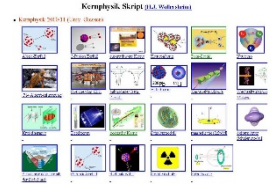


Outline: Radioactivity

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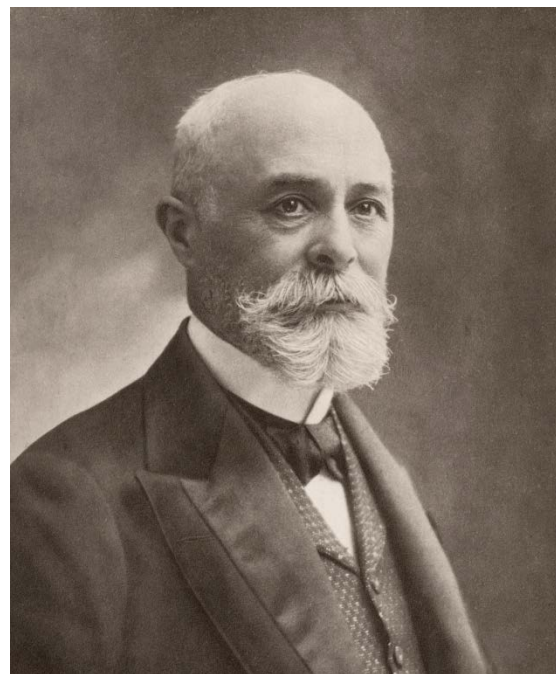
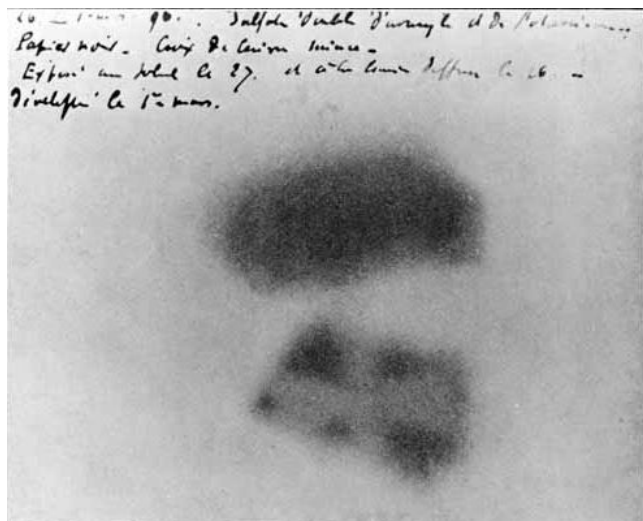


1. history
2. primary modes of radioactive decays
3. radioactive decay law
4. radioactive activity
5. radiation protection

Radioactive decay

Radioactive Decay: is the process by which a nucleus of an unstable atom decreases its total energy by spontaneously emitting radiation.

- Discovered in 1897 by the French scientist Henri Becquerel.
- He placed various phosphorescent salts on photographic plates and noticed that only uranium salts exposed the plates.



Henri Becquerel

Characterization of radioactive material



Marie Curie

Henri Becquerel's discovery inspired Marie and Pierre Curie to investigate the phenomena further.

- They found that the mineral pitchblende was more active than uranium, and concluded that it must contain different radioactive substances
- From this, they discovered polonium and radium, both of them more radioactive than uranium.

