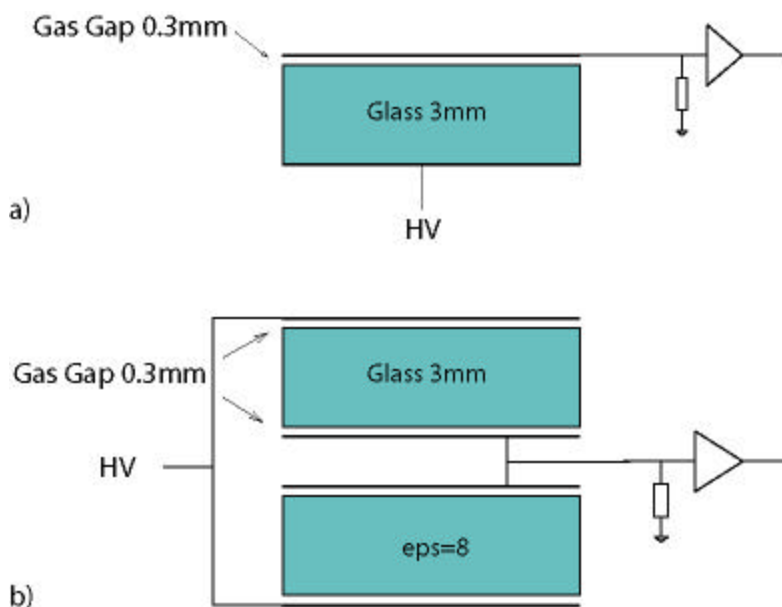




Detector Physics of Resistive Plate Chambers

Christian Lippmann (CERN), Werner Riegler (CERN)

Resistive plate chambers are gaseous parallel plate avalanche detectors with time resolutions down to 50 ps, making them attractive for trigger and time of flight applications.



- ◆ We simulate Timing RPCs in one and four gap configurations as in:

P. Fonte et. al., NIM A449 (2000) 295-301

A. Akindinov, P. Fonte et. al., CERN-EP 99-166

P. Fonte and V. Peskov, preprint LIP/00-04

- ◆ 0.3 mm gap(s); glass resistive plates ($\epsilon=8$, $r=2 \times 10^{12} \text{ Wcm}$)
- ◆ $\text{C}_2\text{F}_4\text{H}_2$ / $i\text{-C}_4\text{H}_{10}$ / SF_6 (85/5/10)
- ◆ HV: 3(6)kV, E: 100kV/cm



