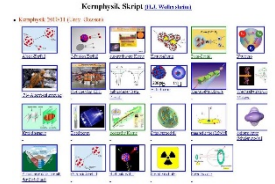


Outline: Interacting Boson Model

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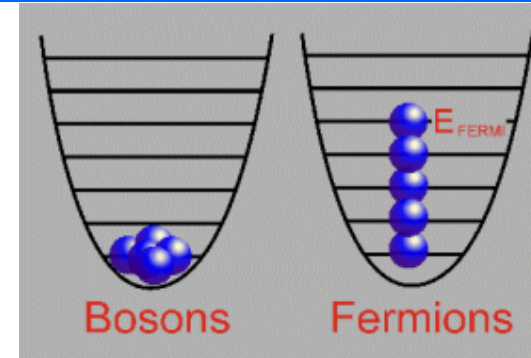


1. particle classification
2. parameters in IBM Hamiltonian
3. vibrational limit $U(5)$
4. rotational limit $SU(3)$
5. extension of IBM

Particle classification

➤ Fermions:

- Spin 1/2 particles
- Obey Pauli Exclusion Principle
 - No more than one fermion can occupy the same quantum state
- Antisymmetric wave functions
- Two classifications
 - Hadrons - interact via strong nuclear force; composed of quarks; neutrons, protons, etc.
 - Leptons - fundamental particles, ie no substructure; electron, muon, tau, neutrino.



➤ Bosons:

- Integer spins
- Obey Bose-Einstein Statistics
 - Any number of bosons can share the same quantum state
- Symmetric wave functions
- Force mediators- Photons, gluons, etc.

Modeling a nucleus

$^{162}_{66}\text{Dy}_{96}$ → shell model →

$\sim 7 \cdot 10^{19}$ 2^+ states



one needs a truncation

Interacting Boson Model

assumptions:

1. only valence nucleons
2. Fermions → Bosons



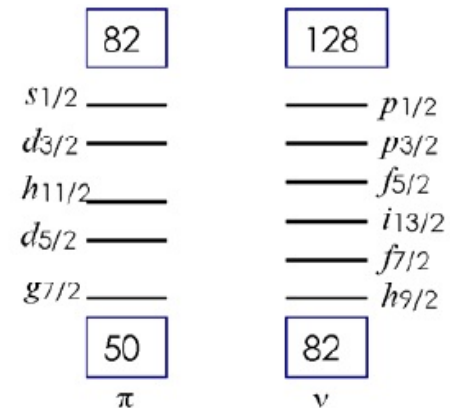
30 basic states 2^+ states



$J = 0$ (s bosons)

$J = 2$ (d bosons)

$^{132}_{50}\text{Sn}_{82}$ inert core



Dimensions

- Assume Ω available 1-fermion states. Number of n -fermion states is

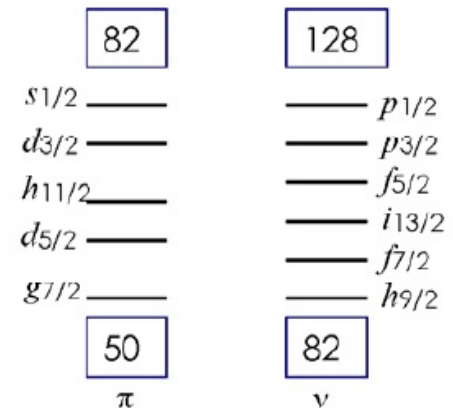
$$\binom{\Omega}{n} = \frac{\Omega!}{n! \cdot (\Omega - n)!}$$

- Assume Ω available 1-boson states. Number of n -boson states is

$$\binom{\Omega + n - 1}{n} = \frac{(\Omega + n - 1)!}{n! \cdot (\Omega - 1)!}$$

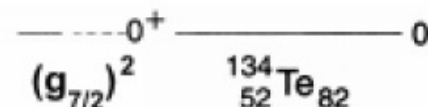
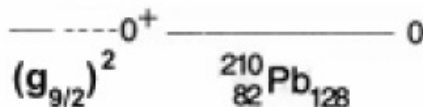
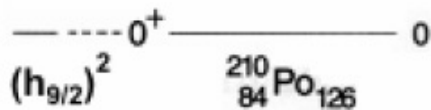
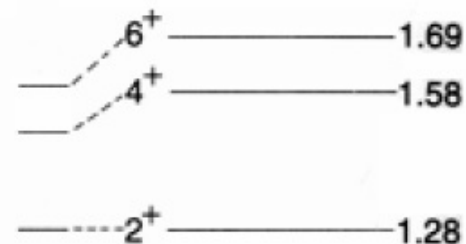
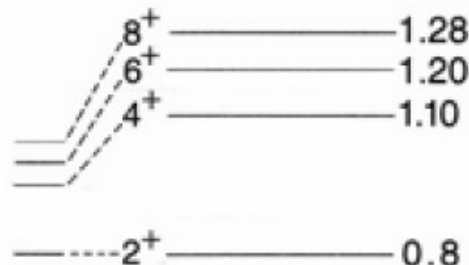
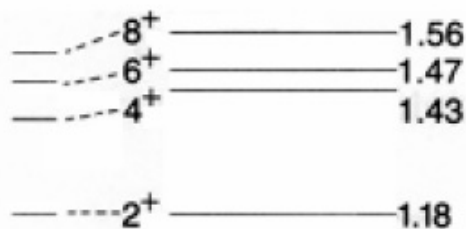
- Example: ${}^{162}_{66}\text{Dy}_{96}$ with 14 neutrons ($\Omega = 44$) and 16 protons ($\Omega = 32$) (${}^{132}\text{Sn}_{82}$ inert core).