The Nobel Prize in Physics 1903



Antoine Henri
Becquerel

Prize share: 1/2



Pierre Curie

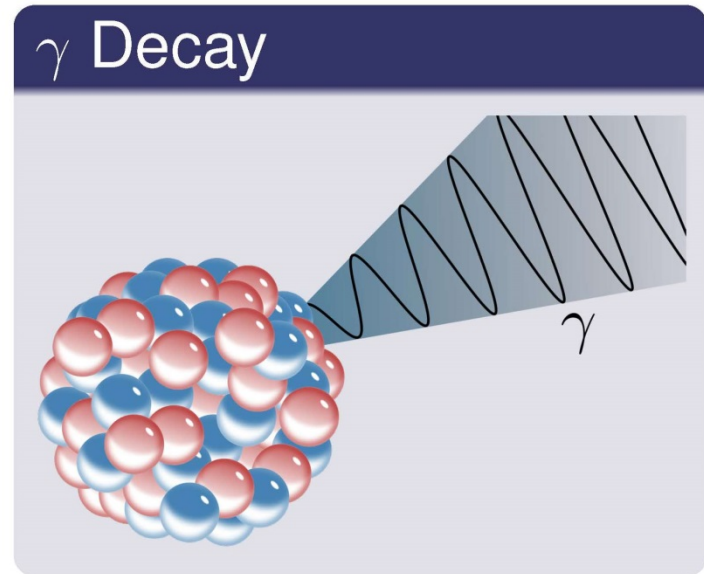
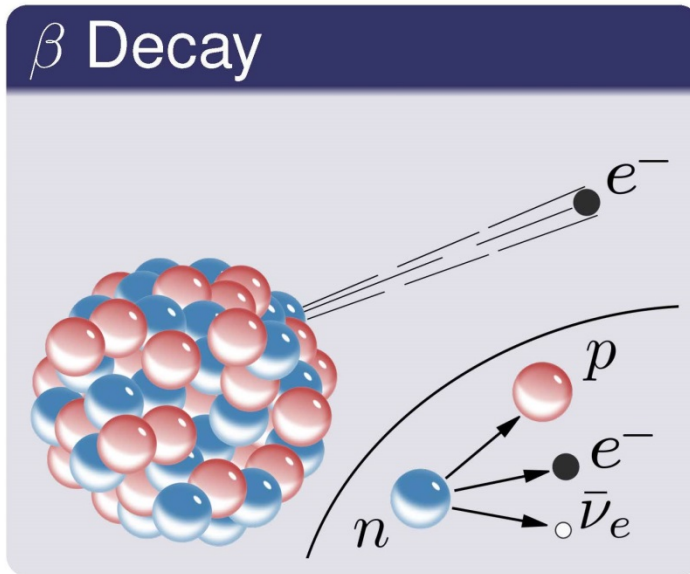
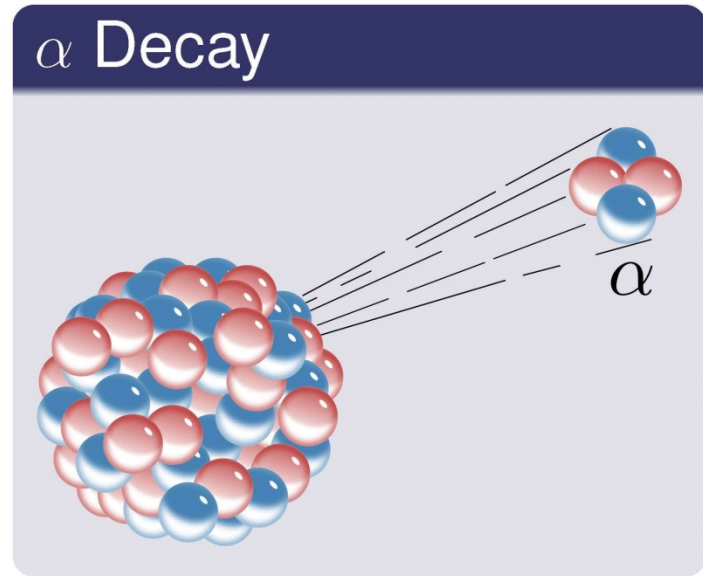
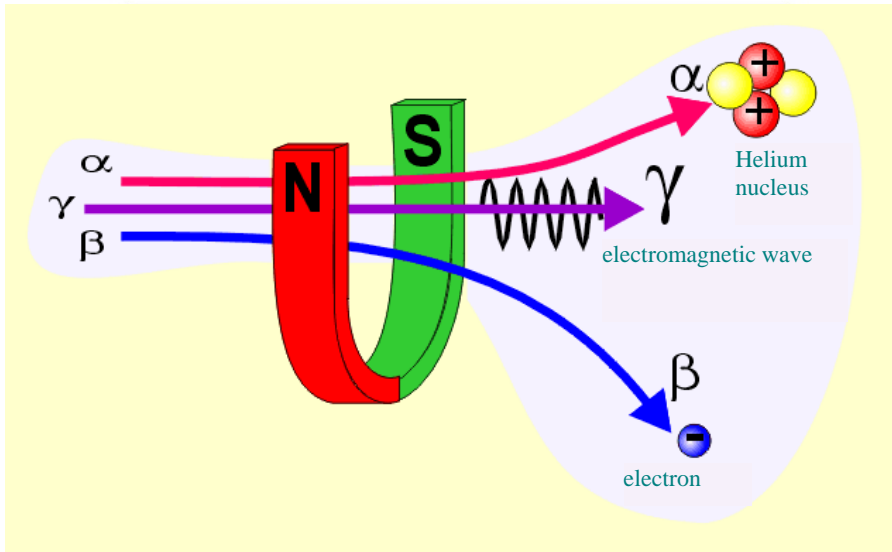
Prize share: 1/4



