

Experiments Lamb shift, hyperfine structure, quasimolecules and MO spectra

Lecture 3

April 29th, 2014



Preliminary plan of the lectures

- 1 15.04.2015 Preliminary Discussion / Introduction
- 2 22.04.2015 Experiments (discovery of the positron, formation of antihydrogen, ...)
- 3 29.04.2015 Experiments (Lamb shift, hyperfine structure, quasimolecules and MO spectra)
- 4 06.05.2015 Theory (from Schrödinger to Dirac equation, solutions with negative energy)
- 5 13.05.2015 Theory (photon, quantum field theory (just few words), Feynman diagrams, QED corrections)
- 6 20.05.2015 Theory (matrix elements and their evaluation, radiative decay and absorption)
- 7 27.05.2015 Experiment (photoionization, radiative recombination, ATI, HHG...)
- 8 03.06.2015 Theory (single and multiple scattering, energy loss mechanisms, channeling regime)
- 9 10.06.2015 Experiment (Kamiokande, cancer therapy, ...)
- 10 17.06.2015 Experiment (Auger decay, dielectronic recombination, double ionization)
- 11 24.06.2015 Theory (interelectronic interactions, extension of Dirac (and Schrödinger) theory for the description of many-electron systems, approximate methods)
- 12 01.07.2015 Theory (atomic-physics tests of the Standard Model, search for a new physics)
- 13 08.07.2015 Experiment (Atomic physics PNC experiments (Cs,...), heavy ion PV research)



Hydrogenic spectrum

- One needs three quantum numbers to define the state of a hydrogen (hydrogen-like) atom:

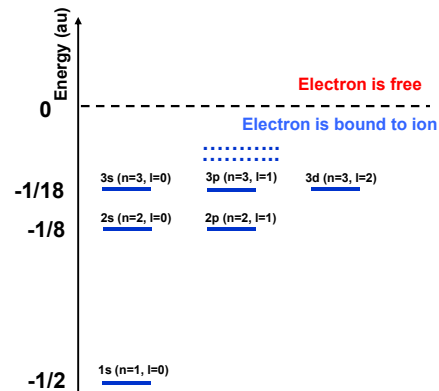
- $n = 1, 2, 3, \dots$ (principal)
- $l = 0, \dots, n-1$ (orbital)
- $m = -l, \dots, +l$ (magnetic)

- The energy depends only on the principal quantum number:

$$E_n = -\frac{Z^2}{2n^2}$$

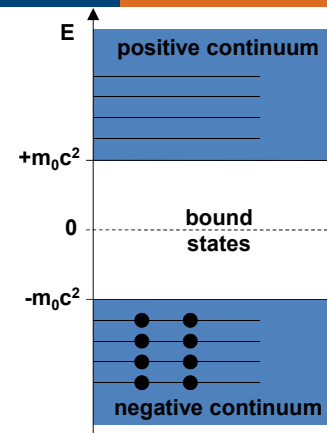
- i.e. in nonrelativistic theory the states are degenerate (l, m)!

$$\psi(\mathbf{r}) = \psi(r, \theta, \varphi) = R_{nl}(r)Y_{lm}(\theta, \varphi)$$



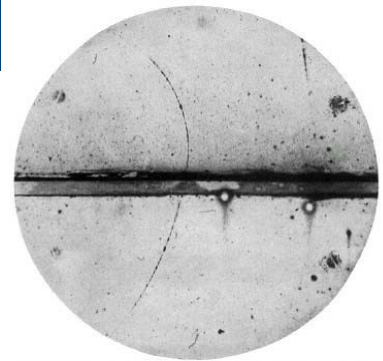
Why set (n, l, m) is good for the description of bound hydrogenic states?
Let us re-consider the set of (nonrelativistic) quantum numbers one more time!

Prediction of anti-matter



Dirac, Anderson, the Positron and the anti-matter. In his famous equation Paul Dirac combined (1929) the fundamental equation of quantum mechanics, the Schrödinger-equation with the theory of special relativity. He did not discard the negative energy –solutions of his equation as unphysical but interpreted them as states of the anti-particle of the electron (positron, having the same mass but opposite charge). In 1932 Carl Anderson discovered the positron the first time in the cosmic radiation. This was the proof of the existence of anti-matter, with incalculable consequences for the future of physics.





The first confirmation of a positron

Cloud chamber photograph by Andersen
 □ Phys. Rev. 43, 491 (1933)
 □ Nobel prize 1936 together with Victor Hess !

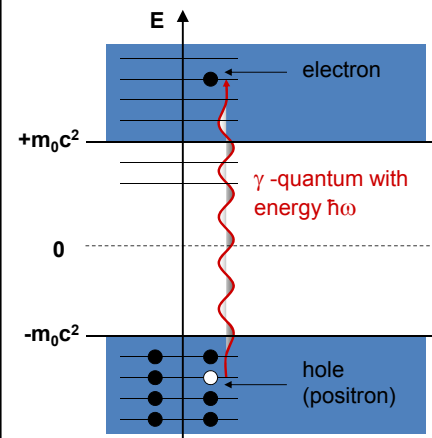
The first **'fingerprint' of anti-matter**. Anderson discovers the trace of a positron in his cloud chamber (in the middle one can see a lead-plate of 6mm thickness).

1. The **upper** part of the bending gives information about the **incoming-direction**.
2. The **lower** part gives the **positive charge** of the particle by its bending-direction.
3. By analyzing the **radius of curvature** before and after the transition the momentum can be estimated



Electron-positron pair production

In order to produce electron-positron pairs we would need: at least two times the rest mass of the electron !!!



+ energy of about $\hbar \omega \approx mc^2 \approx 1 \text{ MeV}$ to induce pair production

In case a whole in the negative continuum excites, an electron will immediately fill the vacancy and two 511 keV quanta are emitted



Experiments

Lamb shift in one and two-electron ions, hyperfine structure

Content

QED: Radiative correction viewed by an experimentalist

Experiments on Hydrogen

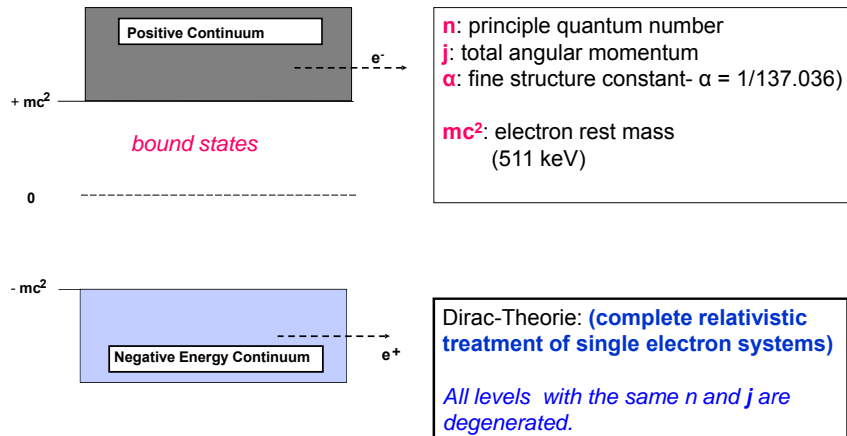
Strong fields: High-Z ions versus exotic atomic systems

The lamb shift in uranium

The hyperfine structure at high-Z

DIRAC Theory

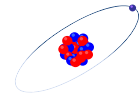
$$E_{nj} = mc^2 \sqrt{1 + \left(\frac{Z\alpha}{n - |j + 1/2| + \sqrt{(j + 1/2)^2 - (Z\alpha)^2}} \right)^2}$$



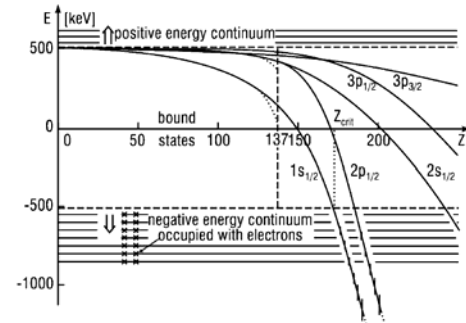
Critical electromagnetic fields

- ▶ Let us come back to Dirac energy of a *single* hydrogen-like ion:

$$E_{1s} = mc^2 \sqrt{1 - (Z\alpha)^2}$$



- ✦ What happens if we increase the nuclear charge Z ?
- ▶ If nuclear charge of the ion is greater than Z_{crit} the ionic levels can "dive" into Dirac's negative continuum.
- ▶ Physical vacuum becomes unstable: creation of pairs may take place!



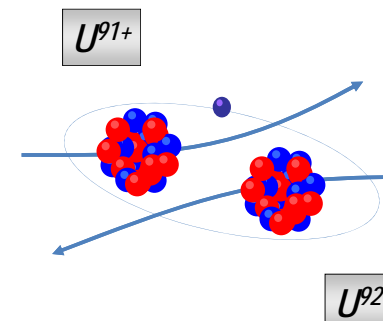
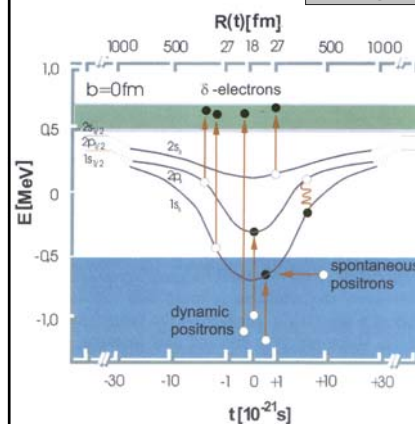
Relativistic quantum numbers and spectroscopic notations

| n | j | parity | spectroscopic notation |
|---|-----|--------|------------------------|
| 1 | 1/2 | + | $1s_{1/2}$ |
| 2 | 1/2 | + | $2s_{1/2}$ |
| 2 | 1/2 | - | $2p_{1/2}$ |
| 2 | 3/2 | - | $2p_{3/2}$ |

- Finally, we know how to characterize bound states of (relativistic) hydrogen.
- What are the energies of these states?