- SM dimension: $\sim 7 \cdot 10^{19}$
- IBM dimension: 15504



Why s- and d-bosons?

- ❖ **s-boson** lowest state of all even-even nuclei is 0^+ .
 δ -force gives 0^+ ground state
- ❖ **d-boson** first excited state in non-magic even-even nuclei almost always 2^+
 δ -force gives 2^+ (occupation $\mu=-2,-1,0,1,2$) next above 0^+



Interacting boson model

- Describe the nucleus as a system of N interacting s and d bosons.

Hamiltonian:

$$\hat{H}_{\text{IBM}} = \sum_{i=1}^6 \varepsilon_i \hat{b}_i^+ \hat{b}_i + \sum_{i_1 i_2 i_3 i_4 = 1}^6 v_{i_1 i_2 i_3 i_4} \hat{b}_{i_1}^+ \hat{b}_{i_2}^+ \hat{b}_{i_3} \hat{b}_{i_4} + \dots$$

Since boson number is conserved for a given nucleus, H can only contain “bilinear” terms: 36 of them.

$$s^+ s, s^+ d_m, d_m^+ s, d_m^+ d_m$$



group theoretical
classification of Hamiltonian

U(6)

- Justification from

- Shell model: s and d bosons are associated with S and D fermion (*Cooper*) pairs.
- Geometric model: for large boson number the IBM reduces to a liquid-drop Hamiltonian.

- 36 Operators of IBM all conserve total boson number

$$\begin{aligned} N &= s^+s + d^+\tilde{d} \\ &= n_s + n_d \end{aligned}$$

N is an example of a **Casimir** Operator – Operator that commutes with all the generators of a group (Lie algebra: $[a, b] = ab - ba$).

e.g:

$$\begin{aligned} [N, s^+\tilde{d}]\Psi &= [N(s^+\tilde{d}) - s^+\tilde{d}N]\Psi \\ &= Ns^+\tilde{d}\Psi - s^+\tilde{d}N\Psi \\ &= Ns^+\tilde{d}\Psi - Ns^+\tilde{d}\Psi = 0 \end{aligned}$$

Set of states of the IBA with given N are an irreducible representation of the group $U(6)$ spanned by these 36 operators

Concept of group theory

For IBA, the 36 operators s^+s, d^+s, s^+d, d^+d are generators of the group $U(6)$

$$\begin{aligned}[d^+s, s^+s]|n_d n_s \rangle &= (d^+s s^+s - s^+s d^+s)|n_d n_s \rangle \\ &= d^+s n_s |n_d n_s \rangle - s^+s d^+s |n_d n_s \rangle \\ &= (n_s - s^+s) d^+s |n_d n_s \rangle\end{aligned}$$

Subsets of generators that commute among themselves.

e.g: d^+d 25 generators—span $U(5)$

They conserve n_d (# d bosons)

Set of states with same n_d are the representations of the group $U(5)$

General IBM Hamiltonian

- Most general rotationally invariant IBM Hamiltonian:

$$\hat{H}_{\text{IBM}} = E_0 + \hat{H}_1 + \hat{H}_2 + \hat{H}_3 + \dots$$

$$\hat{H}_1 = \varepsilon_s \hat{n}_s + \varepsilon_d \hat{n}_d$$

$$\hat{H}_2 = \sum_{l_1 l_2 l'_1 l'_2, L} \tilde{v}_{l_1 l_2 l'_1 l'_2}^L \left(b_{l_1}^+ \times b_{l_2}^+ \right)^{(L)} \cdot \left(\tilde{b}_{l'_1} \times \tilde{b}_{l'_2} \right)^{(L)}$$

$$\hat{H}_3 = \sum_{l_1 l_2 l_3 l'_1 l'_2 l'_3, L} \tilde{v}_{l_1 l_2 l_3 l'_1 l'_2 l'_3}^L \left(b_{l_1}^+ \times b_{l_2}^+ \times b_{l_3}^+ \right)^{(L)} \cdot \left(\tilde{b}_{l'_1} \times \tilde{b}_{l'_2} \times \tilde{b}_{l'_3} \right)^{(L)}$$

Parameters in the IBM Hamiltonian

- Spectrum of a single nucleus: 0+1+5+10 parameters.
- Overall binding energy: 1+1+2+7 parameters.
- \Rightarrow A total of **27 parameters** if all interactions up to **three-body** are included (cfr. 63 2-body *sd*-shell model matrix elements).

Order	Number of interactions		
	total	type I ^a	type II ^b
$n = 0$	1	1	0
$n = 1$	2	1	1
$n = 2$	7	2	5
$n = 3$	17	7	10

^aInteraction energy is constant for all states with the same N .

^bInteraction energy varies from state to state.

