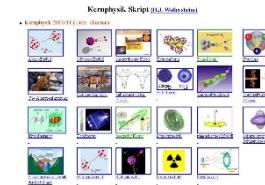


Outline: Fusion reactions

Lecturer: Hans-Jürgen Wollersheim

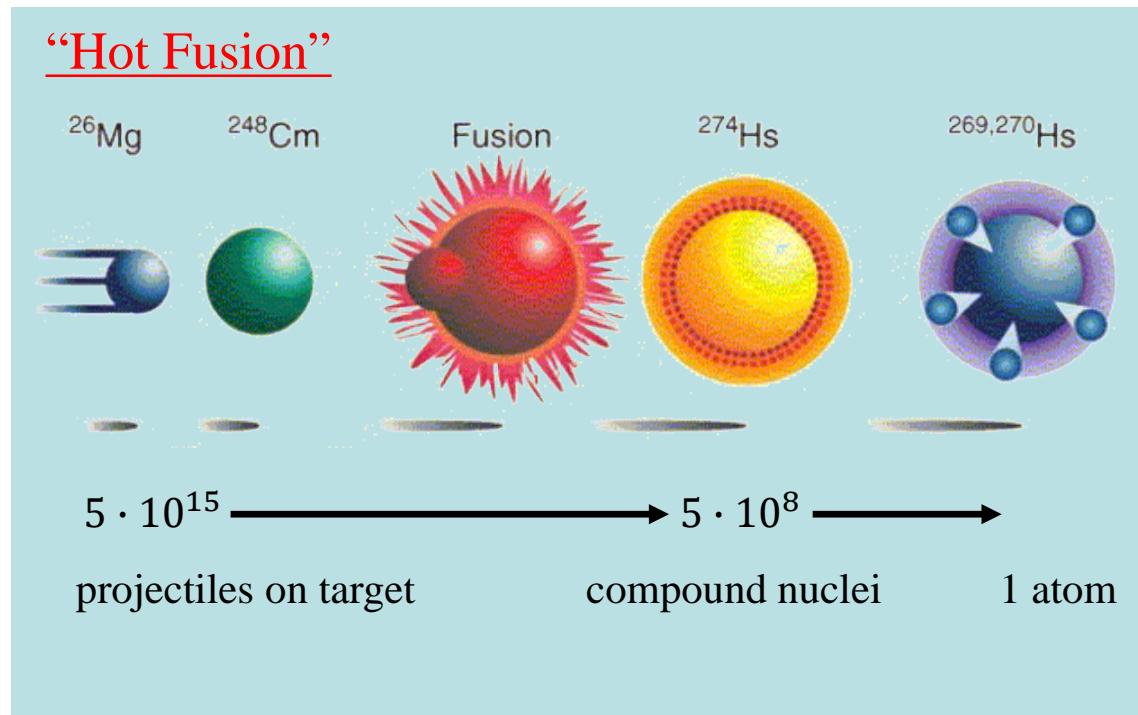
e-mail: h.j.wollersheim@gsi.de

web-page: <https://web-docs.gsi.de/~wolle/> and click on



1. fusion cross section
2. de-excitation of a hot compound nucleus
3. nuclear temperature, level density, evaporation particles
4. limiting nuclear angular momentum
5. production of super heavy elements SHE

Fusion reactions



Hot fusion (~ 1952) successful up to element 106 (Seaborgium)

- Coulomb barrier V_C between projectile and target nucleus has to be exceeded

$$V_C = \frac{Z_p \cdot Z_t \cdot e^2}{R_{int}} = 126.2 \text{ MeV} \quad ({}^{26}\text{Mg} + {}^{248}\text{Cm})$$