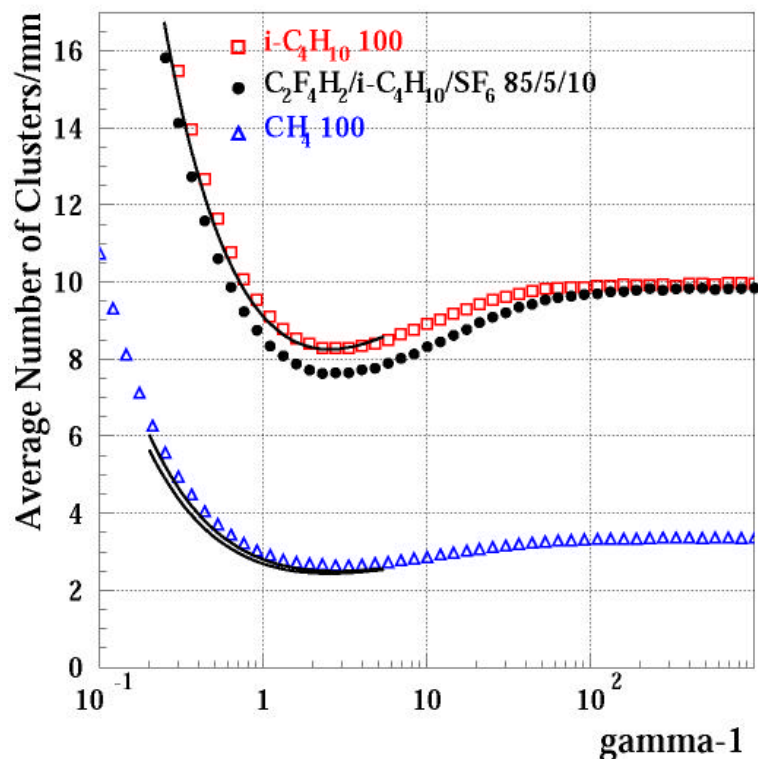
Simulation Input

- ◆ **Primary ionization:** **HEED** (Igor Smirnov)
- ◆ **Townsend, attachment coefficient:** **IMONTE** (Steve Biagi)
- ◆ **Diffusion, drift velocity:** **MAGBOLTZ 2** (Steve Biagi)
- ◆ **Avalanche fluctuations:** **Werner Legler** (Die Statistik der Elektronenlawinen in elektronegativen Gasen bei hohen Feldstaerken und bei grosser Gasverstaerkung,1960)
- ◆ **Space Charge Field:** **CERN-OPEN-2001-074**, "Static electric fields in an infinite plane condensor with one or three homogeneous layers"

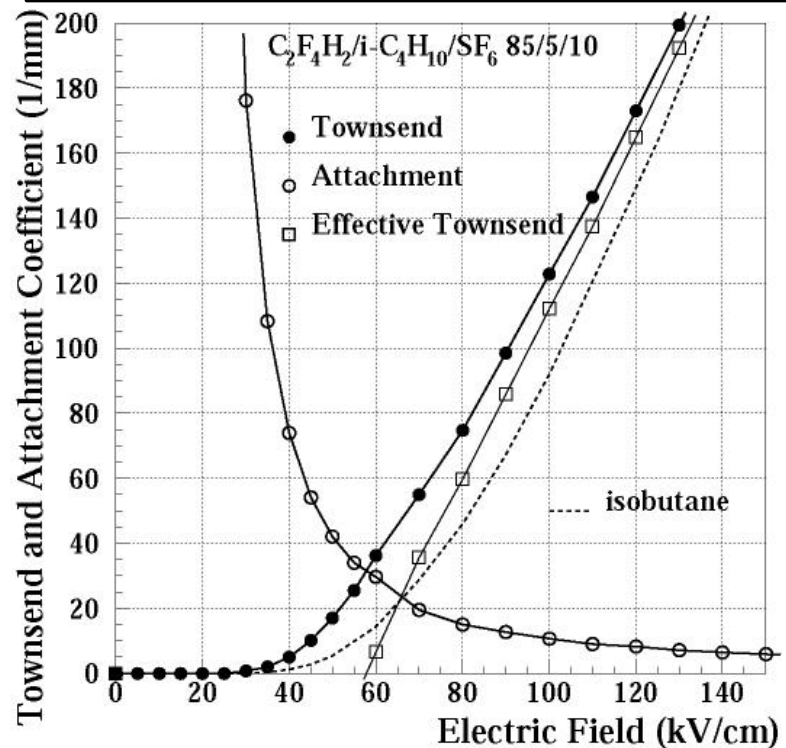


Gas Parameters

- ◆ Primary Ionization:
Number of Clusters per mm
- ◆ $p_0=1013\text{mbar}$, $T_0=296.15\text{K}$
- ◆ Measurements (lines): F.Rieke et al., Phys.Rev. A6, 1507 (1972)