- Coherent state:

$$|N; \beta, \gamma\rangle \propto \left[s^+ + \beta \cos \gamma d_0^+ + \sqrt{\frac{1}{2}} \beta \sin \gamma (d_{-2}^+ + d_{+2}^+) \right]^N |0\rangle$$

- Generic form of the potential:

$$\begin{aligned} V(\beta, \gamma) &\equiv \langle N; \beta, \gamma | E_0 + \hat{H}_1 + \hat{H}_2 + \hat{H}_3 + \dots | N; \beta, \gamma \rangle \\ &= E_0 + \sum_{n \geq 1} \frac{N(N-1)\dots(N-n+1)}{(1+\beta^2)^n} \sum_{kl} a_{kl}^{(n)} \beta^{2k+3l} \cos^l 3\gamma \end{aligned}$$

β is elongation of nuclear ellipsoidal shape – related to ratio of major and minor axes.

γ is the deviation from axial symmetry along the major axis. It is measured as an angle from 0 to 30 degrees.

Coefficients up to third order

$$a_{00}^{(1)} = \varepsilon_s, \quad a_{10}^{(1)} = \varepsilon_d,$$

$$a_{00}^{(2)} = \frac{1}{2} \mathbf{v}_{ssss}^0, \quad a_{10}^{(2)} = \sqrt{\frac{1}{5}} \mathbf{v}_{ssdd}^0 + \mathbf{v}_{sdsd}^2, \quad a_{01}^{(2)} = -\sqrt{\frac{2}{7}} \mathbf{v}_{sddd}^2,$$

$$a_{20}^{(2)} = \frac{1}{10} \mathbf{v}_{dddd}^0 + \frac{1}{7} \mathbf{v}_{dddd}^2 + \frac{9}{35} \mathbf{v}_{dddd}^4,$$

$$a_{00}^{(3)} = \frac{1}{6} \mathbf{v}_{sssss}^0, \quad a_{10}^{(3)} = \sqrt{\frac{1}{15}} \mathbf{v}_{ssssdd}^0 + \frac{1}{2} \mathbf{v}_{ssdsd}^2,$$

$$a_{01}^{(3)} = -\frac{1}{3} \sqrt{\frac{2}{35}} \mathbf{v}_{sssddd}^0 - \sqrt{\frac{2}{7}} \mathbf{v}_{ssdsd}^2,$$

$$a_{20}^{(3)} = \frac{1}{10} \mathbf{v}_{sddsdd}^0 + \sqrt{\frac{1}{7}} \mathbf{v}_{ssddd}^2 + \frac{1}{7} \mathbf{v}_{sddsdd}^2 + \frac{9}{35} \mathbf{v}_{sddsdd}^4,$$

$$a_{11}^{(3)} = -\frac{1}{5} \sqrt{\frac{2}{21}} \mathbf{v}_{sdddd}^0 - \frac{\sqrt{2}}{7} \mathbf{v}_{sdddd}^2 - \frac{18}{35} \sqrt{\frac{2}{11}} \mathbf{v}_{sdddd}^4,$$

$$a_{30}^{(3)} = \frac{1}{14} \mathbf{v}_{dddddd}^2 + \frac{1}{30} \mathbf{v}_{dddddd}^3 + \frac{3}{154} \mathbf{v}_{dddddd}^4 + \frac{7}{165} \mathbf{v}_{dddddd}^6,$$

$$a_{02}^{(3)} = \frac{1}{105} \mathbf{v}_{dddddd}^0 - \frac{1}{30} \mathbf{v}_{dddddd}^3 + \frac{3}{110} \mathbf{v}_{dddddd}^4 - \frac{4}{1155} \mathbf{v}_{dddddd}^6.$$

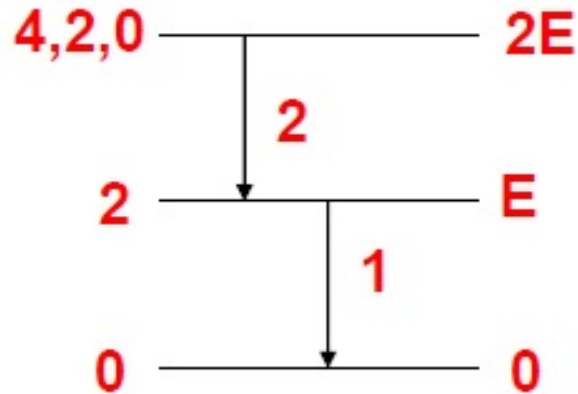
Review of phonon creation and annihilation operators

$$\mathbf{b}|n_b\rangle = \sqrt{n_b} |n_b - 1\rangle$$

$$\mathbf{b}^\dagger|n_b\rangle = \sqrt{(n_b + 1)} |n_b + 1\rangle$$

Why are creation and annihilation operators useful?

Example: Consider a vibrational nucleus with quadrupole phonon excitations, with angular momentum 2. These are vibrational excitations with energy E .