- reaction: $a + A \rightarrow C^* \rightarrow B + b$

$$\Delta m = m_a + m_A - m_{CN}$$



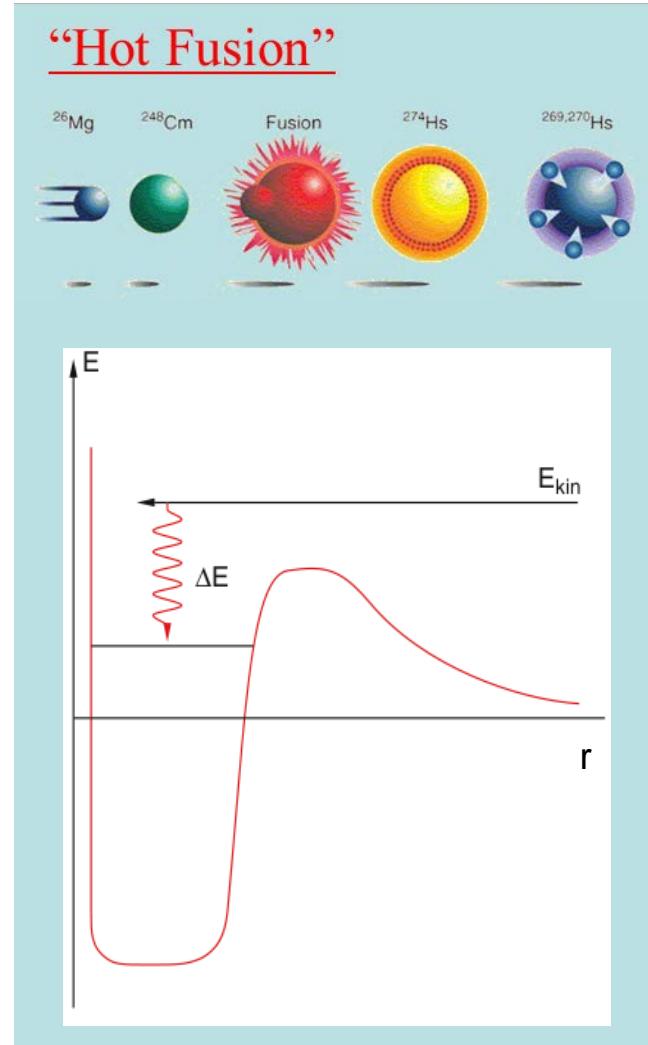
<https://www.nndc.bnl.gov/qcalc/>

$$\begin{aligned} \Delta m &= (25.983 + 248.072 - 274.143) * 931.478 \text{ MeV/c}^2 \\ &= -82.153 \text{ MeV/c}^2 \end{aligned}$$

- excitation energy of compound nucleus

$$\begin{aligned} E^* &= E_{kin} + \Delta m \cdot c^2 \\ &= 126.2 \text{ MeV} - 82.2 \text{ MeV} \\ &= \mathbf{44.0 \text{ MeV}} \end{aligned}$$

- approximate 4 neutrons will be evaporated to avoid fission



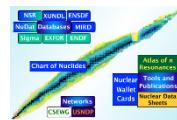
Cold fusion (1981-1996)

- Coulomb barrier V_C between projectile and target nucleus has to be exceeded

$$V_C = \frac{Z_p \cdot Z_t \cdot e^2}{R_{int}} = 223.3 \text{ MeV} \quad (^{58}\text{Fe} + ^{208}\text{Pb})$$

- reaction: $a + A \rightarrow C^* \rightarrow B + b$

$$\Delta m = m_a + m_A - m_{CN}$$



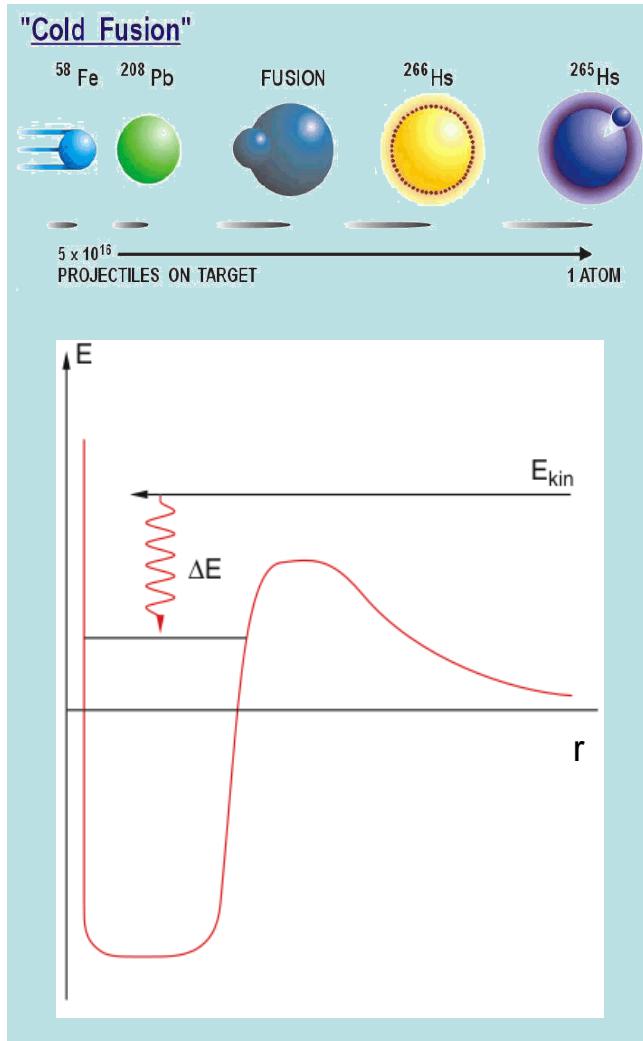
<https://www.nndc.bnl.gov/qcalc/>

$$\begin{aligned} \Delta m &= (57.933 + 207.977 - 266.130) * 931.478 \text{ MeV}/c^2 \\ &= -205.092 \text{ MeV}/c^2 \end{aligned}$$

- excitation energy of compound nucleus

$$\begin{aligned} E^* &= E_{kin} + \Delta m \cdot c^2 \\ &= 223.3 \text{ MeV} - 205.1 \text{ MeV} \\ &= \mathbf{18.2 \text{ MeV}} \end{aligned}$$

- approximate 1-2 neutrons will be evaporated to avoid fission

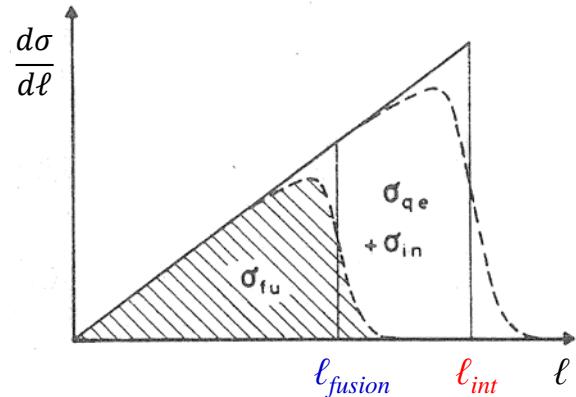


Fusion cross section

Radius for fusion barrier:

$$R_{fusion} = R_{int} - \begin{cases} 0.3117 \cdot (Z_p \cdot Z_t)^{0.2122} & Z_p \cdot Z_t < 500 \\ 1.096 + 1.391 \cdot Z_p \cdot Z_t / 1000 & Z_p \cdot Z_t \geq 500 \end{cases} [fm]$$