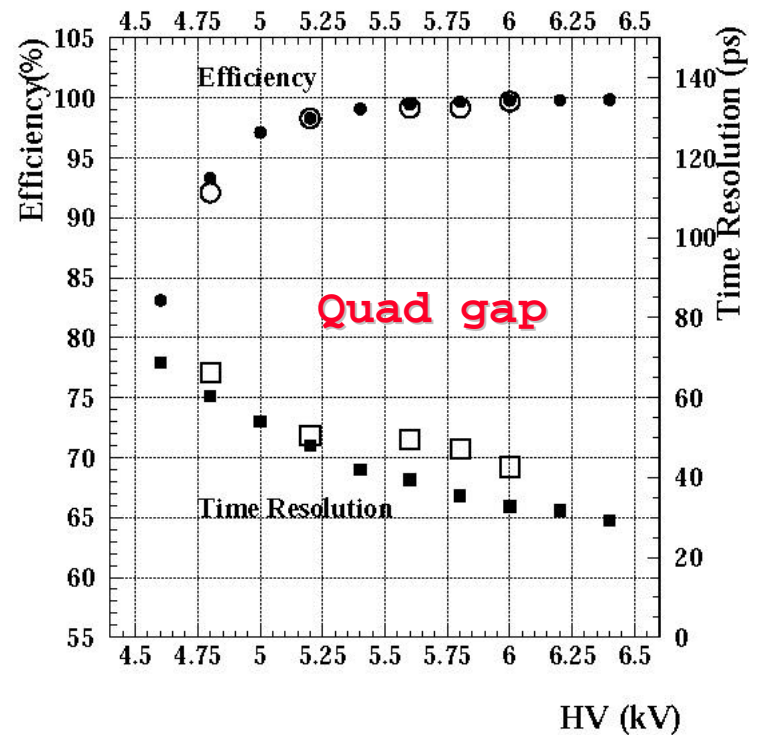
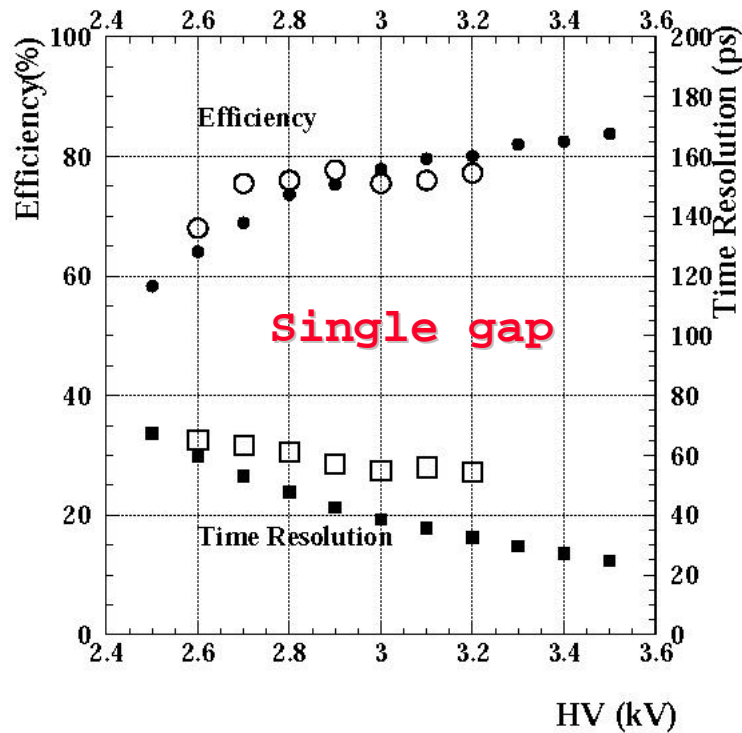
- ◆ Townsend Coefficient
- ◆ Attachment Coefficient
- ◆ $p_0=1013\text{mbar}$, $T_0=296.15\text{K}$