Therefore, the lowest energy state will have $N=0$ of these vibrations. The lowest excited state will have one of them and excitation energy $1 E$.

$$B(E2; 4 \rightarrow 2) = 2 \cdot B(E2; 2 \rightarrow 0)$$

- ❖ $U(5)$ Vibrator
- ❖ $SU(3)$ Axial Rotator
- ❖ $O(6)$ Axially asymmetric Rotor (“gamma-soft”)

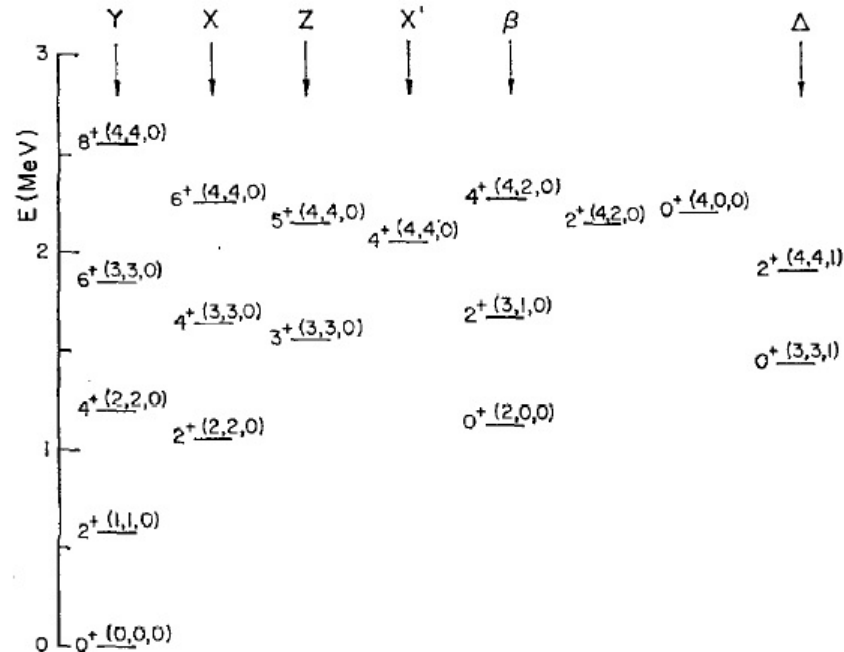
Application to U(5)-like nuclei 4 + 4 parameters

$$E(n_d, \nu, n_\Delta, L, M) = \epsilon \cdot n_d + \frac{\alpha}{2} \cdot n_d(n_d - 1) + \beta(n_d - \nu)(n_d + \nu + 3) + \gamma \cdot [L(L + 1) - 6n_d]$$

eigenvalues do not depend on n_Δ and M .

The parameters α, β, γ can be related to c_0, c_2, c_4 by considering states with $n_d = 2$.

$$\alpha = (1/14)(6c_4 + 8c_2) \quad \beta = (3/70)c_4 - (1/7)c_2 + (1/10)c_0 \quad \gamma = (1/14)(c_4 - c_2)$$



A typical spectrum in the d -boson limit. The numbers in parenthesis are the $SU(5)$ quantum numbers (n_d, ν, n_Δ) . The angular momentum quantum number is explicitly written to the left of each level. The parameters used are $\epsilon = 579$ keV, $c_4 = 39.4$ keV, $c_2 = -95.3$ keV, $c_0 = -27.4$ keV. $\nu = n_d - 2n_\beta$ boson seniority number