Supercritical fields: Formation of Quasi-Molecules

Merged Beams



A. Artemyev, et al. J. Phys. B 43 (2010) 235207

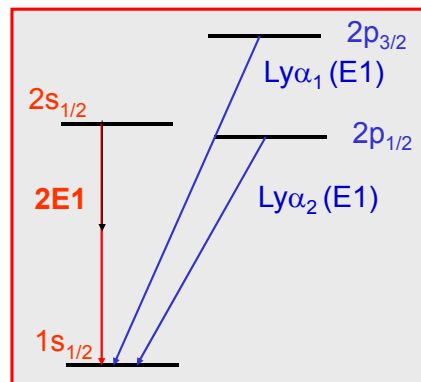


The Lamb shift

Experiment: Lamb-Retherford 1947; Theoretical Explanation: Bethe 1947

This experiment proves that for hydrogen the $2s_{1/2}$ and the $2p_{1/2}$ are not degenerated

=> In contradiction to Dirac-Theory



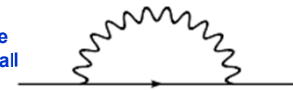
Level scheme for hydrogen

Lamb-Shift
deviations from the Dirac theory for a point like nucleus

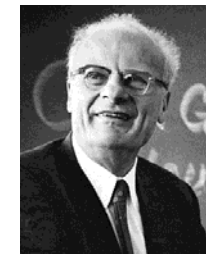
The Lamb shift (LS) is essentially a consequence of the interaction of the bound electron with its own radiation field

Lamb shift: idea of radiation corrections

- From the end of 1930th : interaction of electron with radiation field. But which field?
- Ideas: electron may interact with its own field. The Coulomb potential is therefore perturbed by a small amount and the degeneracy of the two energy levels is removed.
- But: problems with divergence of results!
- In 1947 Hans Bethe has shown how to identify the divergent terms and to subtract them from the theoretical expression.



Development of
Quantum Electrodynamics (QED)
and Quantum Field Theory (QFT)



Hans Bethe



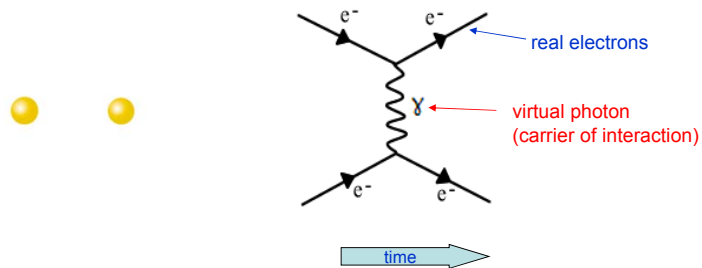
Feynman diagrams

- Feynman diagrams are a nice tool to represent exchange forces and, hence, to perform calculations of scattering processes.

Richard Feynman



Electron-electron repulsion



From: <http://www.slac.stanford.edu/>

The effects of QED increase with the intensity of the electrical field, the electrons are exposed to. That means the s electrons are most affected.

Lamb-Shift
(three effects)

QED

- self-energy

Self Energy

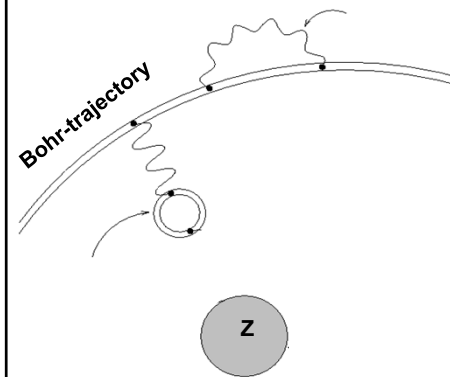


- Vacuum-polarization

Vacuum Polarization



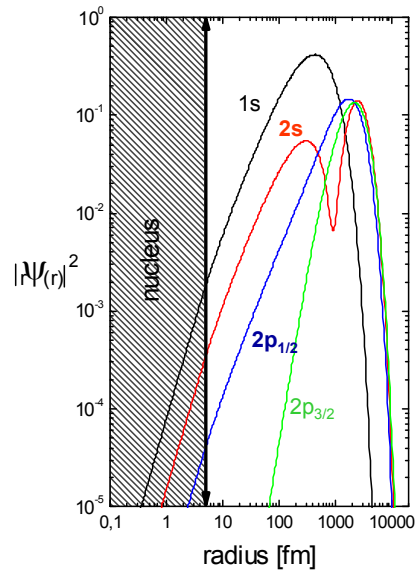
- size of the nucleus



Lamb Shift: An effect of strong fields

QED corrections
 $\Delta E \sim Z^4/n^3$

Z: nuclear charge
n: principal quantum number



Appendix: probability density of the electrons at the nucleus (non-relativistic)

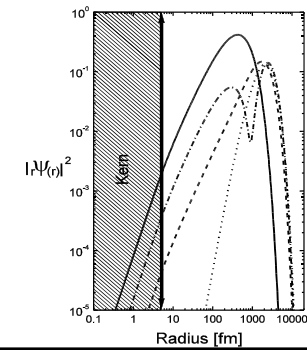
$r = Z/n$ in atomic units. With the normalization condition

$$\int_{-\infty}^{\infty} |\Psi(r)|^2 d^3r = 1$$

follows: $\Psi(r) = \left(\frac{Z}{n}\right)^{3/2}$

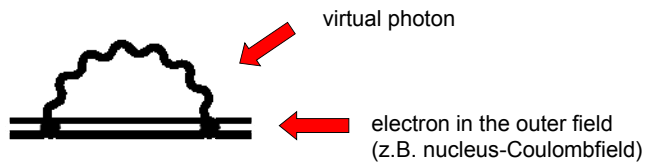
What is the electron density at the nucleus ($r=0$)?

for $l = 0$ $|\Psi(r=0)|^2 = \frac{Z^3}{n^3}$



Self energy

Self-energy: Describes the emission and reabsorption of a virtual photon (Bethe 1947) **The Self-energy decreases the binding energy.**



Mass-renormalization

$$m = m_e - \hat{\partial}m$$

m : measurable electron mass. It already includes the radiation correction.

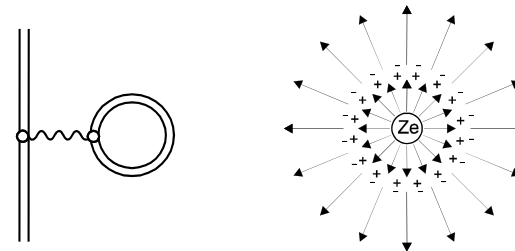
$$\hat{\partial}m$$

$$m_e$$

electron mass of the free electron without interaction with a field.

The self-energy is the most important radiation correction for light systems and decreases the binding energy. Here the electron emits and reabsorbs one and the same virtual photon. In every classical theory the consequence would be an infinitely high mass corresponding to an infinitely high self energy. The first calculations on the self-energy which explains the Lamb shift of Hydrogen were made by Bethe. His success is based on an idea of Kramer, called **mass-renormalization**. The interaction with the virtual photon leads to a modification of the electron mass. This effect is always present and the measured electron mass m already includes the radiation correction. Therefore the self-energy correction for a bound state can be regarded as the mass difference between the bound electron and a free electron of the mass m_e that is not liable to any interaction with the electromagnetic field at the time t

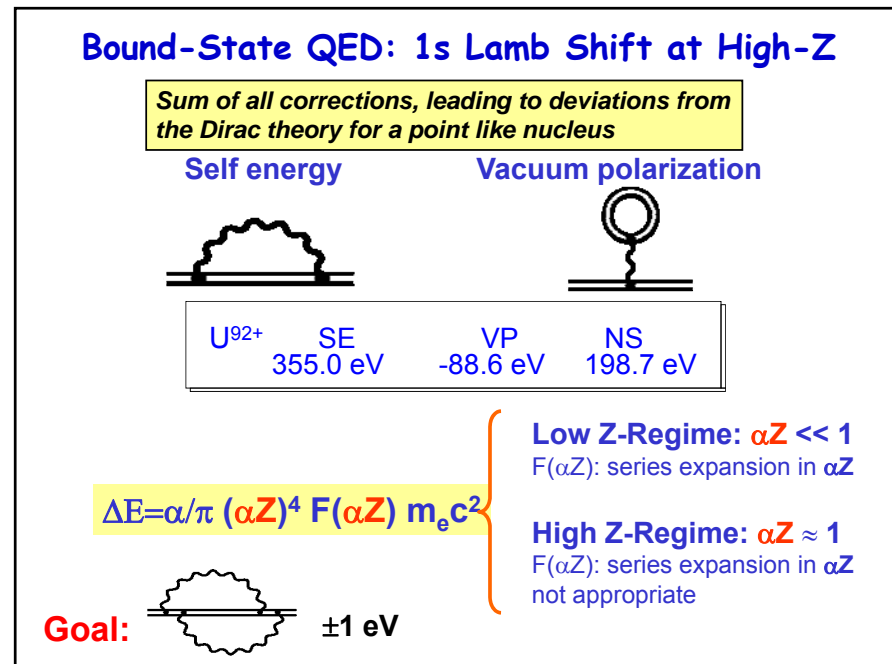
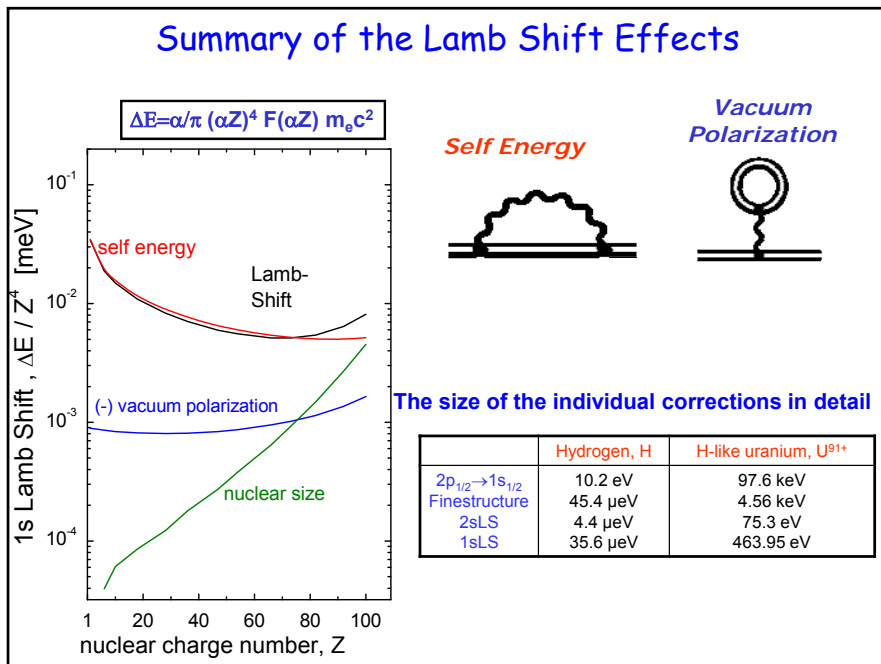
Vacuum Polarization (VP)