Efficiency and Time Resolution

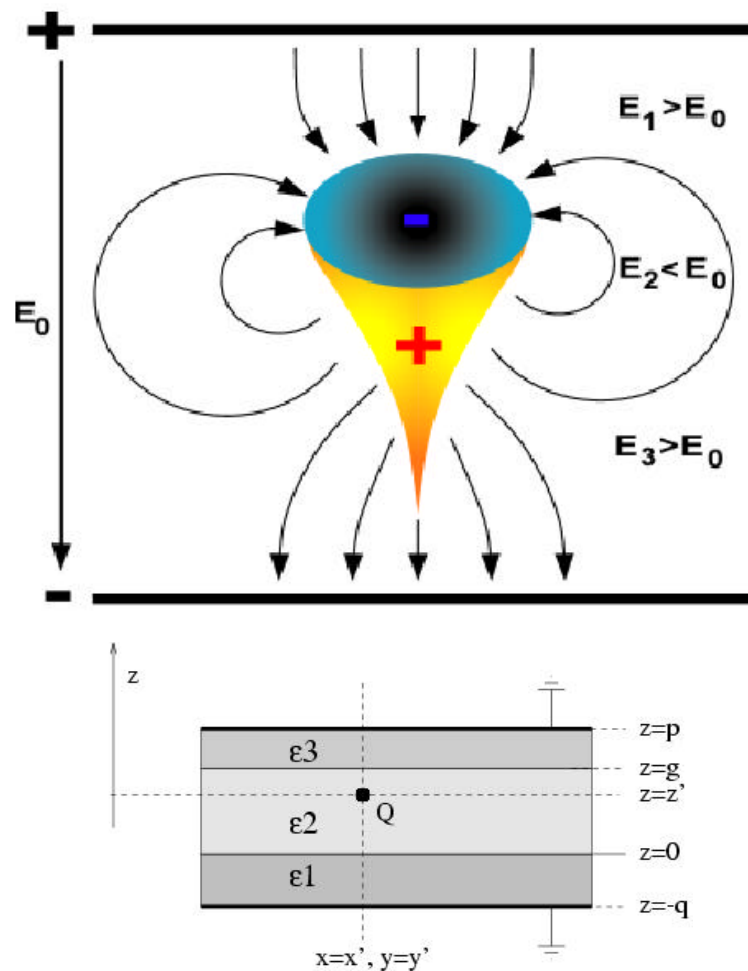
- ◆ At threshold crossing: space charge effect has no influence (for very fast amplifiers)
- ◆ Measurements: open symbols
- ◆ 7GeV pions (9.13 clusters/mm), 20fC threshold, 200ps amplifier peaking time, 1fC noise, T=296.15K and p=970mb





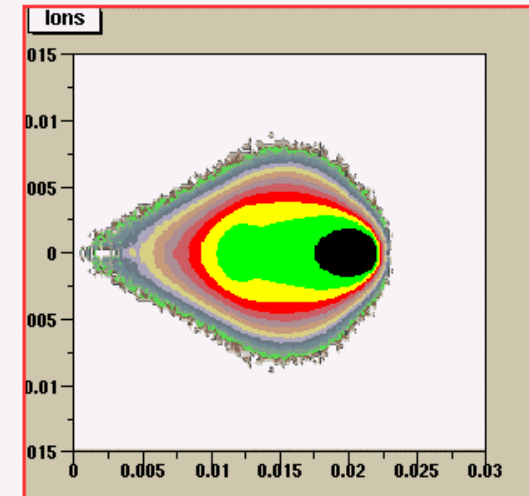
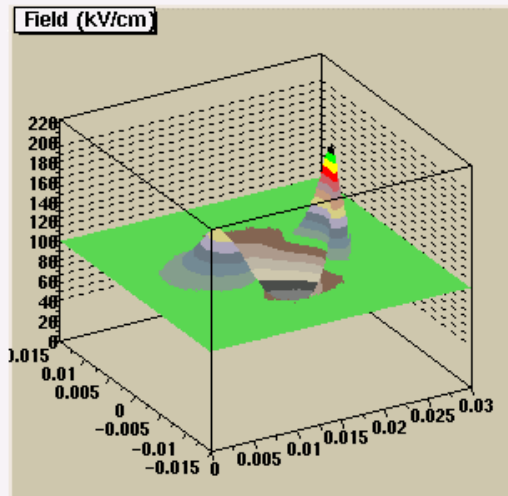
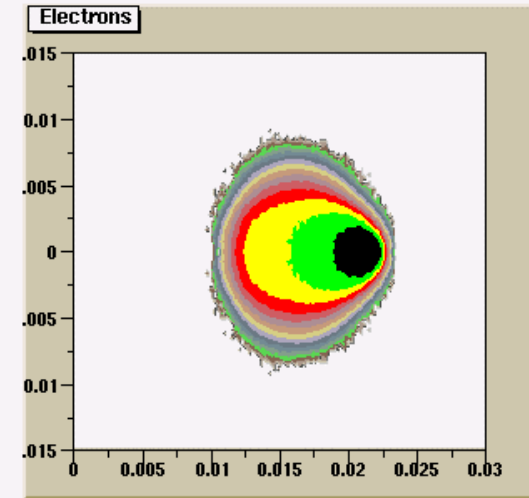
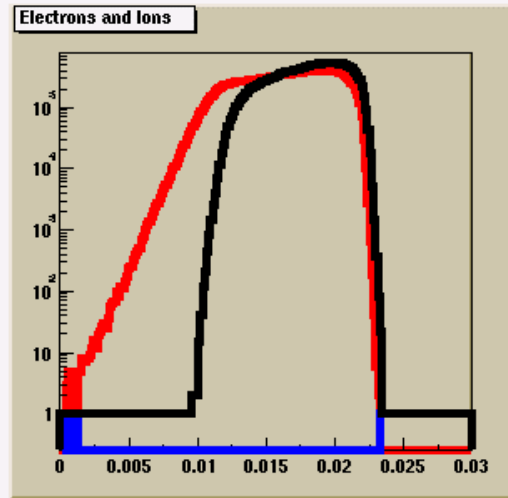
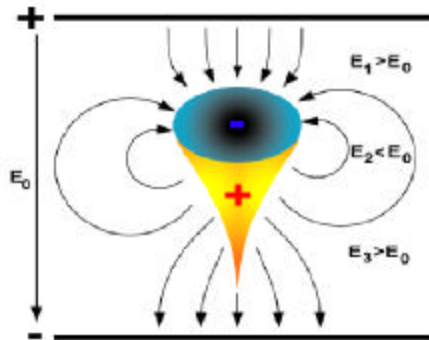
Include Space Charge Effect