Application to U(5)-like nuclei 4 + parameters

B(E2) branching ratios in the Xe-isotopes

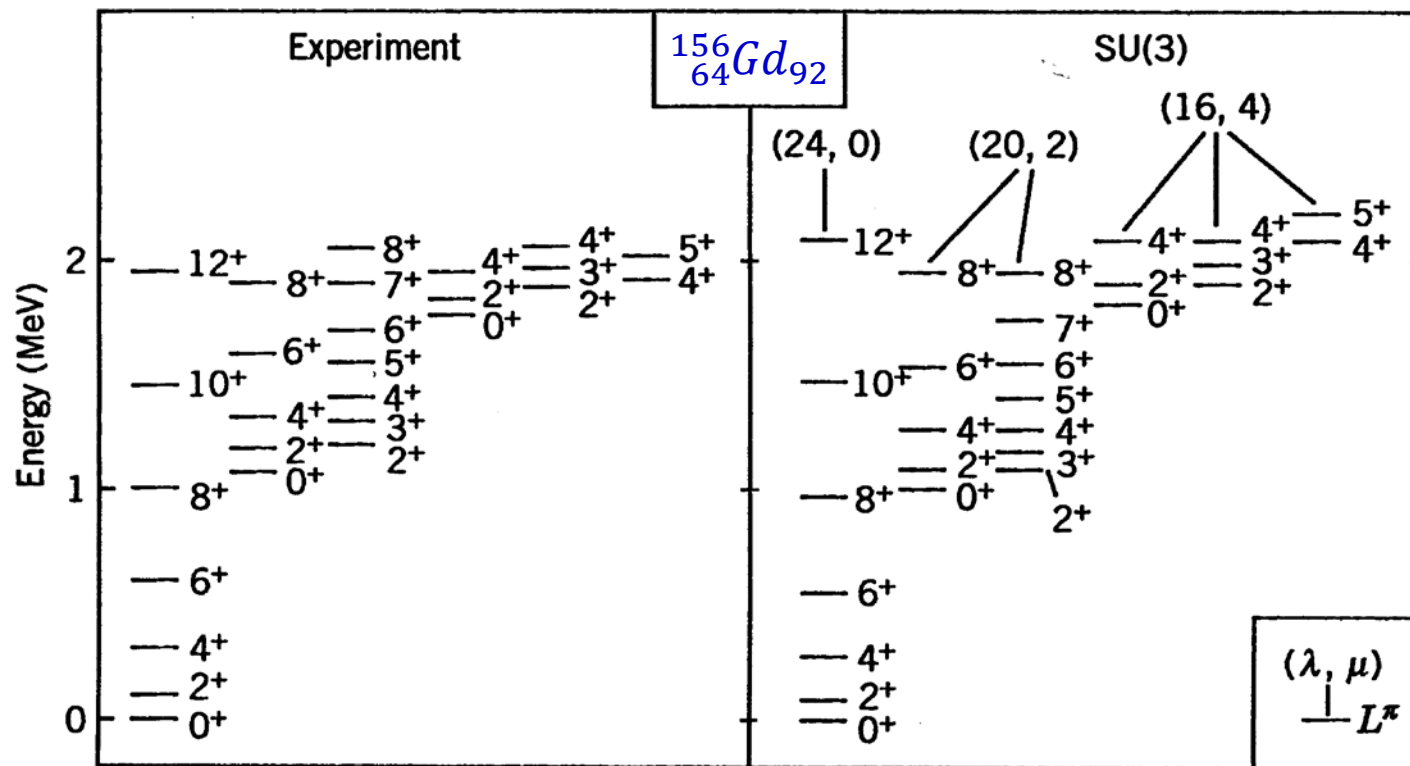
	$^{132}\text{Xe}_{54}^{78}$	$^{130}\text{Xe}_{54}^{76}$	$^{128}\text{Xe}_{54}^{74}$	$^{126}\text{Xe}_{54}^{72}$	$^{124}\text{Xe}_{54}^{70}$	<i>d</i> -boson limit
$\frac{2_2^+ \rightarrow 0_1^+}{2_2^+ \rightarrow 2_1^+} (\times 10^{-2})$	0.17	0.57–0.64	1.01	1.48–1.99	1.86 – 3.89	0
$\left\{ \begin{array}{l} 3_1^+ \rightarrow 2_1^+ \\ 3_1^+ \rightarrow 2_2^+ \end{array} \right. (\times 10^{-2})$	1.03	1.41–1.54	—	1.99	2.99	0
	0.51	0.24–0.25	—	0.46–0.72	0.16	$\frac{2}{5} = 0.400$
$\left\{ \begin{array}{l} 4_2^+ \rightarrow 2_1^+ \\ 4_2^+ \rightarrow 2_2^+ \end{array} \right. (\times 10^{-2})$	—	3.11–3.42	—	1.28	—	0
	—	0.95–1.05	—	0.94	0.90	$\frac{10}{11} = 0.909$
$\left\{ \begin{array}{l} 5_1^+ \rightarrow 4_1^+ \\ 5_1^+ \rightarrow 3_1^+ \end{array} \right. (\times 10^{-2})$	—	3.38	—	4.97	3.90	0
	—	0.46	—	1.26	0.98	$\frac{5}{11} = 0.454$
$\left\{ \begin{array}{l} 5_1^+ \rightarrow 6_1^+ \\ 5_1^+ \rightarrow 3_1^+ \end{array} \right.$	—	0.64	—	—	—	$\frac{104}{231} = 0.450$

Application to SU(3)-like nuclei 2 + parameters

$$E([N](\lambda, \mu)KLM) = \alpha \cdot L(L + 1) - \beta \cdot C(\lambda, \mu) \text{ with } C(\lambda, \mu) = \lambda^2 + \mu^2 + \lambda\mu + 3(\lambda + \mu)$$

eigenvalues do not depend on K and M. $C(\lambda, \mu)$ is the quadratic Casimir operator.

The parameters α, β can be related to κ and κ' : $\alpha = \frac{3}{4}\kappa - \kappa'$ and $\beta = \kappa$