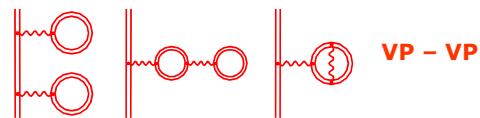
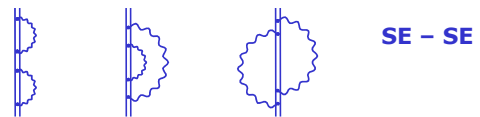
Virtual creation and annihilation of an e^+e^- -pair

Charge renormalization

VP acts attractive



QED correction in second order α



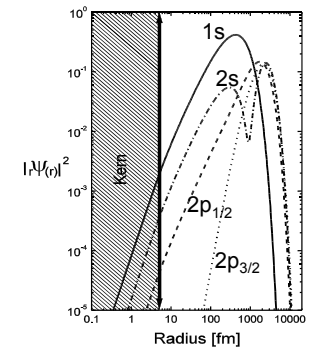
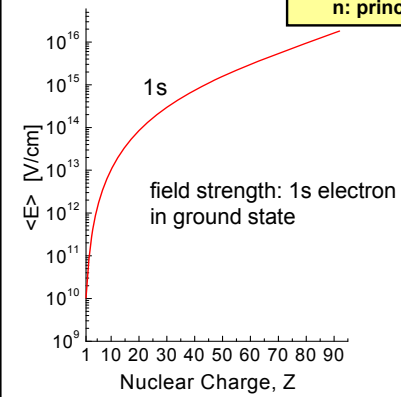
The Lamb-Shift

The effects of QED increase with the intensity of the electrical field, that the electrons are exposed to. That means the s electrons are most affected.

QED-corrections

$$\Delta E \sim Z^4/n^3$$

Z: nuclear charge
n: principal quantum number



Test of Bound-State QED at high-Z

Self Energy



Vacuum Polarization



Experiments

1s Lamb Shift

2eQED for He-like ions

2s-2p transitions in Li-like heavy ions

hyperfine-structure

g-factor of bound electrons

Content

QED: Radiative correction viewed by an experimentalist

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The hyperfine structure at high-Z

Hydrogen Atom

(<http://www.mpq.mpg.de/~haensch/hydrogen/h.html>)

$$f(1S-2S) = 2\,466\,061\,102\,474\,851(34) \text{ Hz}$$

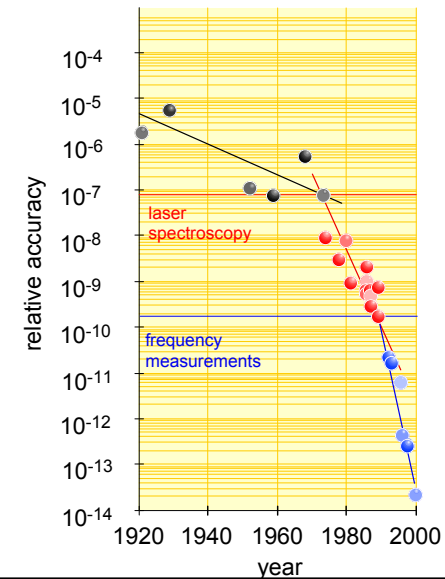
Rydberg Constant


$$R = 10\,973\,731.568\,525(84) \text{ m}^{-1}$$

Lamb Shift




$$L1S = 8\,172.840(22) \text{ MHz}$$

Laser Spectroscopy of Hydrogen



 **The Nobel Prize in Physics 2005**

"for his contribution to the quantum theory of optical coherence" "for their contributions to the development of laser-based precision spectroscopy, including the optical frequency comb technique"

| | | |
|---|--|--|
|  |  |  |
| Roy J. Glauber | John L. Hall | Theodor W. Hänsch |
| 🏆 1/2 of the prize USA | 🏆 1/4 of the prize USA | 🏆 1/4 of the prize Germany |
| Harvard University Cambridge, MA, USA | University of Colorado, JILA; National Institute of Standards and Technology Boulder, CO, USA | Max-Planck-Institut für Quantenoptik Garching, Germany; Ludwig- Maximilians- Universität Munich, Germany |
| b. 1925 | b. 1934 | b. 1941 |

<http://nobelprize.org/physics/laureates/2005/index.html>

Laser spectroscopy

2s

1s

f_0

Laser frequency: f

excitation frequency: f_0

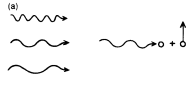
f

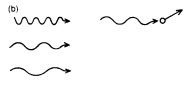
2s

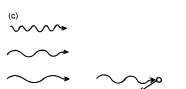
1s

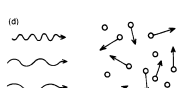
f_0

$v \ll c$

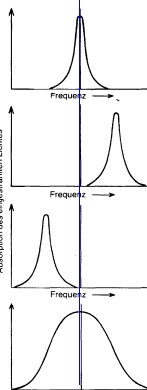
(a) 

(b) 

(c) 

(d) 

Absorption des eingestrahelten Lichtes



Frequenz

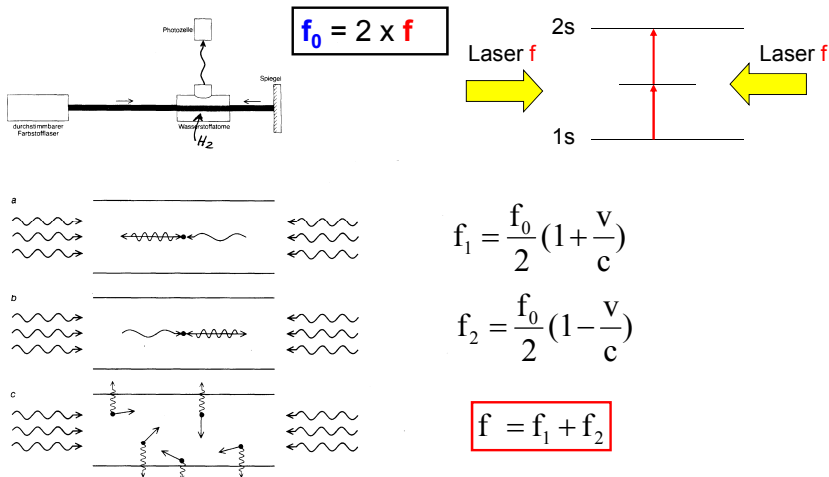
$f = f_0$

$f = f_0 \left(1 + \frac{v}{c}\right)$

$f = f_0 \left(1 - \frac{v}{c}\right)$

Large Doppler broadening by thermal distribution

Two-photon laser spectroscopy



Detection of fluorescence-radiation as a function of the laser frequency

Doppler free-Spectroscopy
Condition for Precision-spectroscopy

Real precision is achieved by relative measurements !!!

Assumption: The corrections (relativistic, QED etc.) are much better known for high n (much smaller !!!) than for n=1 and n=2

Idea: Compare Balmer- β radiation with Lyman- α

Balmer- β : $E_{Ba} = Ry \cdot \left(\frac{1}{2^2} - \frac{1}{4^2} \right)$

Lyman- α : $E_{Ly} = Ry \cdot \left(1 - \frac{1}{2^2} \right)$

$4 \cdot E_{Ba} = 1 \cdot E_{Ly}$
classically correct !!!

Effectively it is:

$4 \cdot E_{Ba} = 1 \cdot E_{Ly} + C$

E_{Ba} and E_{Ly} are measurement categories in the experiment

C is the parameter for the ground state, one has to determine for checking the theories

Ion absorbs photon and is slowed

Spontaneous emission in random direction

On average ion is slowed down (cooled)

$\vec{p}_i = m\vec{v}_i$

$\vec{p} = m\vec{v}_i + \hbar\vec{k}$

$\vec{p} = \hbar\vec{k}'$

http://nobelprize.org/nobel_prizes/physics/laureates/1997/illpres/

The Nobel Prize in Physics 1997

The Nobel Prize in Physics 1997

Steven Chu, Claude Cohen-Tannoudji and William D. Phillips

for their developments of methods to cool and trap atoms with laser light.

Laser light red-detuned from resonance...

... is absorbed by an atom moving counter to the laser beam...

... and through spontaneous decay is re-emitted in a random direction.