- ◆ Divide gas gap into several steps
- ◆ Electrons multiply and drift towards the anode
- ◆ Calculate dynamically the field E of the space charge at each step where electrons are multiplied
- ◆ $a(E)$, $h(E)$, $v_D(E)$
- ◆ Diffusion
- ◆ No photons





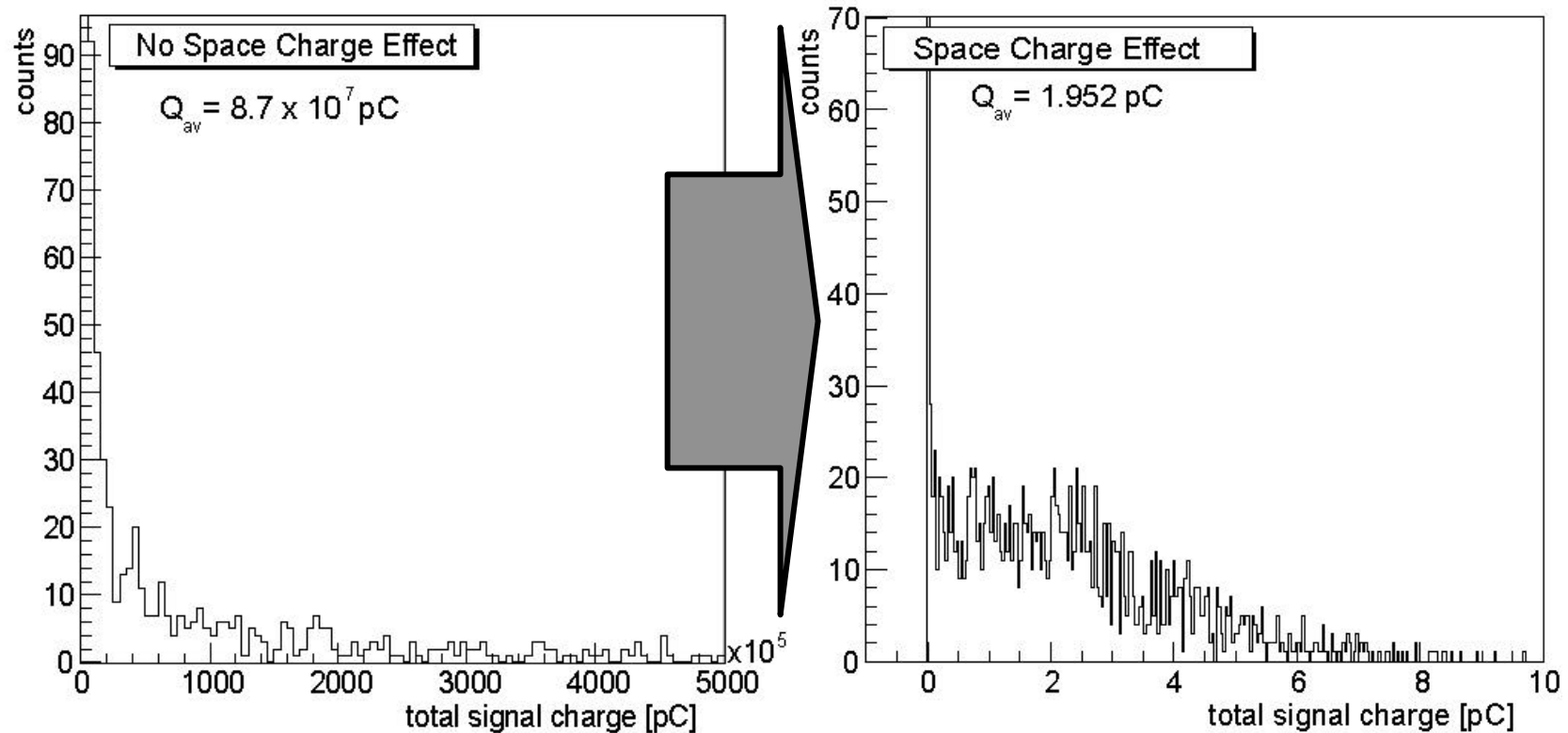
Example Avalanche





Charge Spectra

