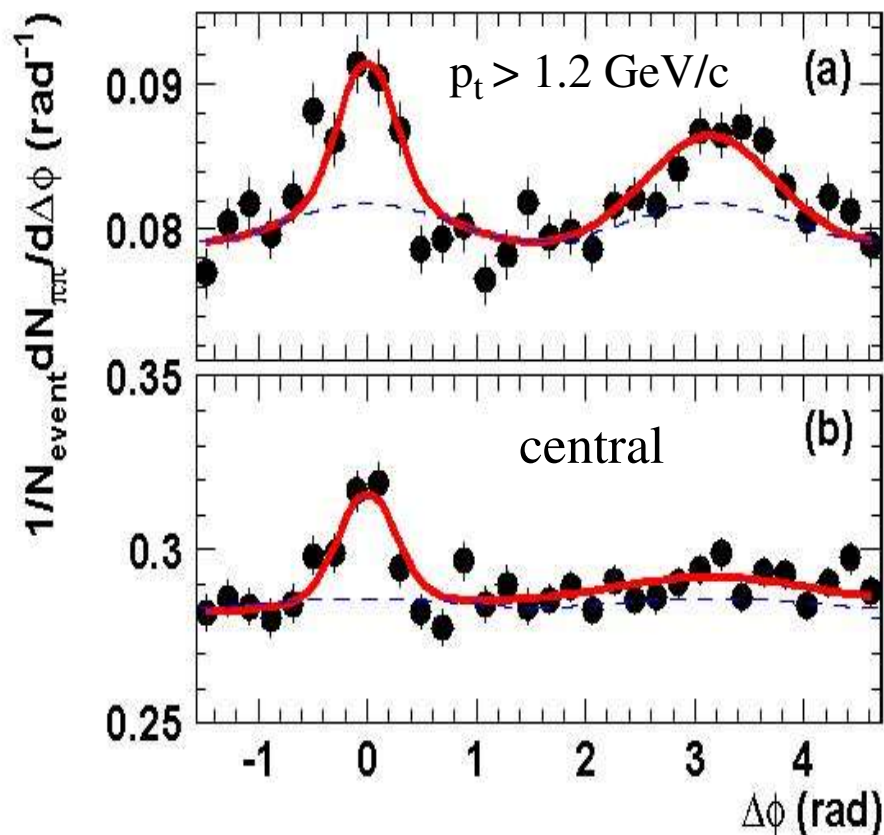
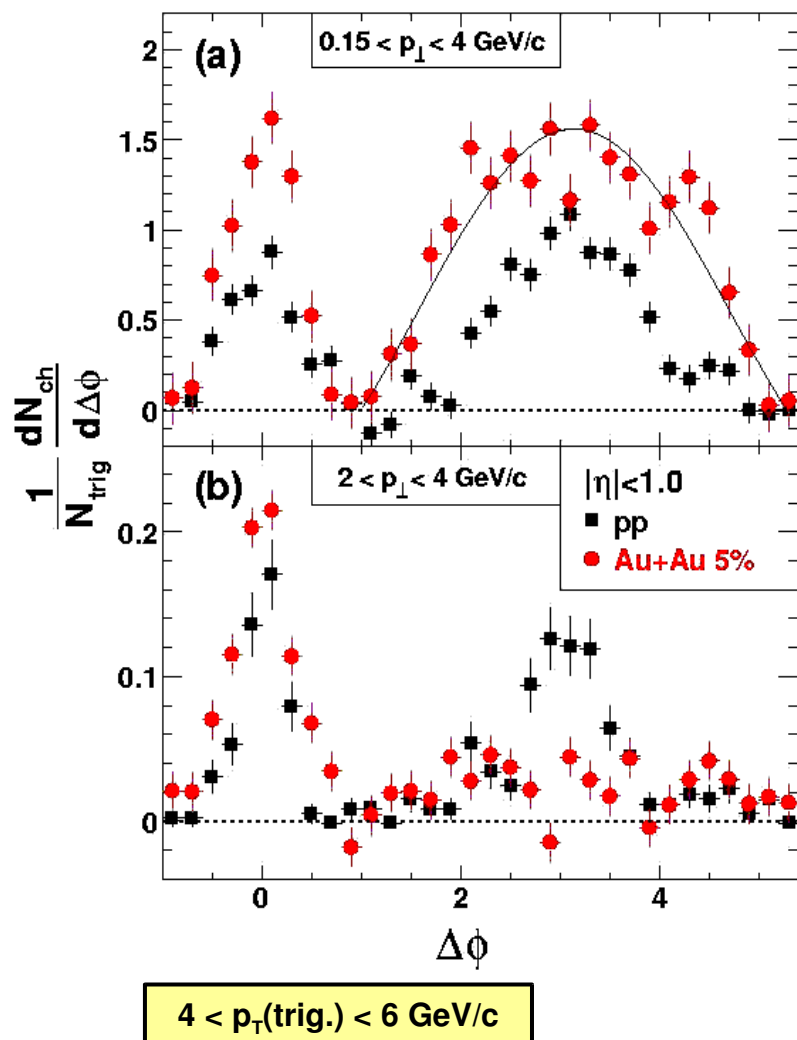
STAR: PRL 91 (2003) 072304