Time ↓

Applications: Bose-Einstein of atoms

http://nobelprize.org/nobel_prizes/physics/laureates/2001/index.html

The Nobel Prize in Physics 2001

"for the achievement of Bose-Einstein condensation in dilute gases of alkali atoms, and for early fundamental studies of the properties of the condensates"

| | | |
|--|--|--|
| <p>Eric A. Cornell</p> <p>1/3 of the prize</p> <p>USA</p> <p>University of Colorado, CO, USA</p> <p>b. 1961</p> | <p>Wolfgang Ketterle</p> <p>1/3 of the prize</p> <p>Federal Republic of Germany</p> <p>Massachusetts Institute of Technology (MIT), Cambridge, MA, USA</p> <p>b. 1957</p> | <p>Carl E. Wieman</p> <p>1/3 of the prize</p> <p>USA</p> <p>University of Colorado, CO, USA</p> <p>Booulder, CO, USA</p> <p>b. 1941</p> |
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Content

QED: Radiative correction viewed by an experimentalist

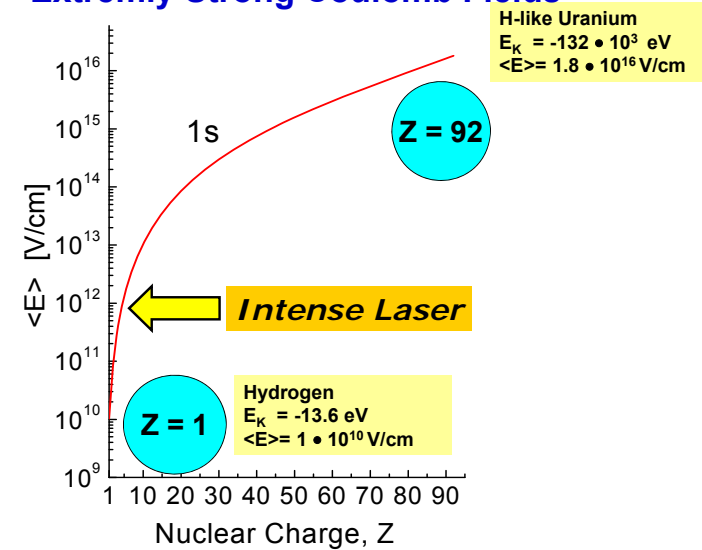
Experiments on Hydrogen

Strong fields: High-Z ions versus exotic atomic systems

The Lamb shift of uranium ($Z=92$)

The hyperfine structure at high-Z

Extremely Strong Coulomb Fields



1s-ground state: increase of the electric field strength by six orders of magnitude

Muonic atoms vs. HCI

Hightly charged ions

H-like ions

- $Z \gg 1$
- Relativistic corrections
- QED effects (self-energy)
- Electron-electron interaction
- High quantity production (ECRIS, EBIT, storage rings,...)

Muonic atoms

muonic hydrogen

- $m \gg m_{\text{electron}}$
- QED effects (vacuum polarization)
- Nuclear structure probe
- Diffcult to produce
- Obtained using intense pion beam (TRIUMF, PSI,...)

$$E_n \propto m \frac{Z^2}{n^2}$$

$$r_0 \propto \frac{1}{mZ}$$

1s Lamb shift in muonic atoms and HCI

Muonic hydrogen (-2.5 keV)

| | |
|---------------------|----------------|
| Lamb shift: | -1.9 eV |
| Self energy | 0.4% |
| Vacuum Polarization | 97.9% |
| Nuclear size | 1.6% |

Z

H-like silicon (-2.6 keV)

| | |
|---------------------|----------------|
| Lamb shift: | 0.48 eV |
| Self energy | 93.6% |
| Vacuum Polarization | 5.9% |
| Nuclear size | 0.5% |

Muonic nitrogen (-136 keV)

| | |
|---------------------|---------------|
| Lamb shift: | 199 eV |
| Self energy | 0.7% |
| Vacuum Polarization | 42.8% |
| Nuclear size | 56.5% |

H-like uranium (-132 keV)

| | |
|---------------------|---------------|
| Lamb shift: | 462 eV |
| Self energy | 54.3% |
| Vacuum Polarization | 13.8% |
| Nuclear size | 30.4% |

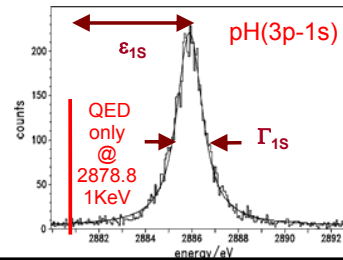
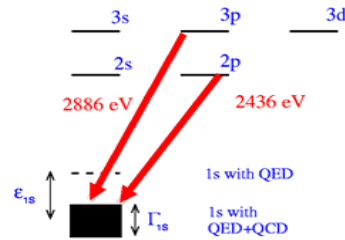
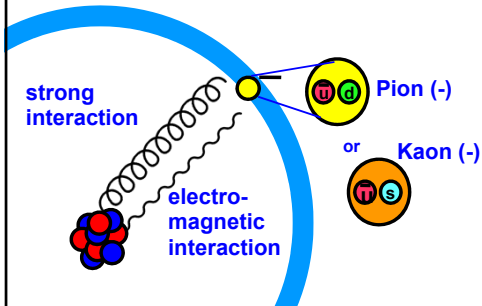
Complementary systems!!



More exotic bound systems...

Pionic and kaonic atoms

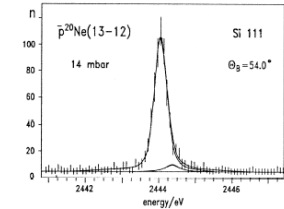
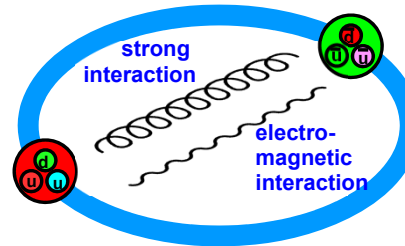
- H-like systems
- EM and **strong interaction**
- Nucleon-meson interaction ($Z=1,2,\dots$)
- Nuclear probe (high Z)



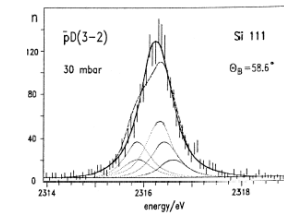
...and more!

Antiprotonic atoms

- EM and **strong interaction**
- Nucleon-antiproton interaction ($Z=1,2,\dots$)
→ nucleon-nucleon interaction
- Nuclear probe (high Z)
- Recoil and QED corrections
- Soon available in **high quantities at FAIR**



Obtained at LEAR (Gotta 1999)



Content

QED: Radiative correction viewed by an experimentalist

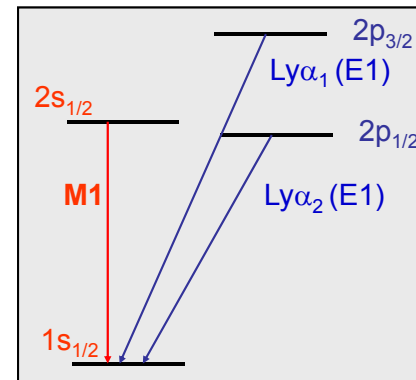
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The hyperfine structure at high-Z