	R_i [fm]	C_i [fm]	R_{int} [fm]	$V_C(R_{int})$ [MeV]	R_{fusion} [fm]	$V_C(R_{fusion})$ [MeV]
^{58}Fe	4.40	4.17	13.75	223.3	12.36	248.4
^{208}Pb	6.96	6.82				



Total cross section for fusion:

$$\sigma_{fusion} = \pi R_{fusion}^2 \cdot \left[1 - \frac{V_C(R_{fusion})}{E_{cm}} \right] \quad \text{with } E_{cm} = \frac{A_t}{A_t + A_p} \cdot E_{lab}$$

$$\sigma_{fusion} = \frac{\pi}{k_\infty^2} \cdot \ell_{fusion} \cdot (\ell_{fusion} + 1) \quad \text{with } k_\infty = 0.2187 \cdot \frac{A_t}{A_t + A_p} \cdot \sqrt{A_p \cdot E_{lab}} [fm^{-1}]$$

Interaction potential

The potential between projectile and target nucleus is given by a function of the relative distance between them

$$V(r) = V_N(r) + V_C(r)$$

nuclear potential + Coulomb potential

$$V_C(r) = \begin{cases} \frac{Z_1 Z_2 e^2}{2 \cdot R_C} \left(3 - \frac{r^2}{R_C^2} \right) & r < R_C \\ \frac{Z_1 Z_2 e^2}{r} & r \geq R_C \end{cases}$$

nuclear proximity potential:

$$V_N(r) = 4\pi \cdot \gamma \cdot \frac{C_p \cdot C_t}{C_p + C_t} \cdot b \cdot \Phi(\xi)$$

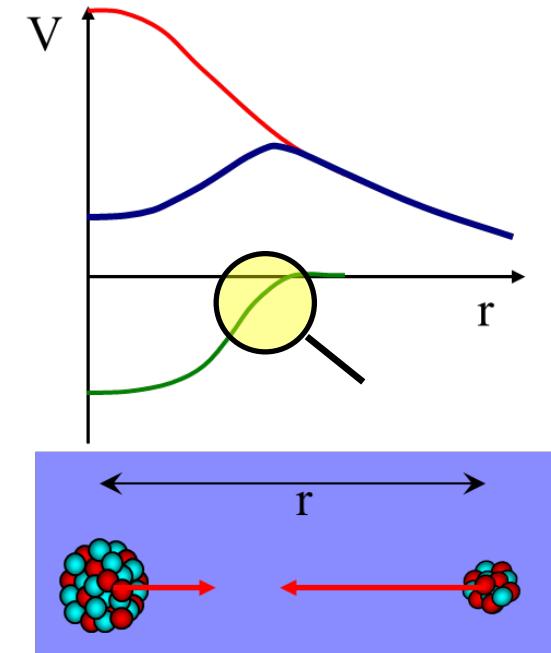
$$\Phi(\xi) = \begin{cases} -0.5 \cdot (\xi - 2.54)^2 - 0.0852 \cdot (\xi - 2.54)^3 & \xi \leq 1.2511 \\ -3.437 \cdot \exp(-\xi/0.75) & \xi \geq 1.2511 \end{cases}$$

$$\xi = (r - C_p - C_t)/b$$

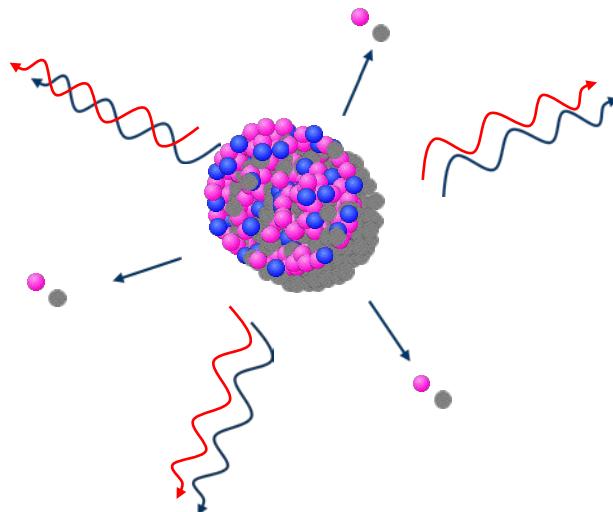
$$b = \frac{\pi}{\sqrt{3}} \cdot a \cong 1 \text{ fm} \quad \text{with } a = 0.55 \text{ fm}$$

$$\gamma = 0.9517 \cdot \left\{ 1 - 1.7826 \cdot \left(\frac{N_c - Z_c}{A_c} \right)^2 \right\} \quad \frac{\text{MeV}}{\text{fm}^2}$$

$$C_i = R_i \cdot (1 - R_i^{-2}) \quad [\text{fm}] \quad R_i = 1.28 \cdot A_i^{1/3} - 0.76 + 0.8 \cdot A_i^{-1/3} \quad [\text{fm}]$$