- ◆ Simulated avalanches at HV=3kV, p=1013mb, single gap RPC
- ◆ Supression Factor 10^7 !!





Reminder: Wire Tube/ Wire Chamber

- ◆ $1/r$ field geometry
 - Space charge region very short ($<100V$)
 - 1.5 orders of magnitude jump to limited streamer region

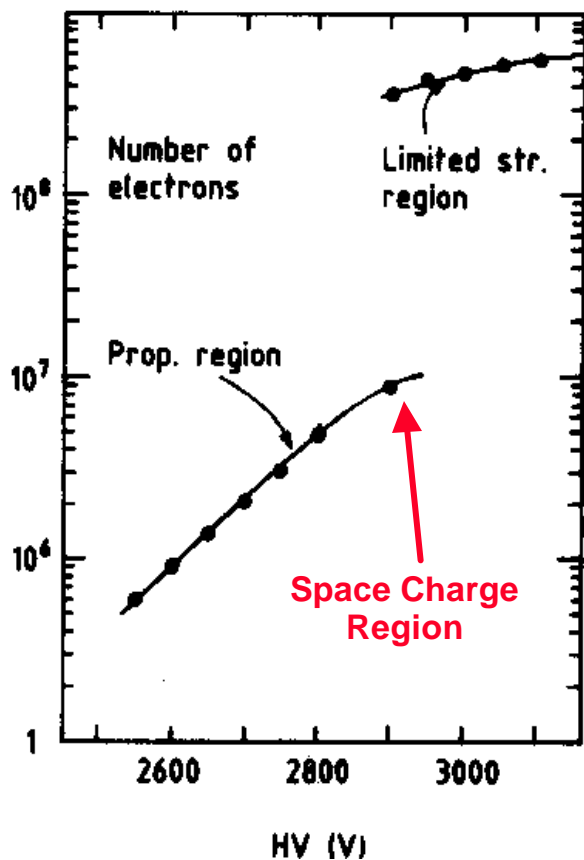
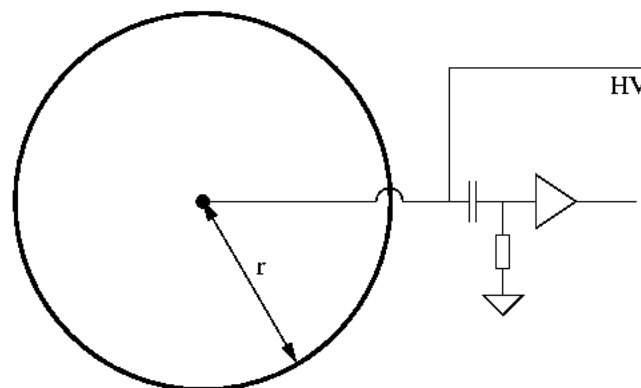