The Structure of One-Electron Systems



QED corrections

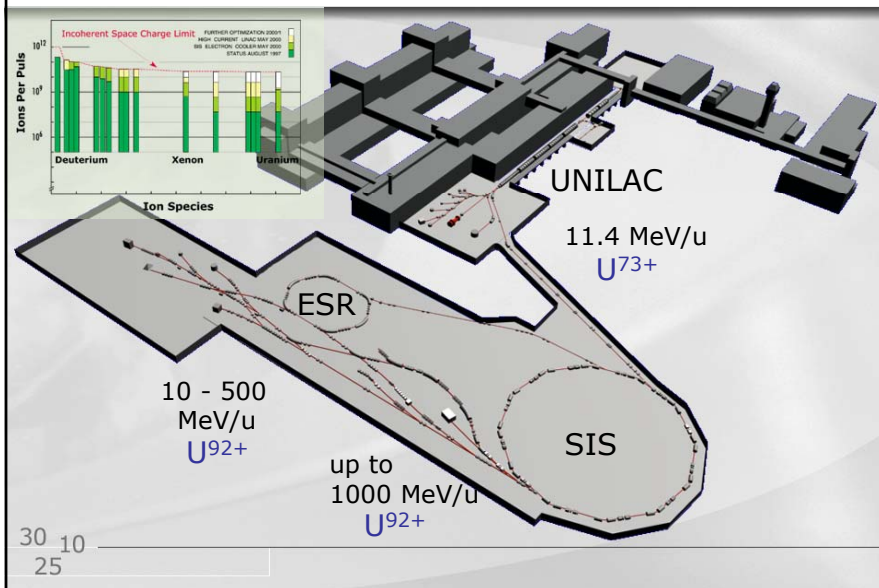
$$\Delta E \sim Z^4/n^3$$

Z: nuclear charge number
n: principal quantum number

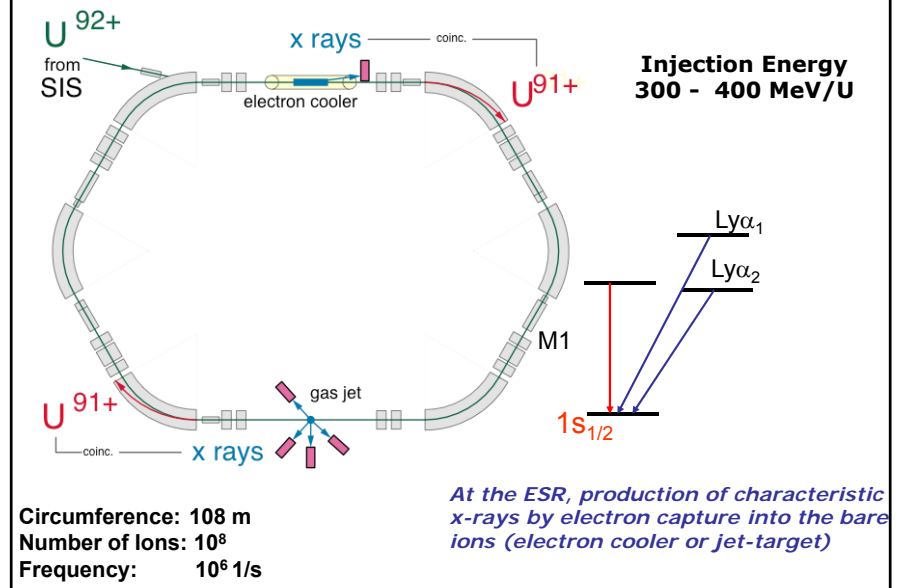
Atomic systems at high-Z

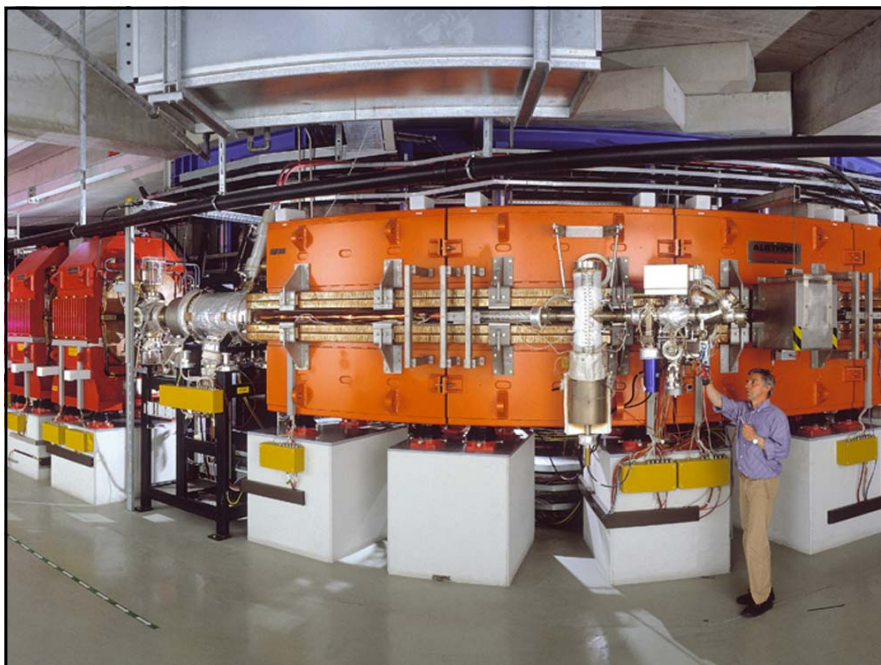
- Large relativistic effects on energy levels and transition rates (e.g. shell and subshell splitting)
- Large QED corrections
- Transition energies close to 100 keV

GSI-Accelerator Facility

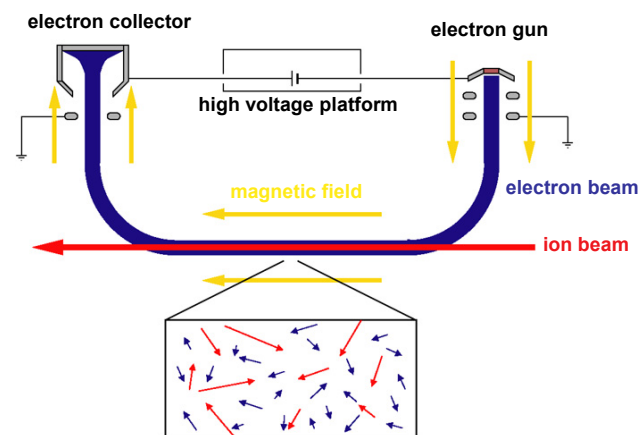


X-Ray Spectroscopy at the ESR Storage Ring

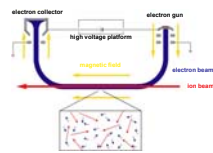




Storage Rings: Cooled Ion Beams



The Electron Cooler



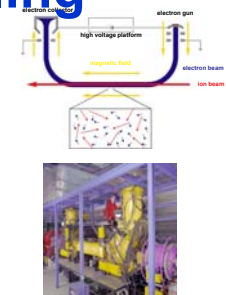
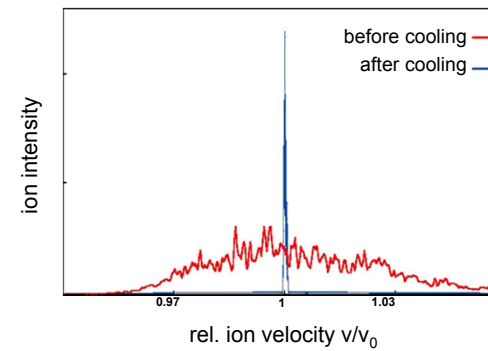
Electron Cooler

2.5 m interaction zone

Voltage: 5 to 200 kV

Current: 10 to 1000 mA

The Effect of Cooling



Ions interact 10^6 1/s with a collinear beam of cold electrons

Properties of the cold ions

Momentum spread $\Delta p/p : 10^{-4} - 10^{-5}$
 Diameter 2 mm

The Experimental Challenge

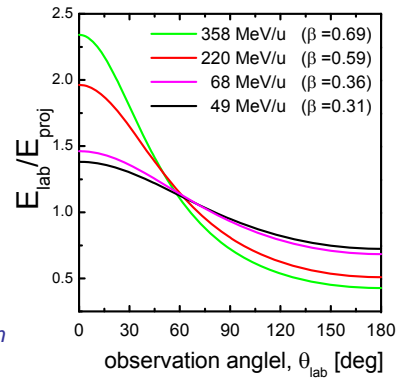
Relativistic Doppler-Transformation

$$E_{lab} = \frac{E_{proj}}{\gamma \cdot (1 - \beta \cdot \cos\theta_{lab})}$$

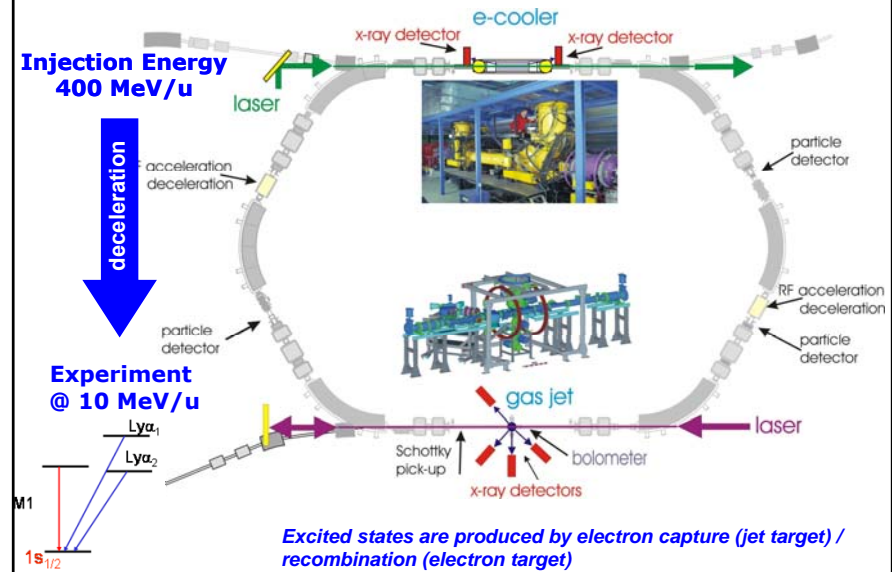
E_{lab} : Photon energy in the laboratory system
 E_{proj} : Photon energy in the emitter system

Doppler-Correction: Strong dependence on velocity and the observation angle θ_{LAB}

$$\gamma = \frac{1}{\sqrt{1 - \beta^2}}; \beta = \frac{v}{c}$$

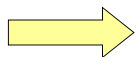


X-Ray Spectroscopy at the ESR Storage Ring



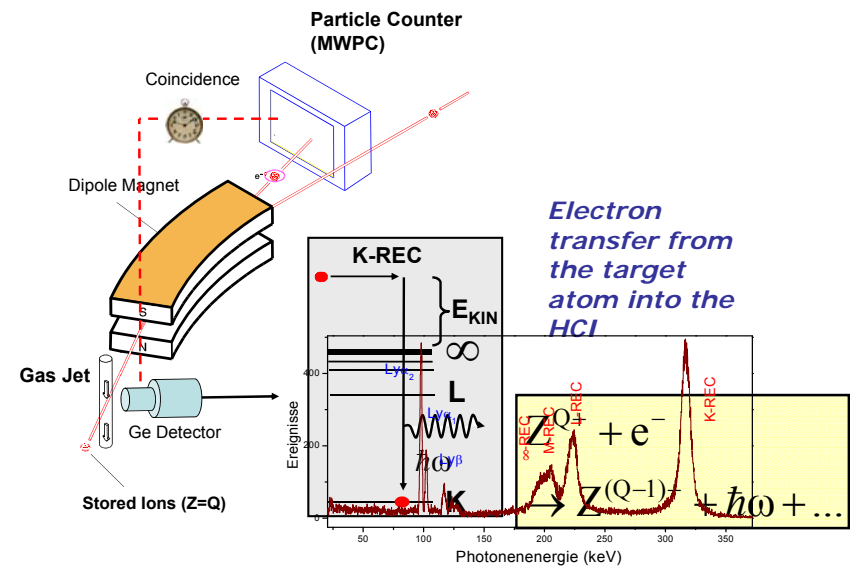
Principle of LS Experiments

1. Production of bare ions
2. Cooling
3. Electron capture in excited states (jet-target or electron cooler)
4. Detection of x-rays
5. Doppler correction
6. Comparison with Dirac theory