The statistical model de-excitation of the hot compound system



$$E^* = E_{kin} + \Delta m \cdot c^2$$

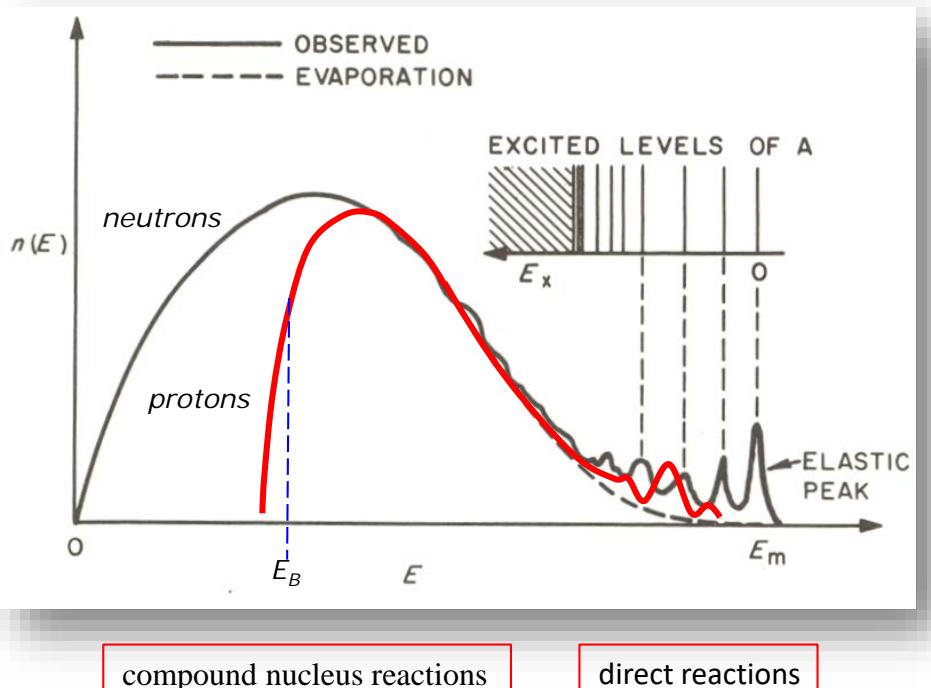
$$E_{kin} > V_C = \frac{Z_a \cdot Z_A \cdot e^2}{R_{int}}$$

$$\Delta m = m_a + m_A - m_{CN}$$



<https://www.nndc.bnl.gov/qcalc/>

Evaporation particles



cm-spectra of particles statistically emitted from CN (evaporation) are of Maxwell Boltzmann type

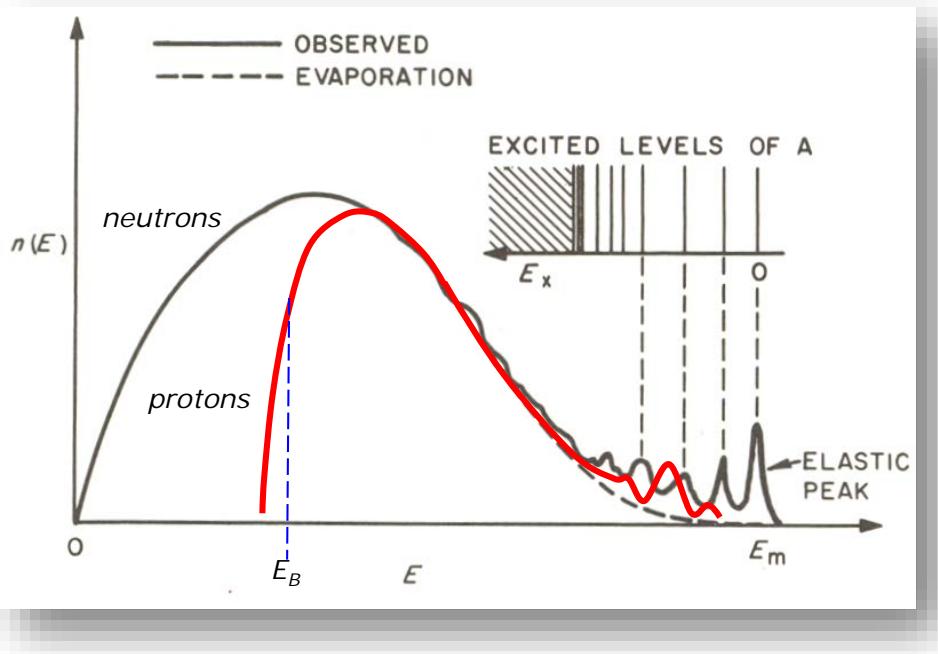
$$\frac{dN}{dE} \propto (E - E_B) \cdot e^{-E/T}$$

E_B = Coulomb barrier

T = effective nuclear temperature

Typical energy spectrum of nucleons emitted at a fixed angle in inelastic nucleon-nucleon reactions.