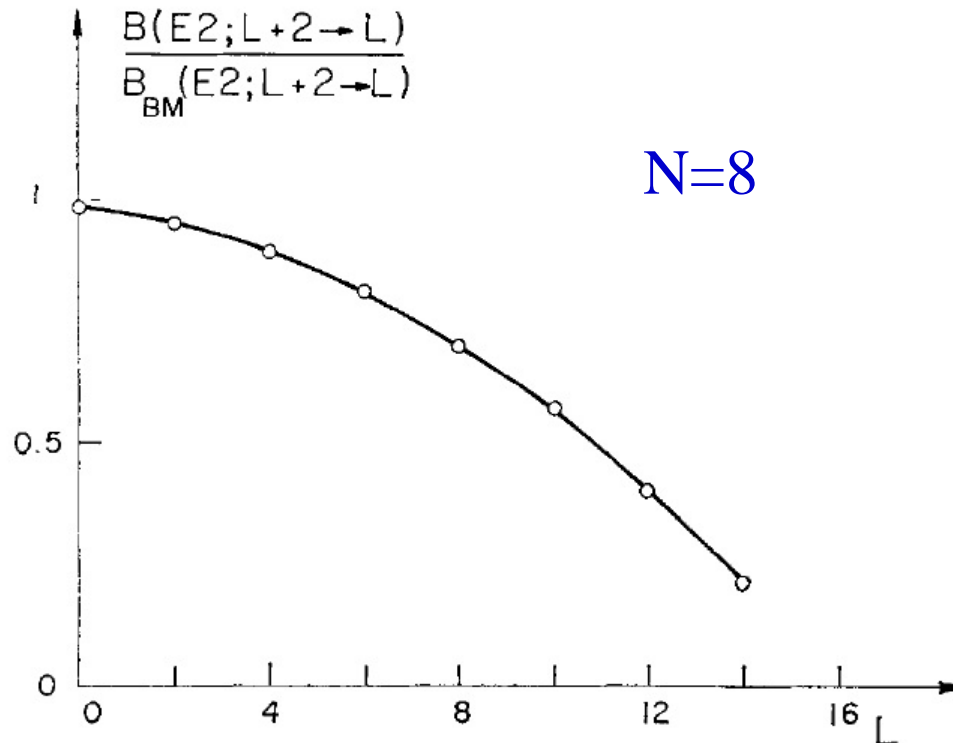


The energy difference between the lowest $(2N, 0)$ and the next lowest $(2N-4, 2)$ representation is in units of κ , $12N-6$

Application to SU(3)-like nuclei 2 + 3 parameters

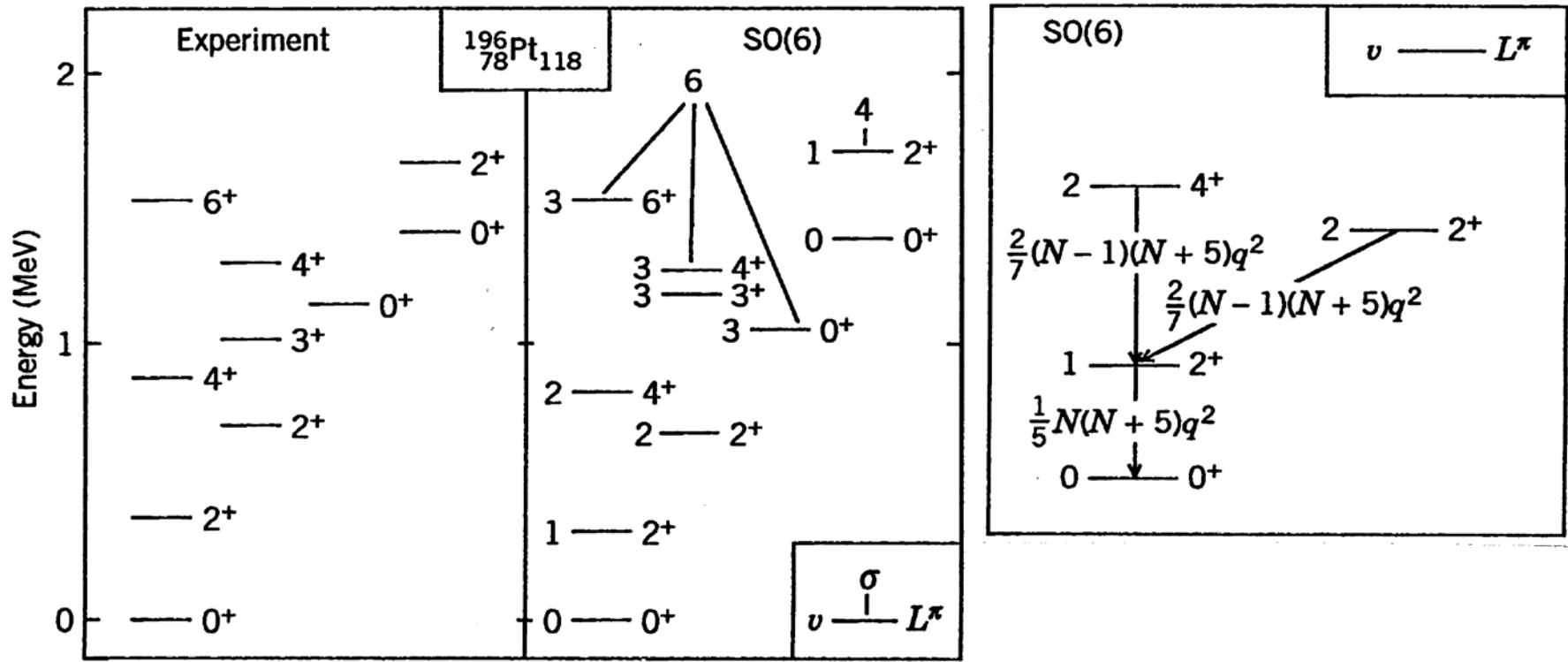
$$B(E2; L + 2 \rightarrow L) = \alpha_2^2 \frac{3}{4} \frac{(L + 2)(L + 1)}{(2L + 3)(2L + 5)} (2N - L)(2N + L + 3)$$

$$B(E2; L + 2 \rightarrow L) = B_{BM}(E2; L + 2 \rightarrow L) \frac{(2N - L)(2N + L + 3)}{(2N + \frac{3}{2})^2}$$