Fig. 4.5. Collected charge as a function of the high voltage, measured by ATAC et al. [ATA 82] on a $100\ \mu\text{m}$ diameter wire in a tube $12 \times 12\ \text{mm}^2$, filled with $\text{Ar}(49.3\%) + \text{C}_2\text{H}_6(49.3\%) + \text{CH}_3\text{CH}_2\text{OH}(1.4\%)$

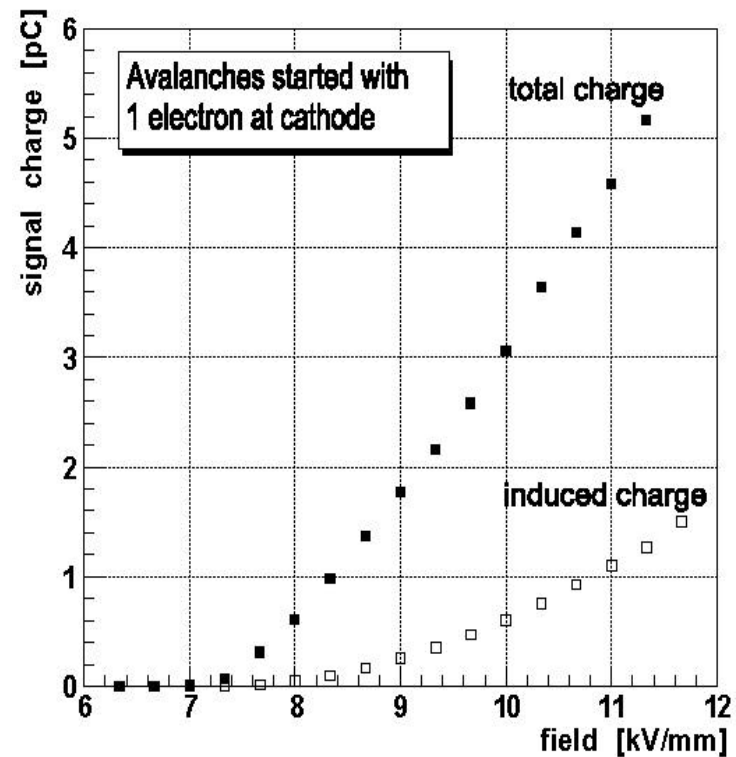
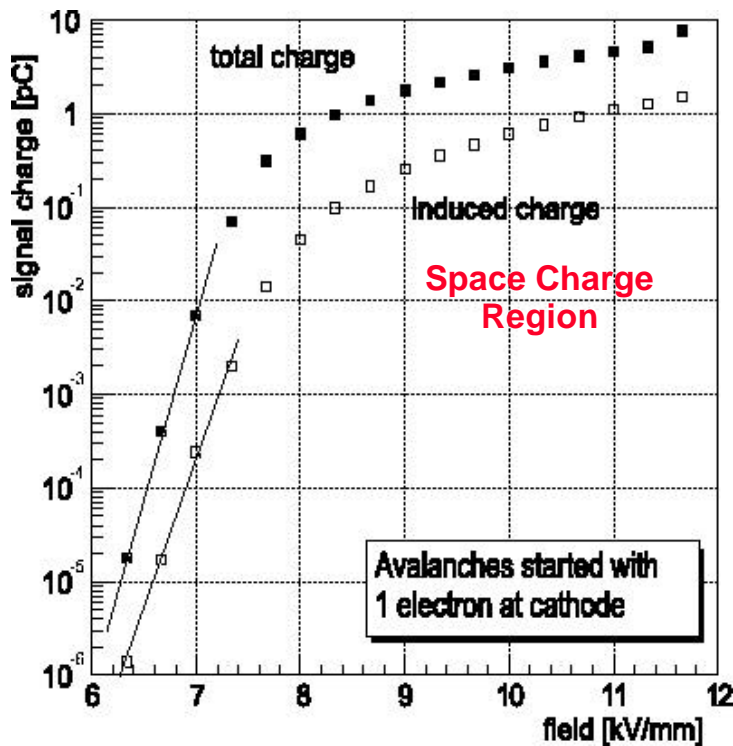
From NIM 200, 345 (1982)