Evaporation particles

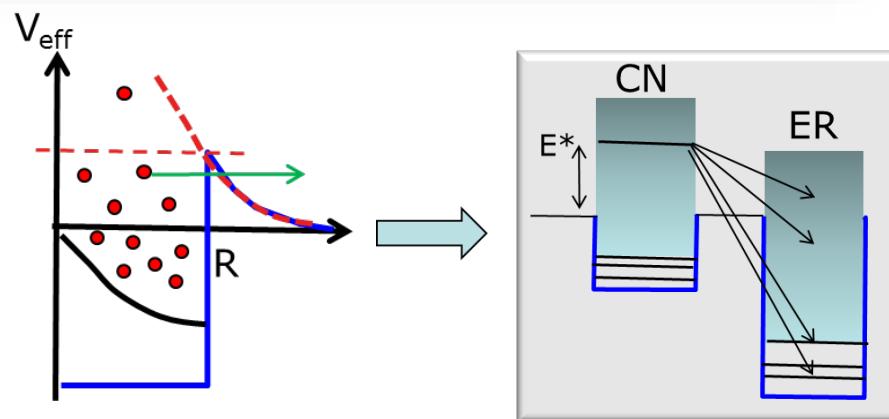


cm-spectra of particles statistically emitted from CN (evaporation) are of Maxwell Boltzmann type

$$\frac{dN}{dE} \propto (E - E_B) \cdot e^{-E/T}$$

E_B = Coulomb barrier

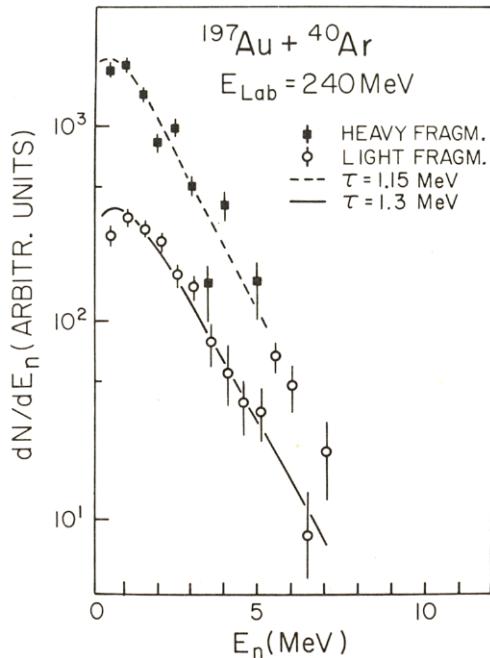
T = effective nuclear temperature



Even for fixed E^* the particle spectrum is continuous (Maxwell Boltzmann), except for transitions to discrete spectrum at low E_{ER}^*

Nuclear temperatures and level densities

de-excitation of the hot compound system



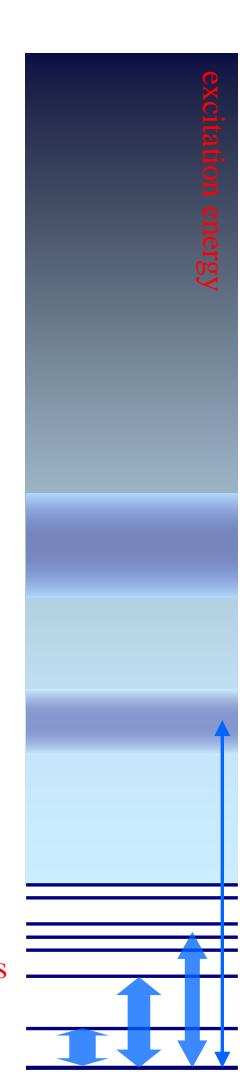
$$\frac{dN}{dE_n} \propto E_n \cdot e^{-E_n/T}$$

spectrum of single neutron

$$\langle E_n \rangle = 2T \quad \max \frac{dN}{dE_n} @ E_n = T$$

$$\frac{dN}{dE_n} \propto \sqrt{E_n} \cdot e^{-E_n/T_{\text{eff}}}$$

$$\langle E_n \rangle = 1.5T \quad T_{\text{eff}} \approx 0.92 \cdot T \quad (1^{\text{st}} \text{ daughter})$$

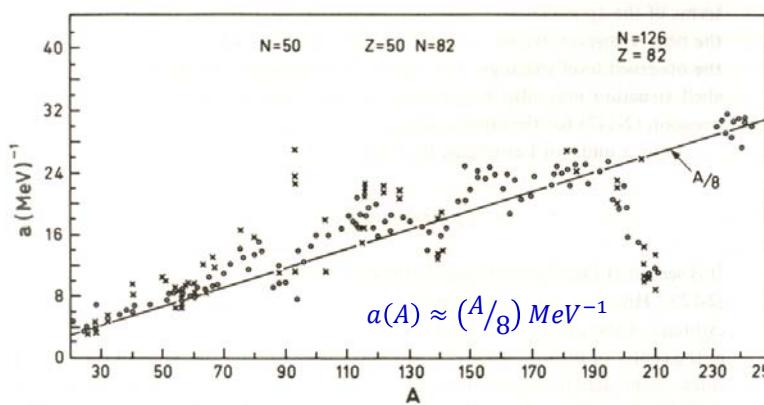


Fermi gas relations:

$$E^* = a \cdot T^2 \quad \text{"little - a"}$$

$$S = \int \frac{dE^*}{E^*} = 2\sqrt{a \cdot E^*}$$

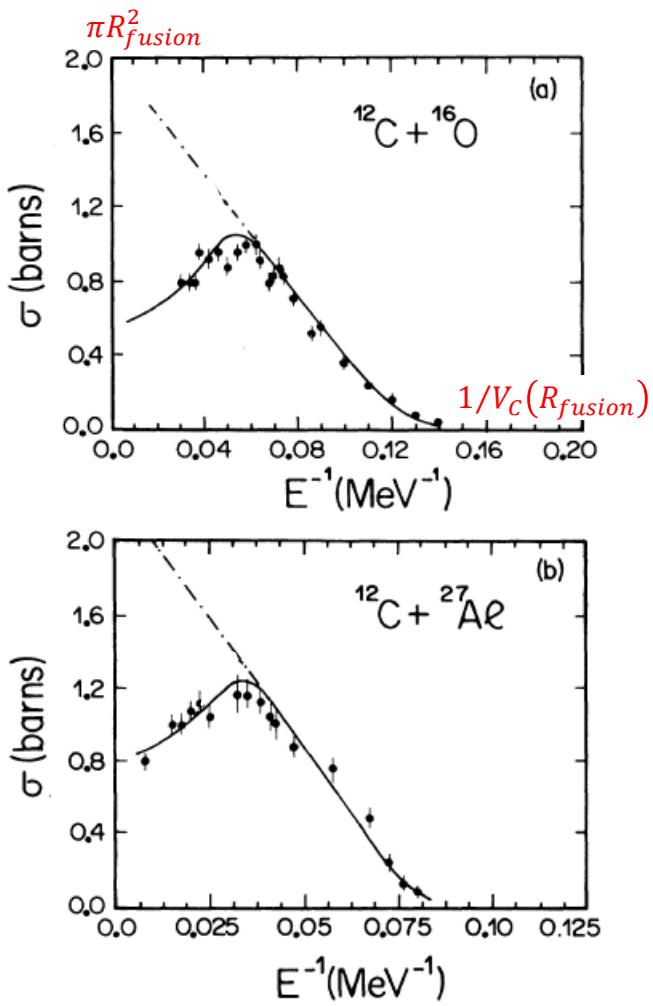
$$\rho(E^*) = \rho_0 \cdot e^{2\sqrt{a \cdot E^*}}$$



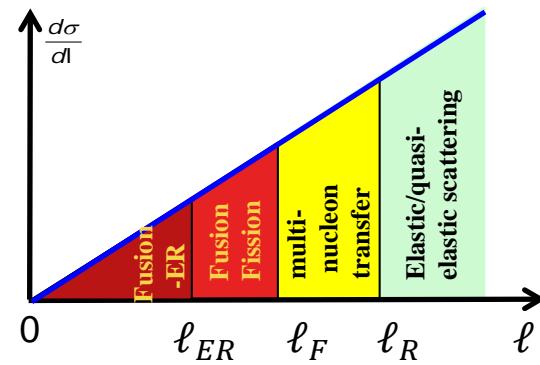
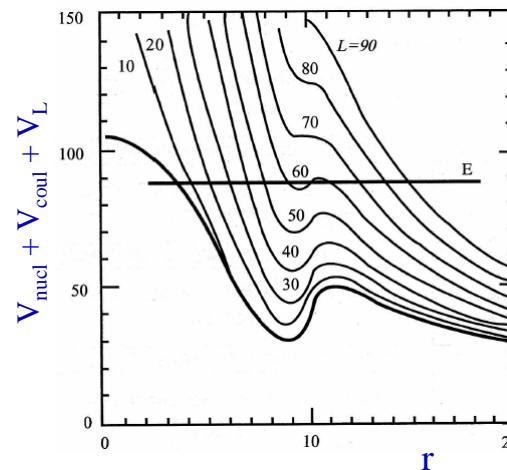
collective states