Marie Curie,
née Sklodowska

Prize share: 1/4

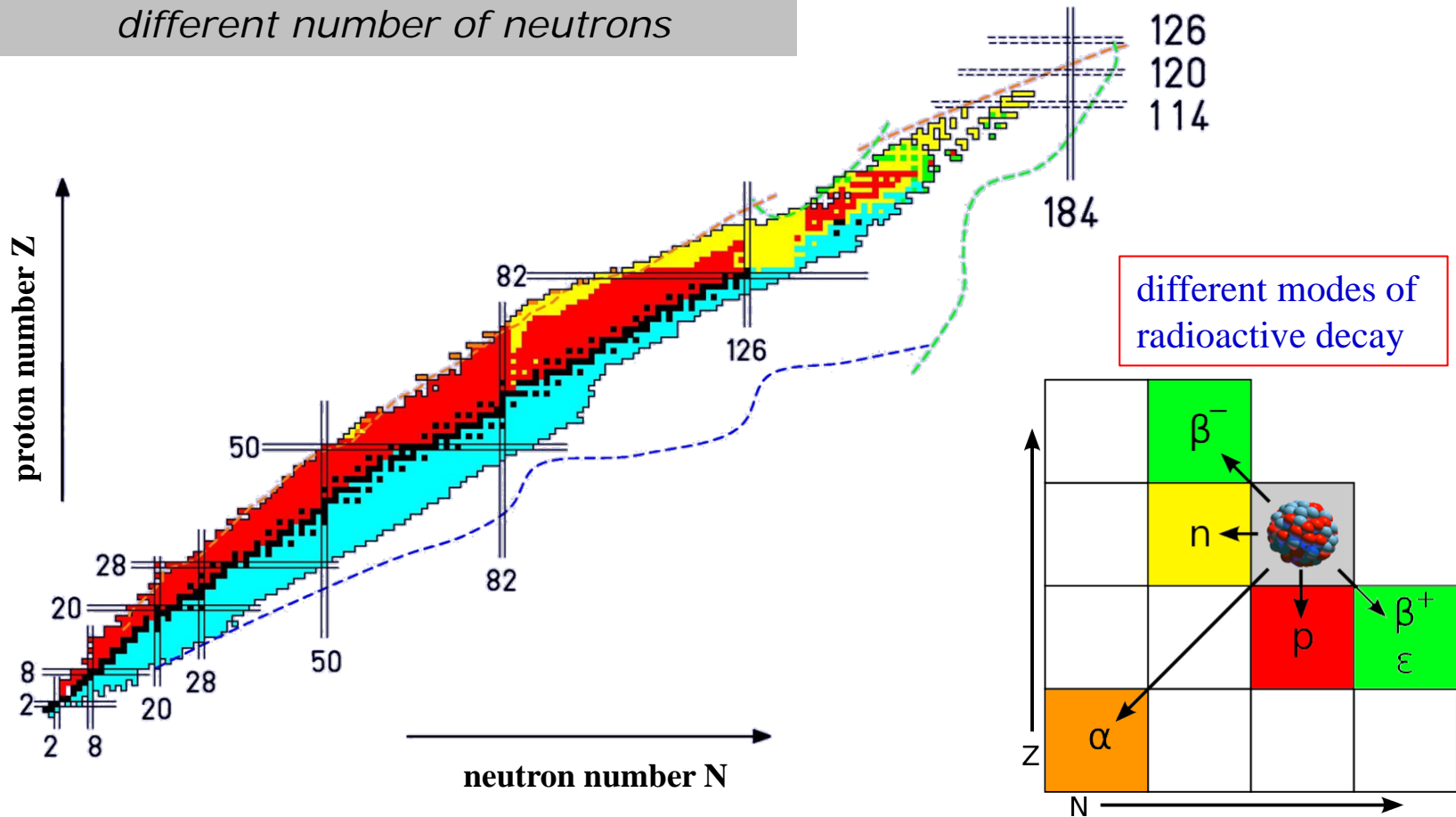
Primary modes of radioactive decay



Primary modes of radioactive decay

chart of nuclides:

- representation of isotopes in the Z - N plane
- isotope: atom (nucleus) of an element with different number of neutrons



The radioactive decay law

The number of nuclei that decay is proportional to the number of radioactive nuclei N in the sample:

$$-\frac{dN}{dt} = \lambda \cdot N$$

λ is the decay constant

- Radioactivity is an exponential decay process

$$N(t) = N_0 \cdot e^{-\lambda \cdot t}$$

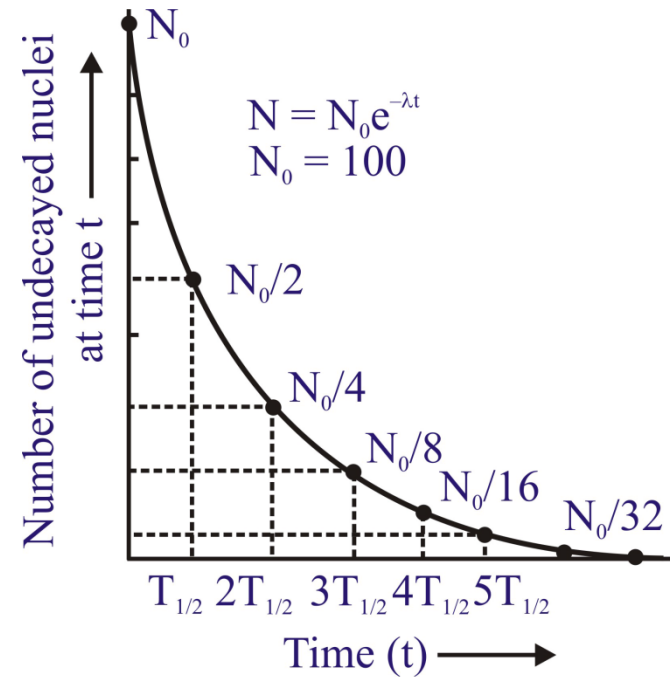
- The sample **activity** or decay rate is given by

$$A = \lambda \cdot N$$

- The lifetime and the half-life is defined as

$$\tau = \frac{1}{\lambda}$$

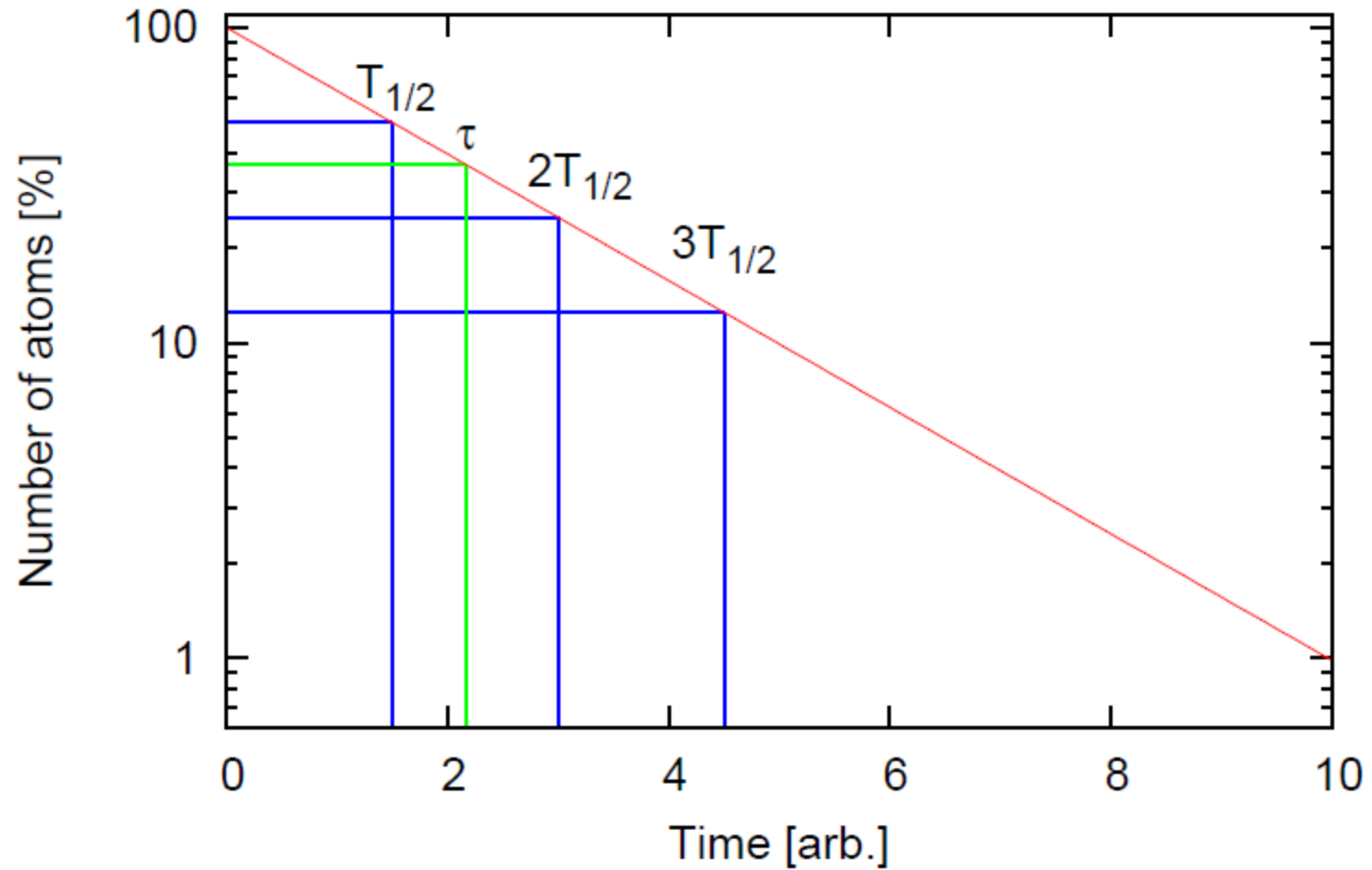
$$T_{1/2} = \frac{\ln 2}{\lambda} = 0.693 \cdot \tau$$