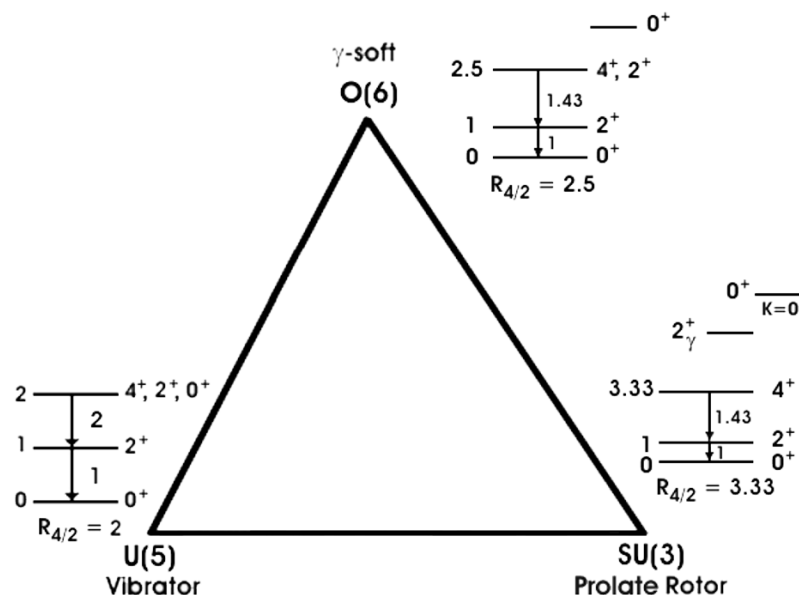
Application to $SO(6)$ γ -unstable nuclei 2 + parameters

- Rotation-vibration spectrum of a γ -unstable body.
- Conserved quantum numbers: σ , ν , L .



Synopsis of IBM symmetries

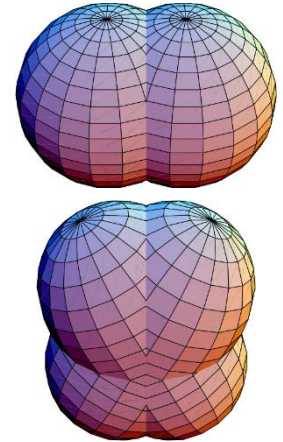
- Symmetry triangle of the IBM:
 - Three standard solutions: U(5), SU(3), SO(6).
 - SU(1,1) analytic solution for U(5) \rightarrow SO(6).
 - Hidden symmetries (parameter transformations).
 - Deformed-spherical coexistent phase.
 - Partial dynamical symmetries.
 - Critical-point symmetries?



Extension of IBM

- Neutron and proton degrees freedom (IBM-2):

- F -spin multiplets ($N_\nu + N_\pi = \text{constant}$)
- Scissors excitations



- Fermion degrees of freedom (IBFM):

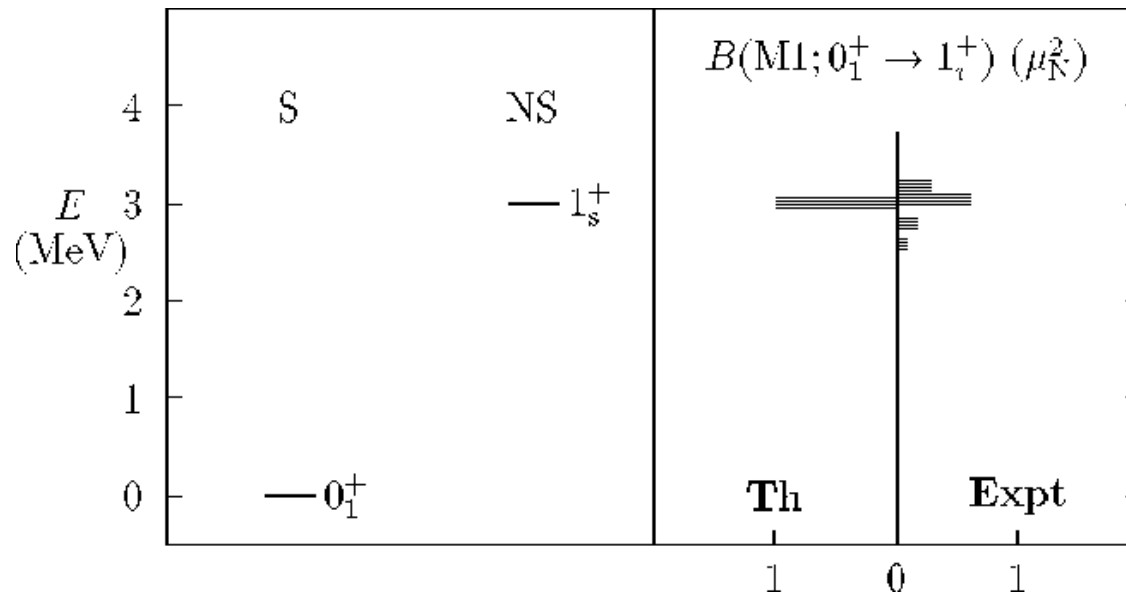
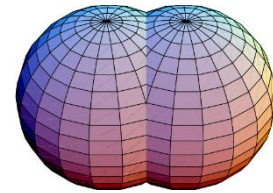
- Odd-mass nuclei
- Supersymmetry (doublets & quartets) ${}^{194}_{78}\text{Pt}_{116}$ ${}^{195}_{78}\text{Pt}_{117}$ ${}^{195}_{79}\text{Au}_{116}$ ${}^{196}_{79}\text{Au}_{117}$