away-side associated hadrons at lower p_t

STAR 200 GeV nucl-ex/0501016

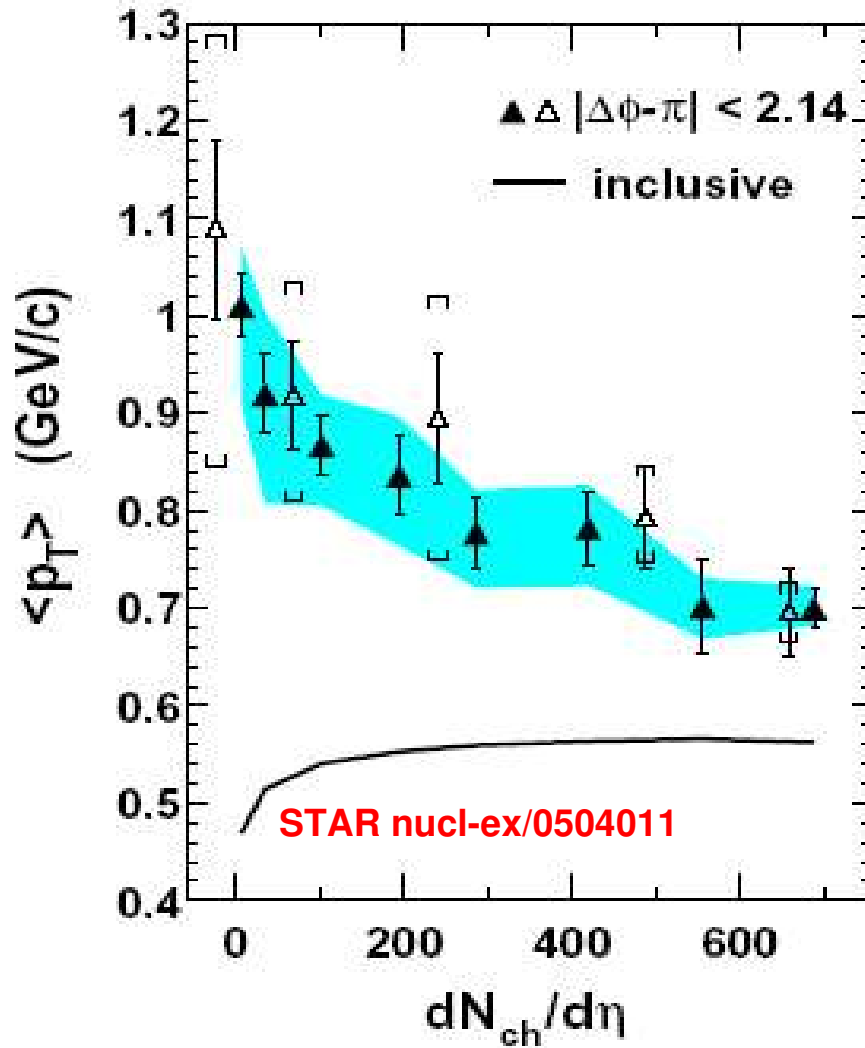
$\sqrt{s} = 17.2 \text{ GeV}$

CERES/NA45 PRL92(2004)032301



shape of away side peak
changes (broadens)
momenta reduced

mean p_t in cone opposite to leading trigger particle



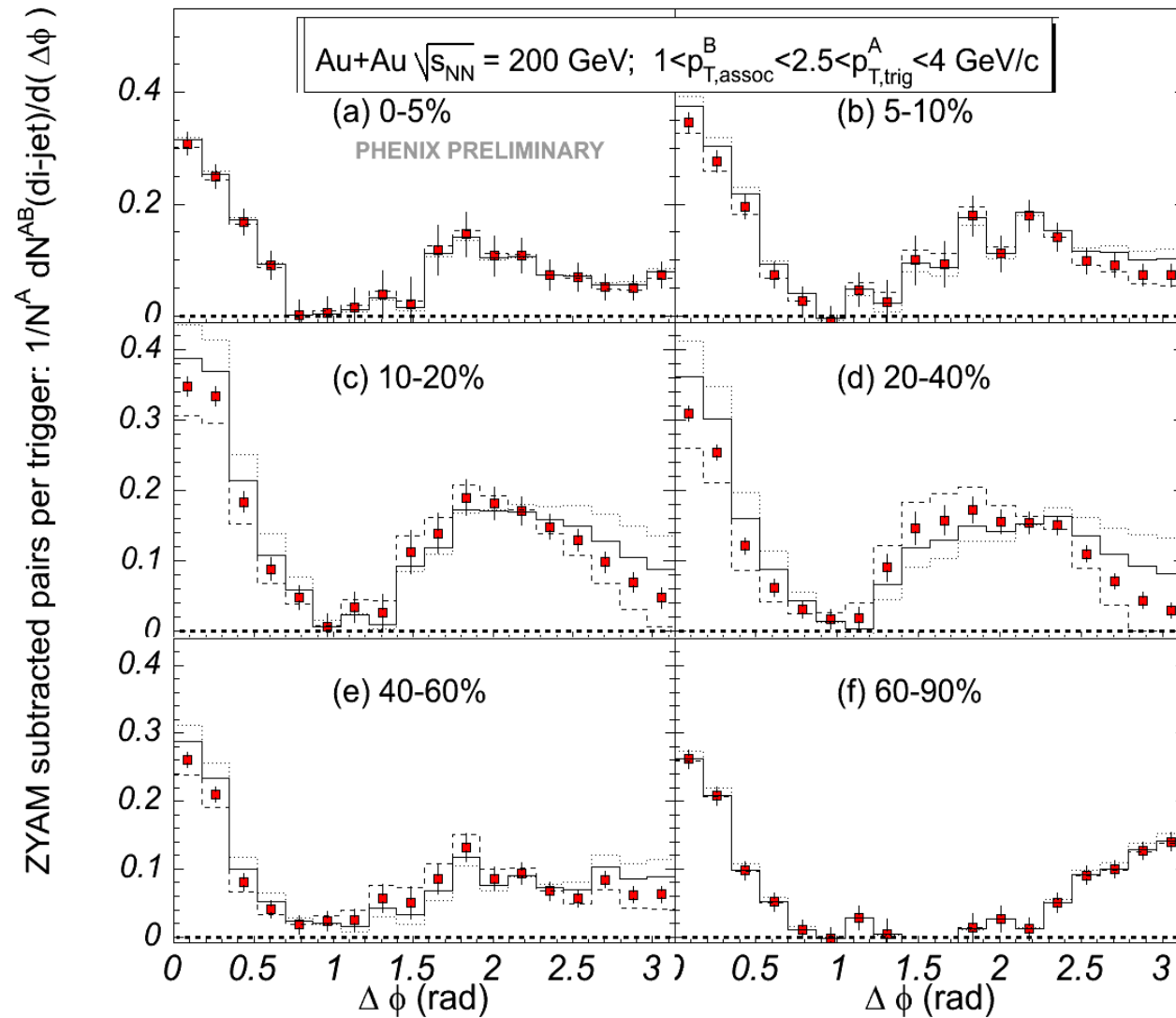
for central collisions
mean p_t on opposite side
looks nearly thermalized

$\sqrt{s_{NN}} = 200$ GeV
Au+Au results:

Closed symbols $\Leftrightarrow 4 < p_T^{trig} < 6$ GeV/c
Open symbols $\Leftrightarrow 6 < p_T^{trig} < 10$ GeV/c

Assoc. particles:
 $0.15 < p_T < 4$ GeV/c

opening angle correlations between high p_t particles



when asking for softer particle opposite hard trigger particle: dip (2σ) at $\Delta \phi = \pi$ except for most peripheral bin