Decay curve for a radioactive element

Nuclear decay

$$T_{1/2} = 1.5, \tau = T_{1/2} / \ln(2) = 2.16$$



Nuclear decay with branching

- ❖ The initial nucleus can decay into either of two products, e.g. in a sample of ^{212}Bi , 36% of the nuclei decay to ^{208}Tl and 64% to ^{212}Po .

Consider the case $A \rightarrow B$ and $A \rightarrow C$ in parallel.
We have for all time t :

$$N = N_A + N_B + N_C$$

which is constant, since the total number of nuclides remain constant.

Differentiating with respect to time:

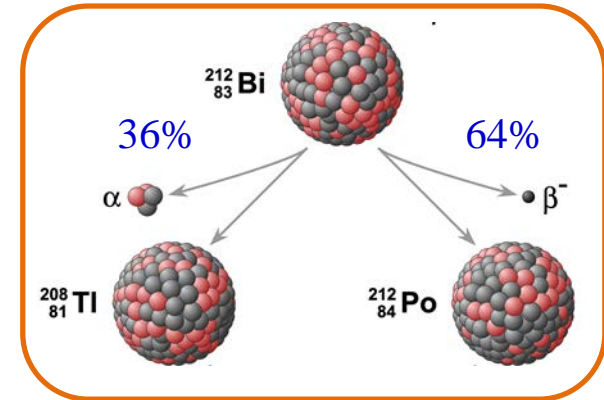
$$\frac{dN_A}{dt} = -\lambda_1 \cdot N_A - \lambda_2 \cdot N_A = N_A(0) \cdot e^{-(\lambda_1 + \lambda_2) \cdot t} \quad \lambda = \lambda_1 + \lambda_2$$

Definition of the **branching ratio**:

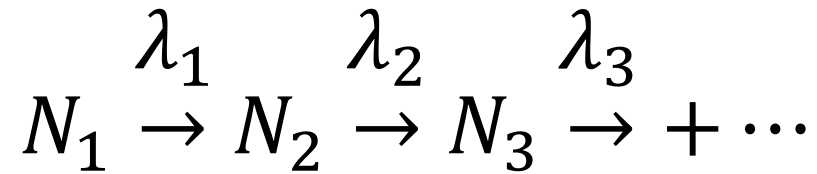
$$f_1 = \frac{\lambda_1}{\lambda} \quad f_2 = \frac{\lambda_2}{\lambda}$$

$$\beta\text{-activity: } A_\beta(t) = N_A \cdot \lambda \cdot f_1 = N_A \cdot \lambda_1$$

$$\alpha\text{-activity: } A_\alpha(t) = N_A \cdot \lambda \cdot f_2 = N_A \cdot \lambda_2$$



Radioactive decay chains



Ultimately, the characterization of multiple decay chains boils down to solving a series of coupled differential equations:

$$\frac{dN_1}{dt} = -\lambda_1 N_1$$

$$\frac{dN_2}{dt} = \lambda_1 N_1 - \lambda_2 N_2$$

$$\frac{dN_3}{dt} = \lambda_2 N_2 - \lambda_3 N_3$$

General decay chains

One can show that a general solution can be obtained as

$$N_1(t) = a_{11}e^{-\lambda_1 t}$$

$$N_2(t) = a_{21}e^{-\lambda_1 t} + a_{22}e^{-\lambda_2 t}$$

$$N_3(t) = a_{31}e^{-\lambda_1 t} + a_{32}e^{-\lambda_2 t} + a_{33}e^{-\lambda_3 t}$$

with initial conditions $a_{11} = N_1(t = 0)$ $\sum_{i=1}^k a_{ki} = 0$

For the particular case of $N_2(t)$ we obtain:

$$\frac{dN_2}{dt} = \lambda_1 N_1 - \lambda_2 N_2$$

$$-\lambda_1 a_{21}e^{-\lambda_1 t} - \lambda_2 a_{22}e^{-\lambda_2 t} = \lambda_1 a_{11}e^{-\lambda_1 t} - \lambda_2 a_{21}e^{-\lambda_1 t} - \lambda_2 a_{22}e^{-\lambda_2 t}$$