- Other boson degrees of freedom:

- Isospin $T=0$ & $T=1$ pairs (IBM-3 & IBM-4)
- Higher multipole (g,...) pairs

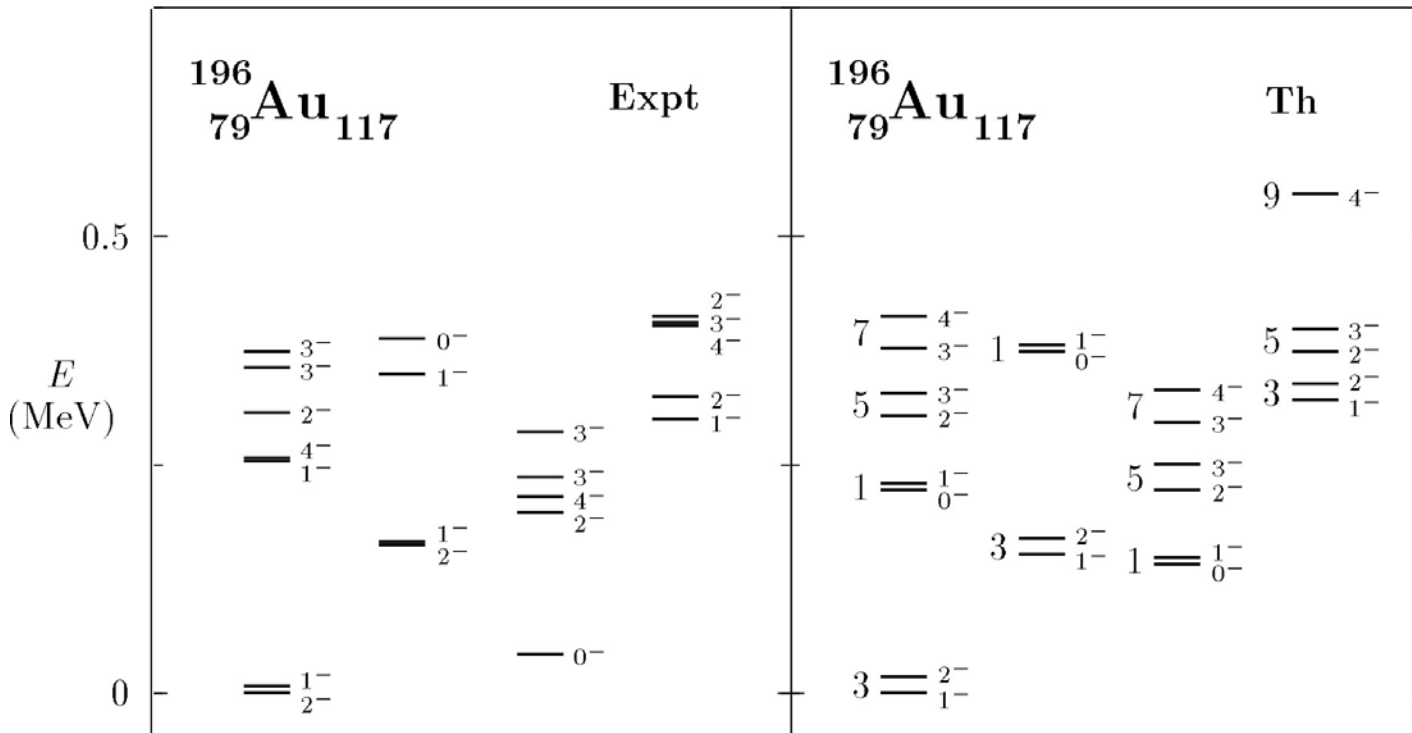
- Collective displacement modes between neutrons and protons:

- *Linear* displacement (giant dipole resonance):
 $R_{\nu} - R_{\pi} \Rightarrow E1$ excitation.
- *Angular* displacement (scissor resonance):
 $L_{\nu} - L_{\pi} \Rightarrow M1$ excitation



Supersymmetry

- A simultaneous description of even- and odd-mass nuclei (*doublets*) or of even-even, even-odd, odd-even and odd-odd nuclei (*quartets*).
- Example of $^{194}_{78}\text{Pt}_{116}$, $^{195}_{78}\text{Pt}_{117}$, $^{195}_{79}\text{Au}_{116}$ & $^{196}_{79}\text{Au}_{117}$



- Nuclear spectra:
 - Systematic use of band-plotting option in NuDat
⇒ important for identification of collective bands.
 - Occasional problems:
 - Gamma-band structure in ^{160}Dy and ^{170}Yb ??
 - Gamma band cut in two in ^{180}W ??
 - 2^+ @ 691 keV in ^{174}Os should be in beta band??
 - Beta band cut in two in ^{180}Os ??
 - In general: confusion concerning beta band in deformed nuclei.
What are the correct criteria to label a $K^\pi = 0^+$ band as collective beta band?