ground state

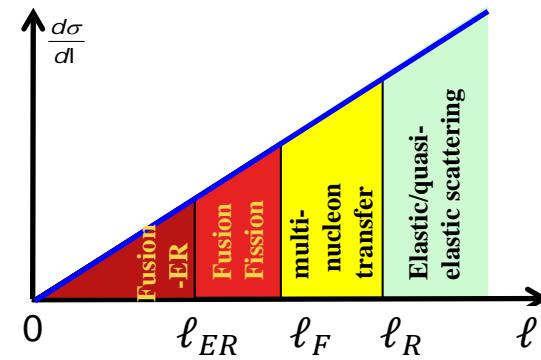
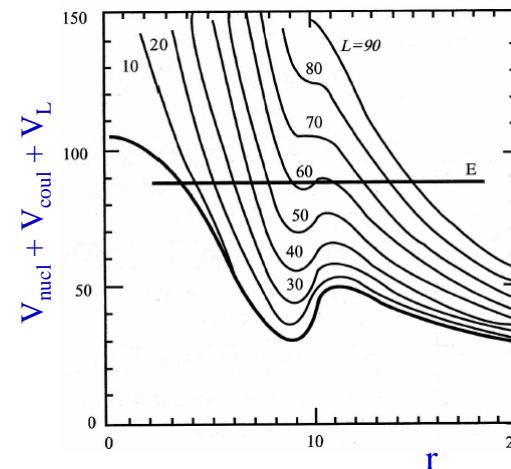
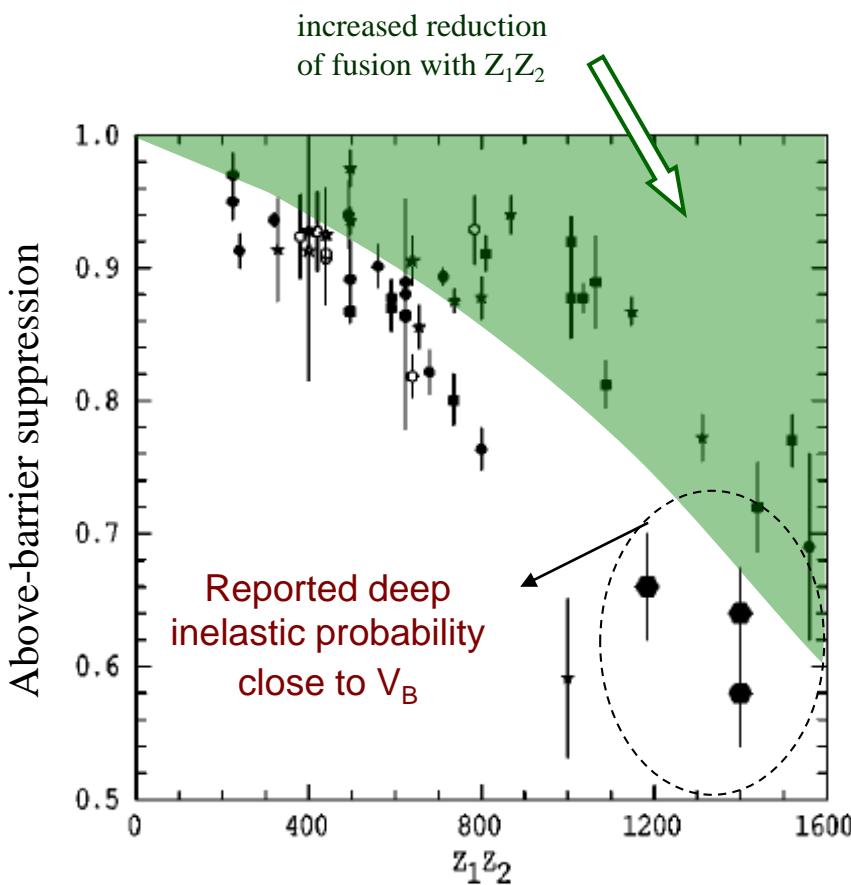
Fusion excitation functions



maximum ℓ_{fusion} due to nuclear centrifugal stability



Reduction in fusion at above barrier energies



A limiting nuclear angular momentum rotating charged liquid drop

surface energy:

$$E_S^{(0)} = 17.9439 \cdot \left[1 - 1.7826 \cdot \left(\frac{N-Z}{A} \right)^2 \right] \cdot A^{2/3} \quad [MeV]$$

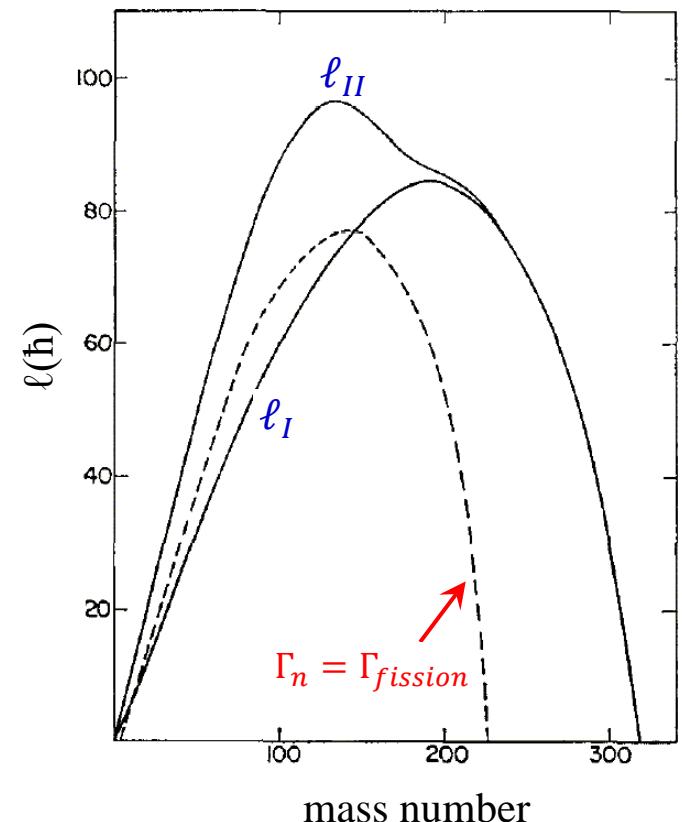
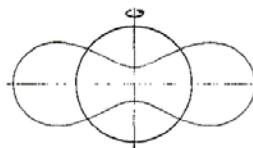
Coulomb energy:

$$E_{Coul}^{(0)} = 0.7053 \cdot (Z^2/A^{1/3}) \quad [MeV]$$

rotational energy:

$$E_{Rot}^{(0)} = \frac{1}{2} \frac{\hbar^2 \cdot \ell^2}{(2/5) \cdot A \cdot m \cdot R^2} = 34.54 \cdot \frac{\ell^2}{A^{5/3}} \quad [MeV]$$