$$a_{21} = \frac{\lambda_1}{\lambda_2 - \lambda_1} a_{11}$$

$$a_{22} = \frac{\lambda_1}{\lambda_1 - \lambda_2} a_{11}$$

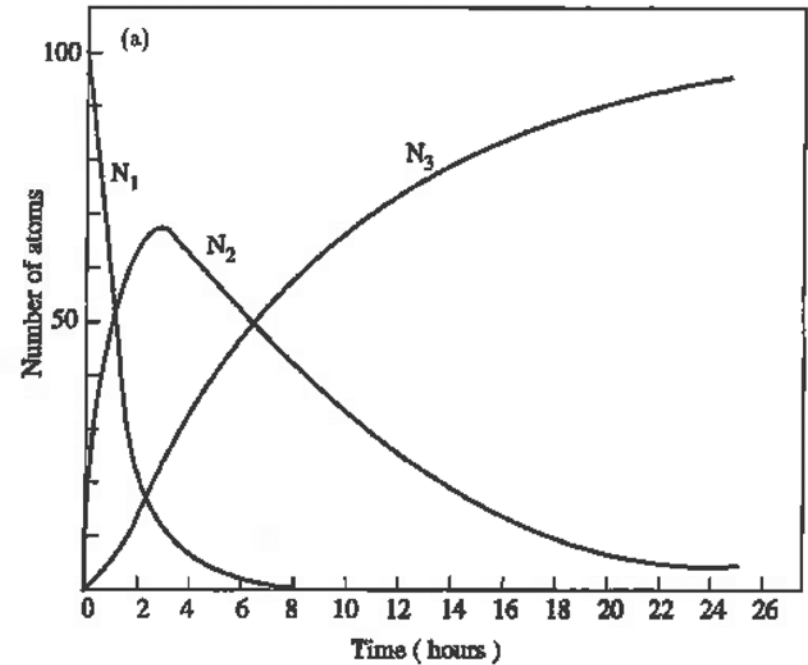
General solution:

$$a_{ki} = \frac{\lambda_{k-1}}{\lambda_k - \lambda_i} a_{k-1,i}$$

$$a_{ki} = \frac{\lambda_1 \lambda_2 \cdots \lambda_{k-1}}{(\lambda_1 - \lambda_i)(\lambda_2 - \lambda_i) \cdots (\lambda_k - \lambda_i)} a_{11}$$

The Bateman Equation

The results for a radioactive series with
 $T=1\text{h}$, $T=5\text{h}$ and $T=\infty$



$$N_1(t) = N_1^0 e^{-\lambda_1 t}$$

$$N_2(t) = \frac{\lambda_1}{\lambda_2 - \lambda_1} N_1^0 (e^{-\lambda_1 t} - e^{-\lambda_2 t})$$

$$N_3(t) = \frac{\lambda_1 \lambda_2}{\lambda_2 - \lambda_1} N_1^0 \left(\frac{1 - e^{-\lambda_1 t}}{\lambda_1} - \frac{1 - e^{-\lambda_2 t}}{\lambda_2} \right)$$

The Bateman Equation

If we extend our solution to k decays, we can generalize our solution
This is referred to as the **Bateman equation**.

$$N_k(t) = \sum_{i=1}^k \left[N_i(0) \cdot \left(\prod_{j=1}^{k-1} \lambda_j \right) \cdot \left(\sum_{J=1}^k \left(\frac{e^{-\lambda_j t}}{\prod_{p=i, p \neq j}^k (\lambda_p - \lambda_j)} \right) \right) \right]$$

Activity

- ❖ For lifetimes comparable or longer than the span of a human life there are no measurable changes in the activity of a sample which prohibits direct decay curve measurements. In these cases the **decay rates** λ are deduced from the ratio of the observed **activity** A to the absolute **number of radioactive atoms** N in a sample.

$$A = \lambda \cdot N = \lambda \cdot \frac{m}{\mu} \cdot N_A = \frac{1}{\tau} \cdot \frac{m}{\mu} \cdot N_A$$

with the molar mass μ and the Avogadro number N_A

- ❖ For example 1 MBq of tritium (${}^3\text{H}$, $T_{1/2}=12.33$ [y]) corresponds to $5.59 \cdot 10^{14}$ atoms or 2.78 [ng] of mass.
- ❖ 1 MBq of ${}^{14}\text{C}$ ($T_{1/2}=5730$ [y]) corresponds to $2.6 \cdot 10^{17}$ atoms or 6 [μg] of mass.