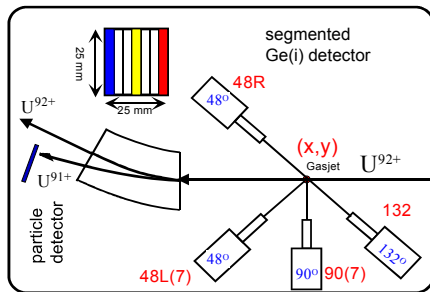


1s Lamb Shift

Experiments at the Jet-Target

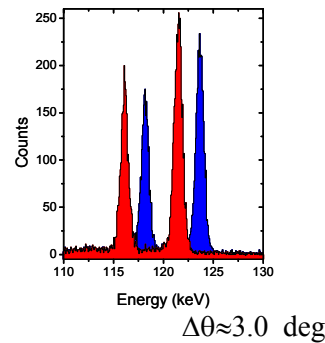


Lamb Shift-Experiment at the Jet-Target



- Simultaneous observation at various angles
- Forward/Backward symmetry
- Left/Right symmetry

$\text{Ly}\alpha_1 (2p_{3/2} \rightarrow 1s_{1/2})$
 $102\,171 \pm 13.2 \text{ eV}$

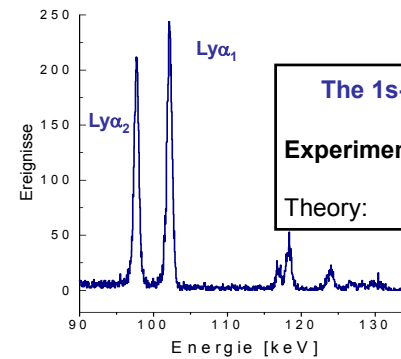


1s-Lamb Shift

Experiment: $468 \text{ eV} \pm 13 \text{ eV}$
 Theory: 466 eV^*

Result of the jet-target experiment using decelerated uranium ions

Geometry $\pm 8.5 \text{ eV}$ **$\Delta\beta$** $\pm 2.6 \text{ eV}$ **Fit** $\pm 9.7 \text{ eV}$ **$\text{Ly}\alpha_1 (2p_{3/2} \rightarrow 1s_{1/2})$** $102\,171 \pm 13.2 \text{ eV}$



The 1s-Lamb Shift

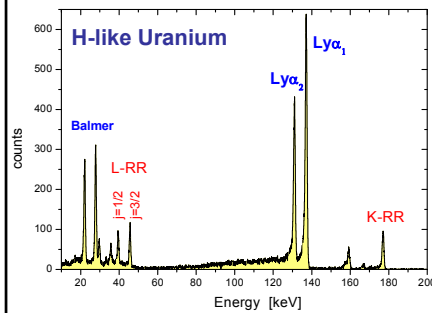
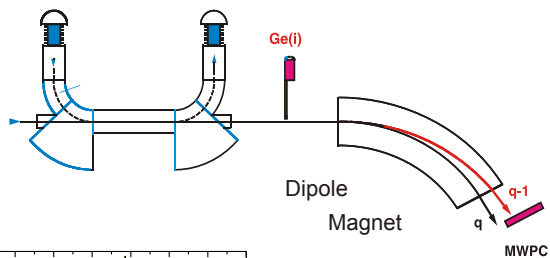
Experiment: $468 \text{ eV} \pm 13 \text{ eV}$

Theory: 466 eV^*

*Theory: T. Beier et al., 1998

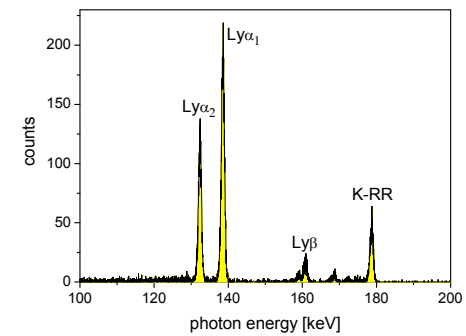
3% Sensitivity to Lamb Shift (466 eV)
4% Sensitivity to Self Energy (355 eV)
15% Sensitivity to Vacuum Polarization (-88.6 eV)

0° Spectroscopy at the Electron Cooler



- Blue shift has its maximum $\beta \approx 0.29 \Rightarrow E_{lab} \approx 1.43 \times E_{proj}$
- $\Delta\theta_{LAB}$ not critical, almost no Doppler width
- Uncertainty caused by $\Delta\beta$ has its maximum

The Ground State Lamb Shift in H-like Uranium



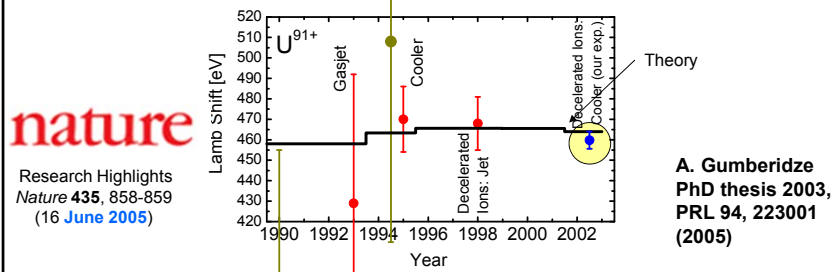
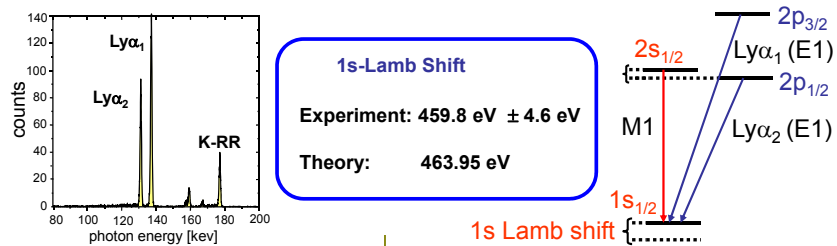
1s Lamb shift in U^{91+}

$$460.2 \pm 2.3 \pm 3.5 \text{ eV}$$

statistical
4.6 eV
uncertainty in the β

Test of Quantum Electrodynamics (1s-LS)

The 1s-LS in H-like Uranium



FOCAL
 H. Beyer, et al

$8 \cdot 10^8 \text{ Bq } ^{168}\text{Yb}$
 0.3 m
 curved crystal Si 220
 R = 2 m
 scanning or position-sensitive detector

$z = \frac{R}{2d}$

Scanning technique versus Bragg micro-strip detectors

$$\lambda = 2 \cdot d \cdot \sin \theta \approx \frac{Z}{R} \cdot 2d$$

Scanning
 FOCAL Spectrometer: $\epsilon \approx 10^{-8}$
 \Rightarrow 3 Events per Hour

Micro-Strip Germanium Detector
 Timing, Energy and Position Resolution

200 Strips
 $\Delta x \approx 200 \mu\text{m}$
 $\Delta E \approx 1.6 \text{ keV}$
 $\Delta t \approx 50 \text{ ns}$

EMMI
 Jilz-Institut Jena, 19.04.2013

There is a need for 2D detectors !

$$E_{\text{Lab}} = \frac{E_{\text{Proj}}}{\gamma(1 - \beta \cos \theta_{\text{Lab}})}$$

fast moving source stationary source

z-direction [mm]

Alexandre Gumberidze Wavelength standards in the x-ray region, Helmholtz-Institut Jena, 19.04.2013 59

Prototype 2D μ STRIP X-Ray Detector

2D μ STRIP planar detector systems for precision x-ray spectroscopy experiments (FOCAL)

energy resolution – timing – 2D position sensitivity

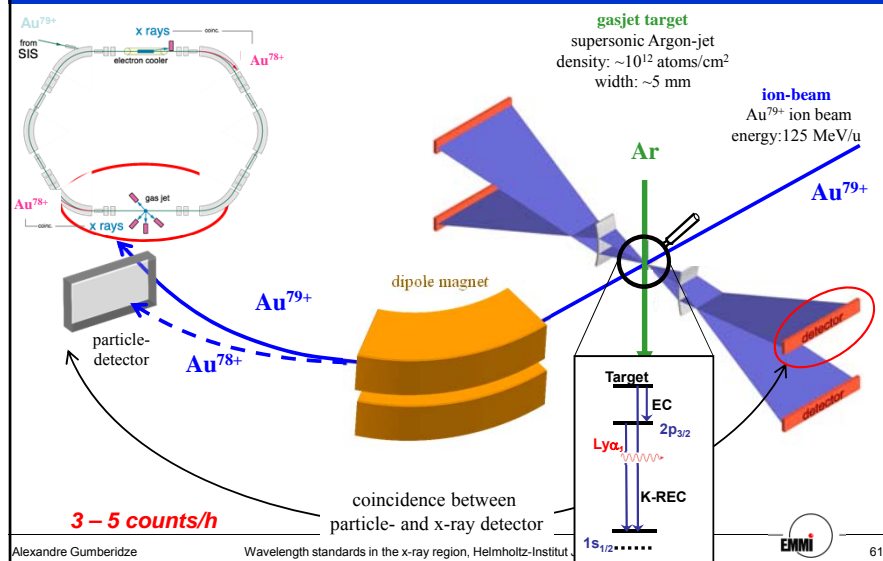
a-Ge-contact
(48 strips)

μ STRIP detector developed by

front: 128 strips pitch \sim 250 μ m
back: 48 strips pitch \sim 1167 μ m
equivalent to 6144 pixel

Alexandre Gumberidze Wavelength standards in the x-ray region, Helmholtz-Institut Jena, 19.04.2013