N.N.Ajitanand, Proc. Int. Conf. Phys. & Astrophys. of QGP, Kolkata Feb. 2005

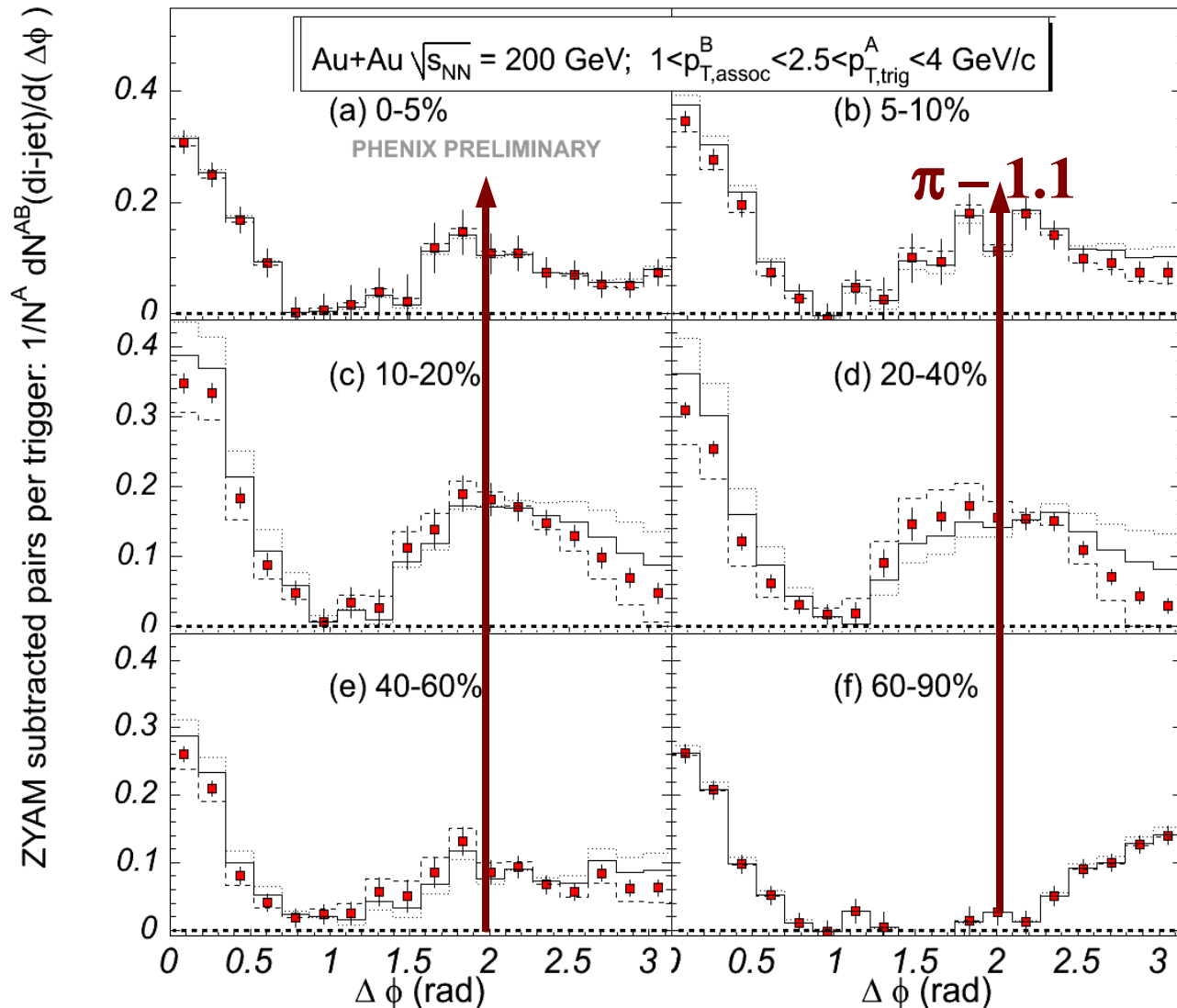
Mach cone due to sonic boom from quenched jets?

original idea:
Stöcker/Greiner 1976
for nuclear reactions;
Stöcker 2004:
60° cone for jets in QGP

J.Casalderrey-Solana, E. Shuryak, D. Teaney, hep-ph/0411315

if this can be established
far reaching consequences:
sensitivity to speed of
sound and EOS

experimental challenge:
can one see cone in 2d?
rel to reaction plane?



jet quenching indicative of gluon rapidity density

	$\tau_0 [fm]$	$T [MeV]$	$\epsilon [GeV / fm^3]$	$\tau_{tot} [fm]$	dN^g / dy
SPS	0.8	210-240	1.5-2.5	1.4-2	200-350
RHIC	0.6	380-400	14-20	6-7	800-1200
LHC	0.2	710-850	190-400	18-23	2000-3500

I. Vitev, JPG 30 (2004) S791

• Consistent estimate with hydrodynamic analysis