Timing RPC: Long Space Charge Mode

- ◆ Homogeneous (applied) electric field
- ◆ Proportional Region is below Threshold
- ◆ Very long space charge Region ($>700V$)
- ◆ Charge grows first exponentially, then linearly with HV (which is also an experimental fact)



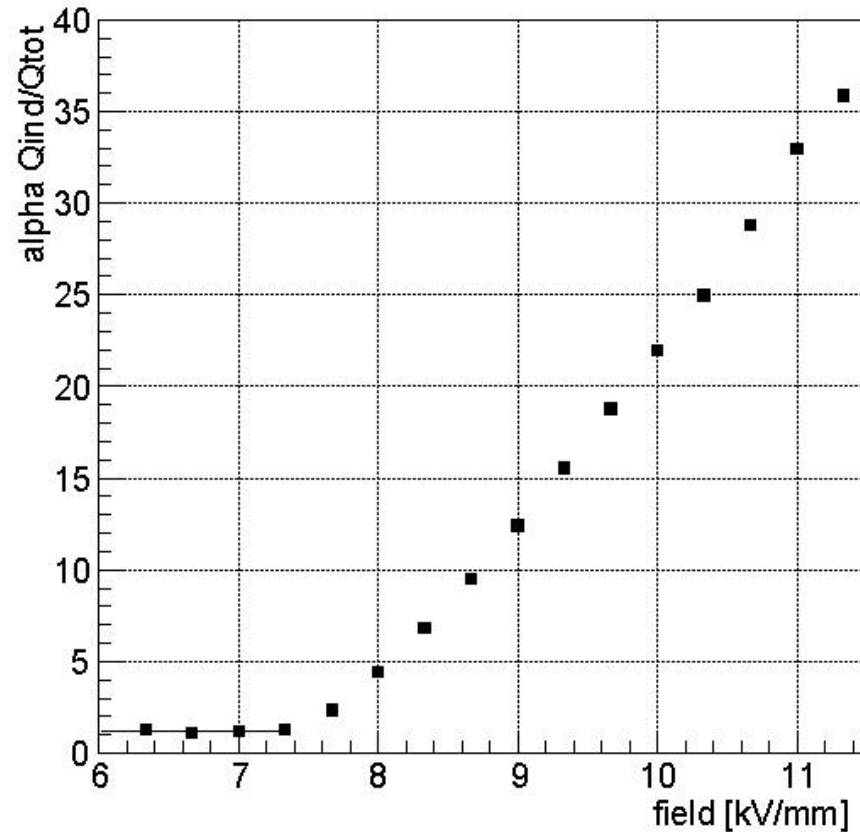


Ratio Q_{ind}/Q_{tot}

- ◆ For avalanches where no space charge effect is present we expect:

$$\frac{Q_{ind}}{Q_{tot}} = \frac{E_w}{\alpha}$$

- ◆ Indicator for a strong space charge effect present for $E > 7.5 \text{ kV/mm}$





Conclusion

- ◆ We have applied **standard detector physics** simulations to Timing RPCs and find **good agreement** with measurements for efficiency, time resolution and charge spectra.
- ◆ The operational mode of timing RPCs is strongly influenced by a space charge effect.
- ◆ The suppression factor is huge (10^7).
- ◆ Details on our work:
 - CERN-EP-2002-024
 - NIM A481(2001) 130-143
 - CERN-EP-2002-046
 - CERN-OPEN-2001-074