The FOCAL setup April 2012



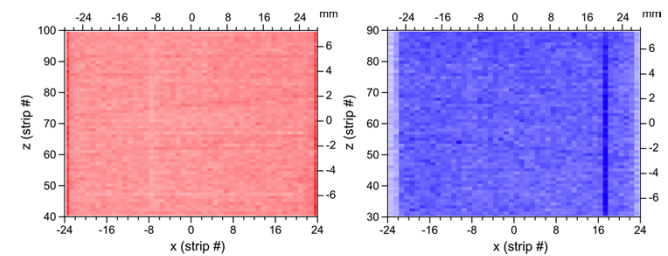
Alexandre Gumberidze

Wavelength standards in the x-ray region, Helmholtz-Institut

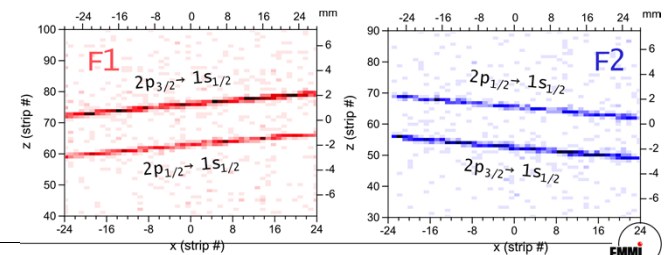
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Preliminary results

Raw 2D
spectrum



2D spectrum
with energy and
time condition



Alexandre Gumberidze

Wavelength standards in the x-ray region, Helmholtz-Institut Jena, 19.04.2013

EMMI

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Content

QED: Radiative correction viewed by an experimentalist

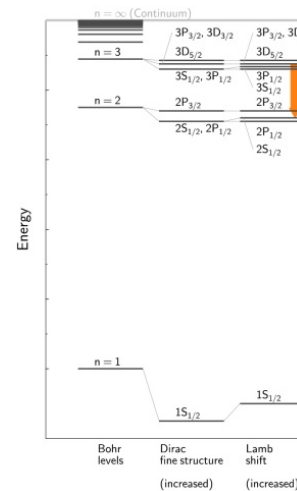
Experiments on Hydrogen

Strong fields: High-Z ions versus exotic atomic systems

The Lamb shift of uranium ($Z=92$)

The hyperfine structure at high-Z

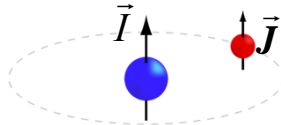
Energy levels of hydrogen atom: From Schrödinger to Dirac to QED



- Do we expect some more effects which can influence the level structure?
- Actually yes!

Hyperfine interaction

- Hyperfine structure is a perturbation in the energy levels of atoms/ions due to the magnetic dipole-dipole interaction, arising from the interaction of the nuclear magnetic moment with the magnetic field of the electron.



From: <http://www.wikipedia.org>

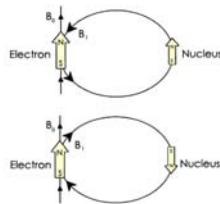
- Coupling of electron and nuclear momenta give rise to the total angular momentum of the system:

$$\vec{F} = \vec{J} + \vec{I}$$

- The HF splitting is:

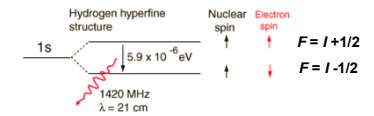
$$\Delta E_{HF} \approx a (F(F+1) - I(I+1) - J(J+1))$$

- Levels are split one more time!

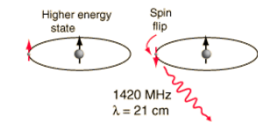


Hyperfine interaction

- For the case of ground $1s_{1/2}$ state ($j=1/2$) of hydrogen, HF interaction results in splitting of energy level into two levels.



- One can observe transition between two HF levels: famous 21 cm line in astrophysics!



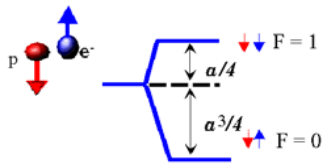
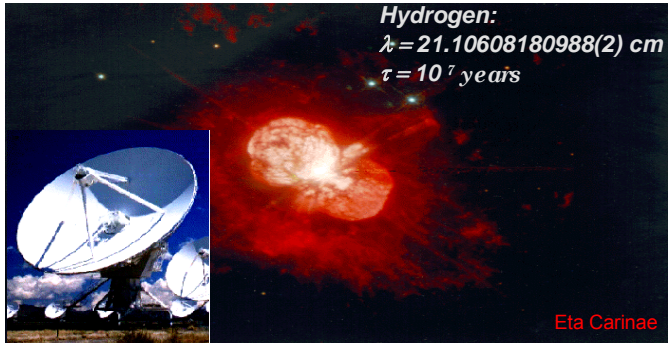
From: <http://hyperphysics.phy-astr.gsu.edu>



From <http://antwrp.gsfc.nasa.gov/apod/ap050825.html>

- 21 cm radiation is used, for example, to measure radial velocities of the spiral arms of the Milky Way.

Hyperfine Structure at High-Z



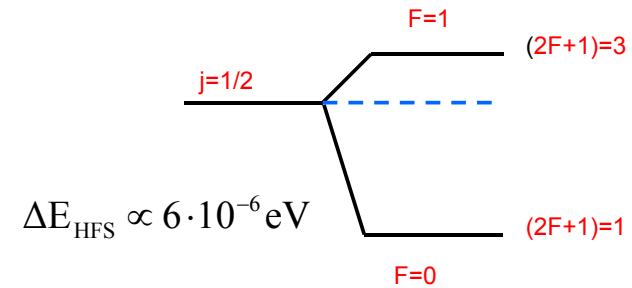
Hydrogen Atom: Groundstate

$$\Delta E_{HFS}^1 = +\frac{A}{4}$$

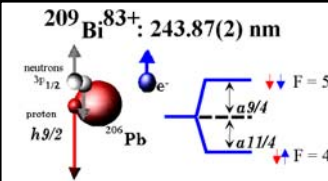
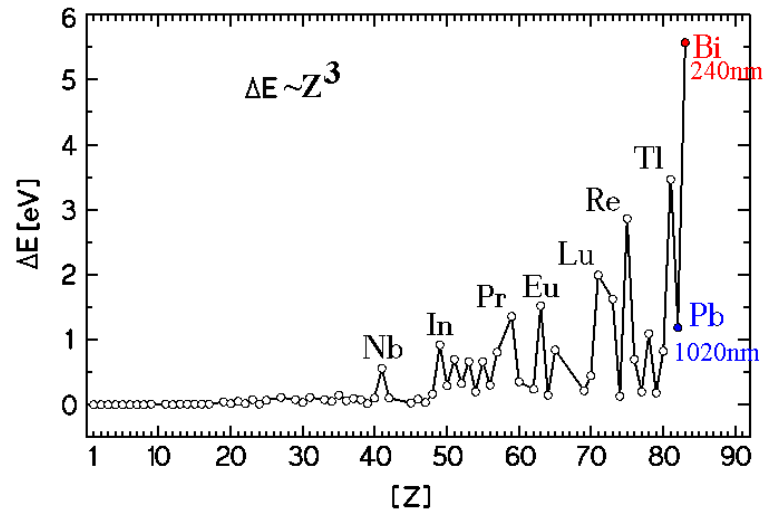
for $j=1/2, s_p=1/2, F=1$

$$\Delta E_{HFS}^0 = -\frac{3A}{4}$$

for $j=1/2, s_p=1/2, F=0$



Hyperfine Structure at High-Z



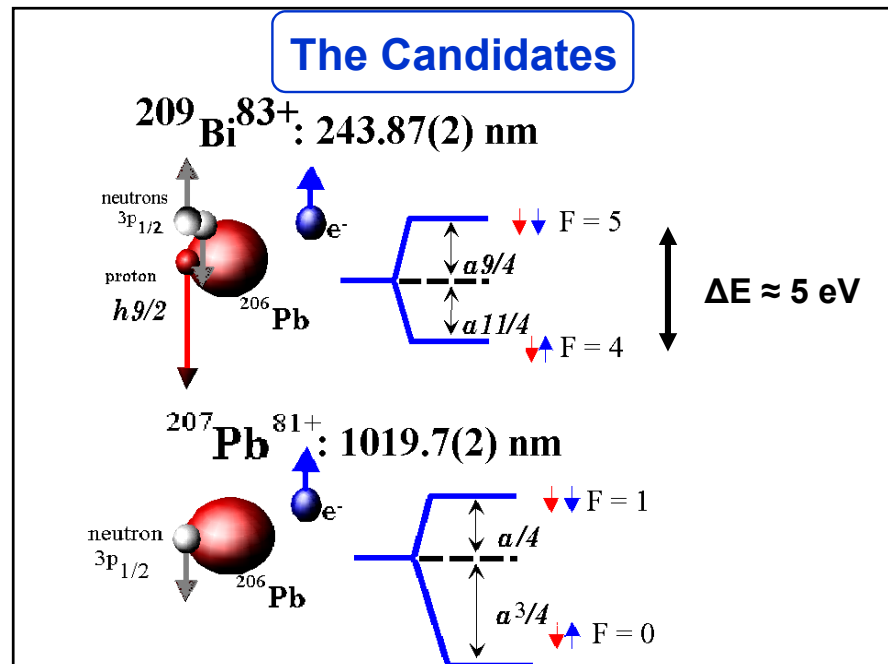
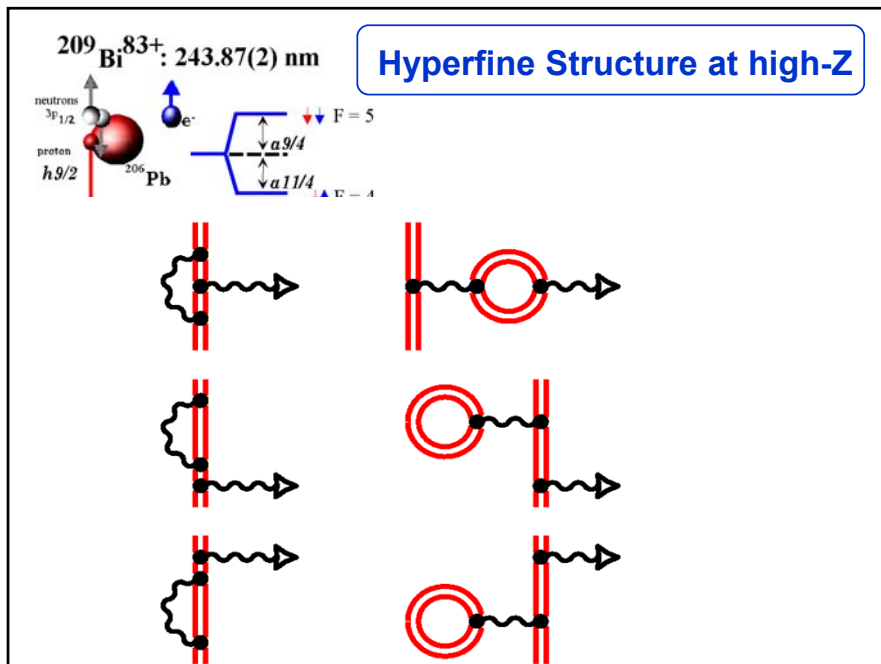
HFS: Test of QED in strong magnetic fields (high-Z)