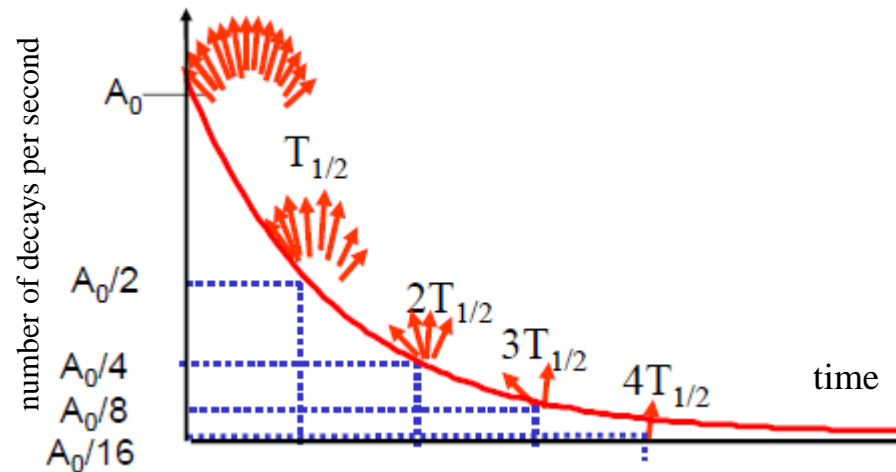
$$A = \frac{\ln 2}{T_{1/2}} \cdot N \quad \rightarrow \quad N = \frac{A \cdot T_{1/2}}{\ln 2} = \frac{10^6 \cdot 12.33 \cdot 365 \cdot 24 \cdot 60 \cdot 60}{0.693} = 5.6 \cdot 10^{14} \text{ atoms}$$

Activity

The number of radioactive nuclei is often unknown. One can only measure the number of decays per second (decay rate). This is the reason why the activity is defined as the number of decays per second:

$$A(t) = -\frac{\Delta N(t)}{\Delta t} = A_0 \cdot e^{-\lambda \cdot t}$$

1 Becquerel (Bq) = 1 decay/s



When at the beginning $N(t=0)$ radioactive nuclei were present, then the activity A after the time t is given by

$$A(t) = -\frac{dN}{dt} = \lambda \cdot N_0 \cdot e^{-\lambda \cdot t} = \lambda \cdot N = \frac{N}{\tau} = A_0 \cdot e^{-\lambda \cdot t}$$

Activity

Activity at the beginning:

$$A_0$$

Activity at time t:

$$A_1 = A_0 \cdot e^{-\lambda \cdot t} \quad \rightarrow \quad \frac{A_1}{A_0} = e^{-\lambda \cdot t}$$

Activity at time 2t:

$$A_2 = A_0 \cdot e^{-\lambda \cdot 2t} = A_1 \cdot e^{-\lambda \cdot t} = A_1 \cdot \left(\frac{A_1}{A_0}\right)$$

Activity at time 3t:

$$A_3 = A_0 \cdot e^{-\lambda \cdot 3t} = A_2 \cdot e^{-\lambda \cdot t} = A_2 \cdot \left(\frac{A_1}{A_0}\right) = A_1 \cdot \left(\frac{A_1}{A_0}\right)^2$$

General expression:

$$A_{n+1} = A_n \cdot \left(\frac{A_1}{A_0}\right) = A_1 \cdot \left(\frac{A_1}{A_0}\right)^n$$

Example: The activity of a sample was measured 1000 [Bq] one hour ago. The present activity is 900 [Bq]. What is the activity in one hour?

$$A_2 = A_1 \cdot \left(\frac{A_1}{A_0}\right) = 900 \cdot \left(\frac{900}{1000}\right) = 810$$

Quantum theory of radiative decays

Consider a general quantum state: $\psi = c_1\psi_1 + c_2\psi_2$

where the probability $|c_1|^2$ is to be found in state 1, and $|c_2|^2$ is to be found in state 2.

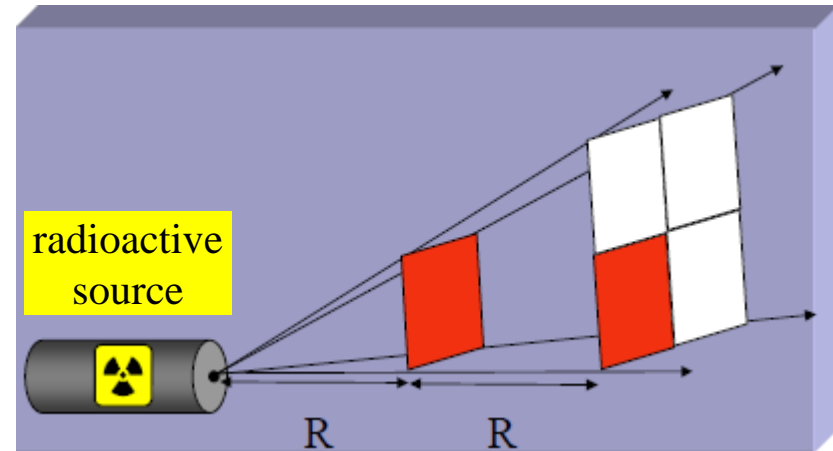
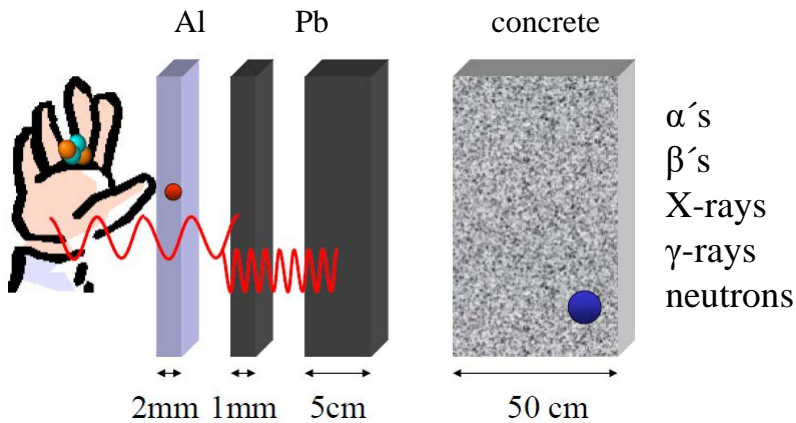
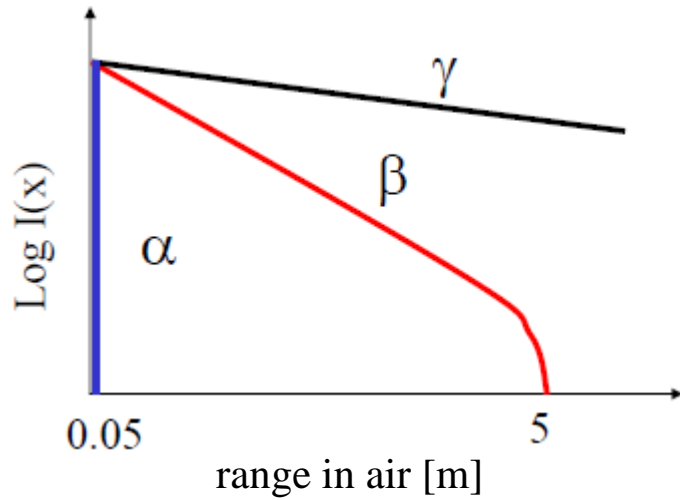
- We assume a time-independent Hamiltonian that generates stationary states.
- For radioactive decay we need to adopt a perturbative approach
- Our potential is of the form $V + V'$, where V' is a very weak addition to our potential that allows for transitions between the two states.

Transition probability is therefore just the decay constant λ , and is given by **Fermi's Golden Rule**:

$$\lambda = \frac{2\pi}{\hbar} |V'_{fi}|^2 \rho(E_f)$$

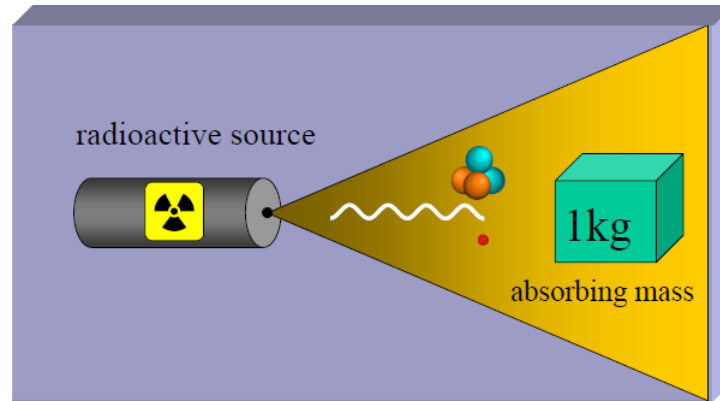
- The transition matrix element that connects the initial and final states $V'_{fi} = \langle f|V'|i\rangle$
- The density $\rho(E_f)$ of final states

Radiation protection



$$\Omega = \frac{a^2}{4\pi \cdot R^2} = \frac{\pi \cdot r^2}{4\pi \cdot R^2}$$

Radiation protection



$$\text{absorbed dose} = \frac{\text{absorbed energy}}{\text{mass}}$$

$$1 \frac{\text{Joule}}{\text{kg}} = 1 \text{ Gray (Gy)}$$

The biological dose, sometimes also known as the dose equivalent is expressed in units of Sieverts [Sv]. This dose reflects the fact that the biological damage caused by a particle depends not only on the total energy deposited but also on the rate of energy loss per unit distance traversed by the particle.

$$\text{equivalent dose [Sv]} = \text{energy dose} \cdot \text{quality factor Q}$$

radiation	quality factor Q
X-ray, γ , β	1
thermal neutrons	2.3
fast neutrons	10
α -particles, heavy ions	20