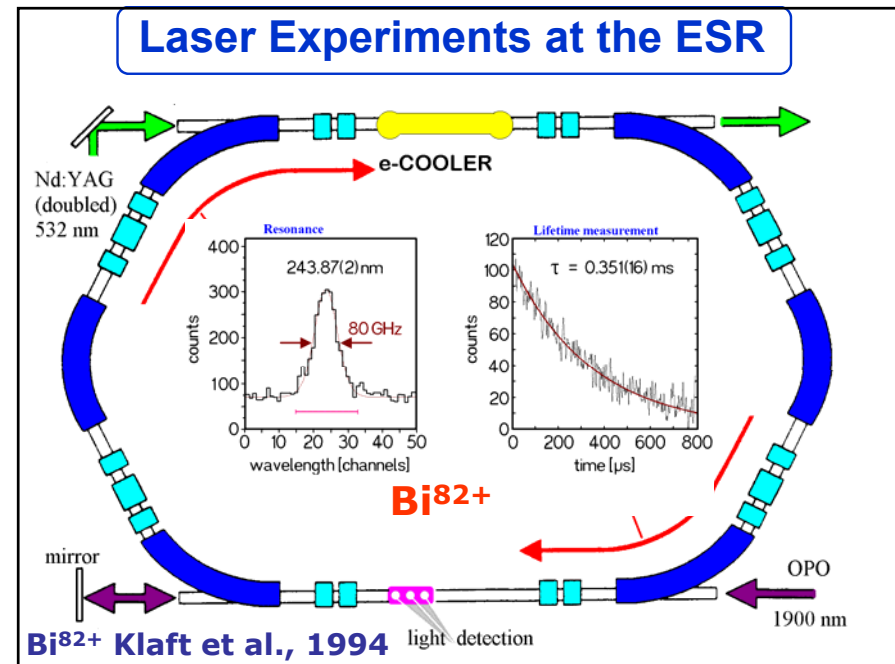
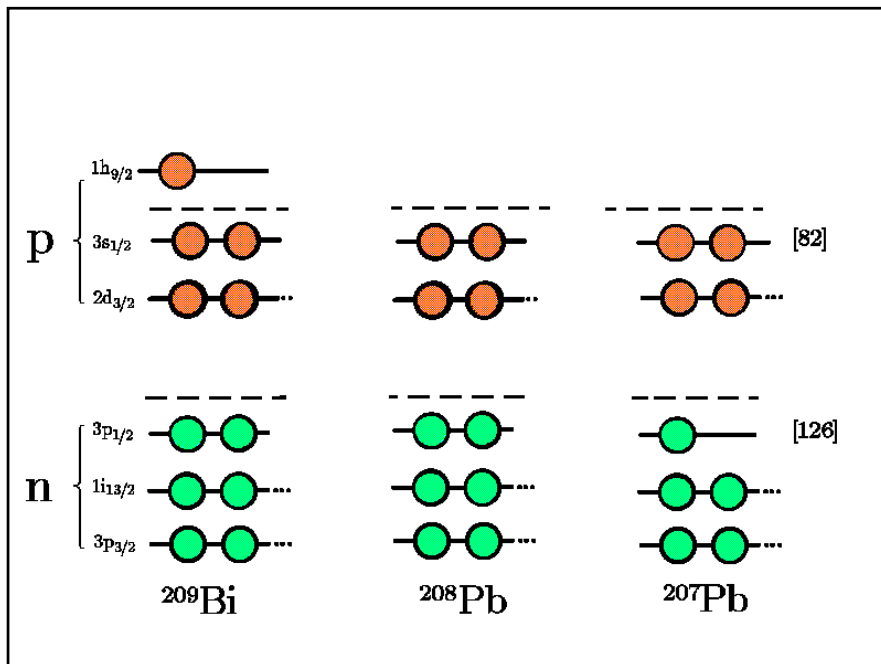
strong magnetic fields
 $\approx 10^{12}$ Gauss

Fermiformula for the HFS

$$\Delta \nu = \left(I + \frac{1}{2} \right) \cdot \frac{8}{3} \frac{h \cdot c}{n^3} \cdot Z^3 \cdot Ry \cdot \alpha^2 \cdot g_I \cdot \frac{M}{m_e} \times$$

- $F(j, Z)$ relativistic correction
- $(1 - \delta)$ Nuclear-Charge distribution (Breit-Rosental-Corwford-Effect)
- $(1 - \varepsilon)$ current distribution (Bohr-Weisskopf-Effect)
- QED ???

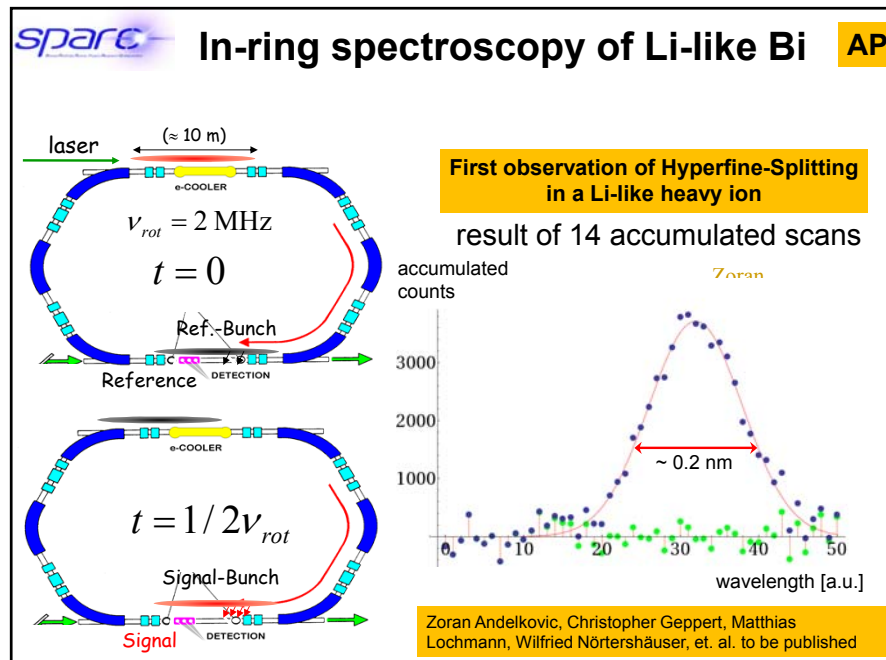




We have a problem !

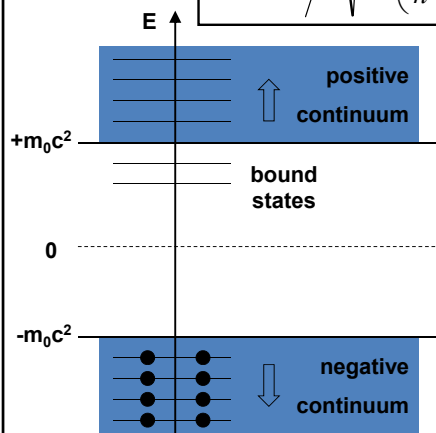
| | $^{209}\text{Bi}^{82+}$ | $^{207}\text{Pb}^{81+}$ |
|-------------------------------|-------------------------|-------------------------|
| <i>RMS radius</i> | 5.519 fm | 5.497 fm |
| <i>Magn.Mom.(corr.)</i> | 4.1106 n.m. | 0.58219 n.m. |
| Point nucl.(Dirac) | 212.320 nm | 880.017 nm |
| Breit-Schawlow | + 26.561(50) nm | +109.64(1) nm |
| <u><i>Ext. Nucleus</i></u> | <u>238.888(50) nm</u> | <u>989.65(1) nm</u> |
| Bohr-Weisskopf | + 5.025(330) nm | + 29.5(20) nm |
| <u><i>Theory, no QED</i></u> | <u>243.91(38) nm</u> | <u>1019.1(21) nm</u> |
| Vac. polarisation | - 1.64 nm | - 6.83 nm |
| Self energy | + 2.86 nm | + 11.9 nm |
| Total QED | + 1.22 nm | + 5.08 nm |
| <u><i>Theory incl.QED</i></u> | <u>245.13(58) nm</u> | <u>1024.2(25) nm</u> |
| Experimental | 243.87(2) nm | 1019.5(2) |

?



DIRAC Theory (relativistic formulation of quantum mechanics)

$$E_{nj} = mc^2 \sqrt{1 + \left(\frac{Z\alpha}{n - |j + 1/2| + \sqrt{(j + 1/2)^2 - (Z\alpha)^2}} \right)^2}$$



n: principal quantum number
j: total angular momentum
 α : Finestructure- or Sommerfeld-constant $\alpha = 1/137.036$
 mc^2 : electron mass at rest (511 keV)

Dirac-Theory: *Relativistic complete description of quantum mechanics*
All states with same n and j are degenerated