change of the nuclear shape

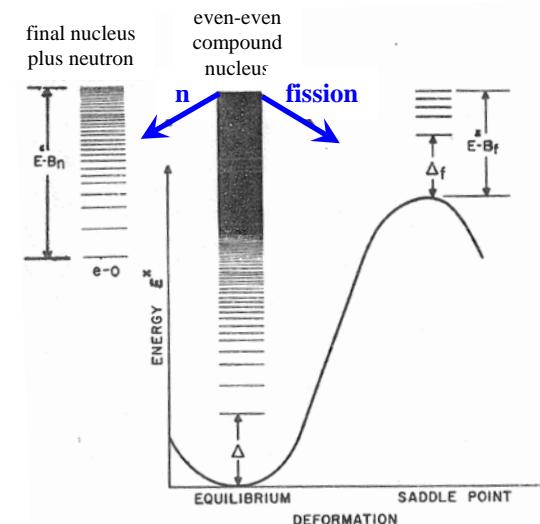
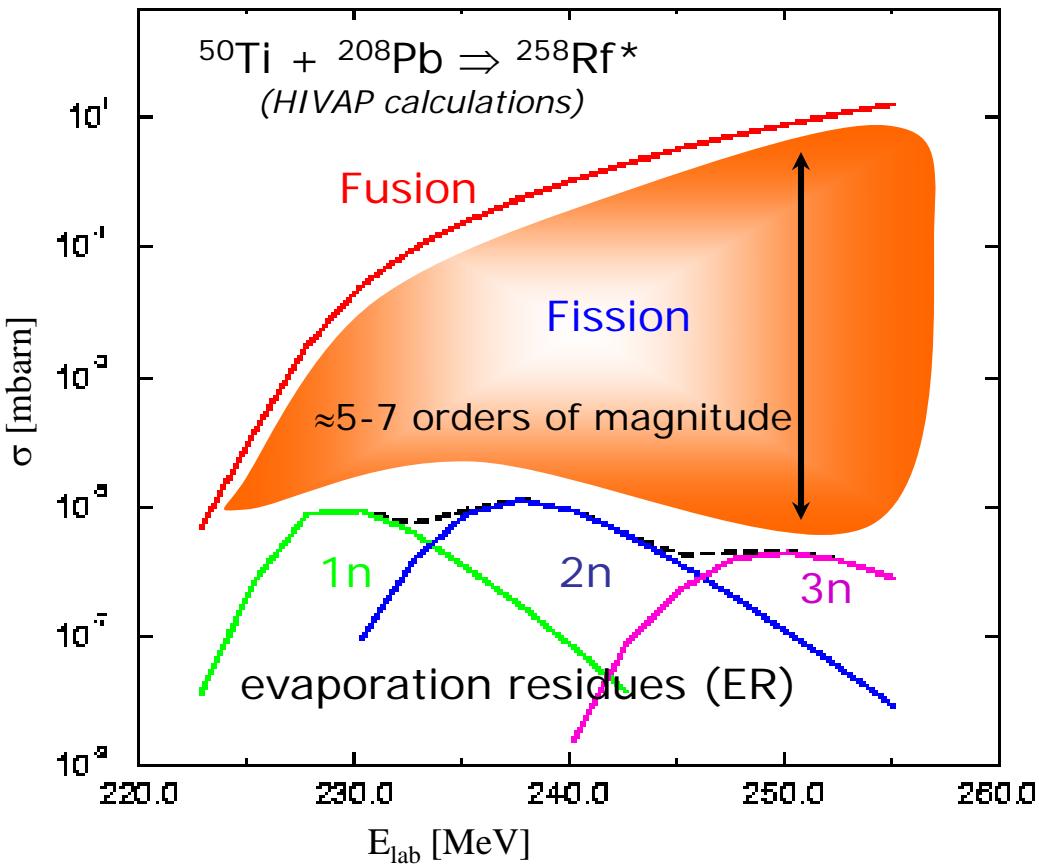


$$\frac{E_{Rot}^{(0)}}{E_S^{(0)}} = \begin{cases} 0.2829 - 0.3475 \cdot X - 0.0016 \cdot X^2 + 0.0501 \cdot X^3 & 0 \leq X \leq 0.75 \\ (7/5) \cdot (1-X)^2 - 4.5660 \cdot (1-X)^3 + 6.7443 \cdot (1-X)^4 & 0.75 \leq X \leq 1.0 \end{cases} \quad \text{with } X = \frac{E_{Coul}^{(0)}}{2 \cdot E_S^{(0)}} \text{ “fissility parameter”}$$

example: $^{127}_{57}La \quad E_S^{(0)} = 444.9 \text{ [MeV]} \quad E_{Coul}^{(0)} = 455.9 \text{ [MeV]} \quad X = 0.512 \quad E_{Rot}^{(0)}/E_S^{(0)} = 0.1112 \quad E_{Rot}^{(0)} = 49.48 \text{ [MeV]} \quad \ell_I = 67.8 \text{ [\hbar]}$

Cohen, Plasil, Swiatecki; Ann. Phys. 82, 557 (1974)

Fusion and evaporation cold fusion



Both decay processes are determined by the level density, either from the residual nucleus or at the saddle point.

level density: $\rho(E^*) = \text{const} \cdot \exp(E^*/T)$

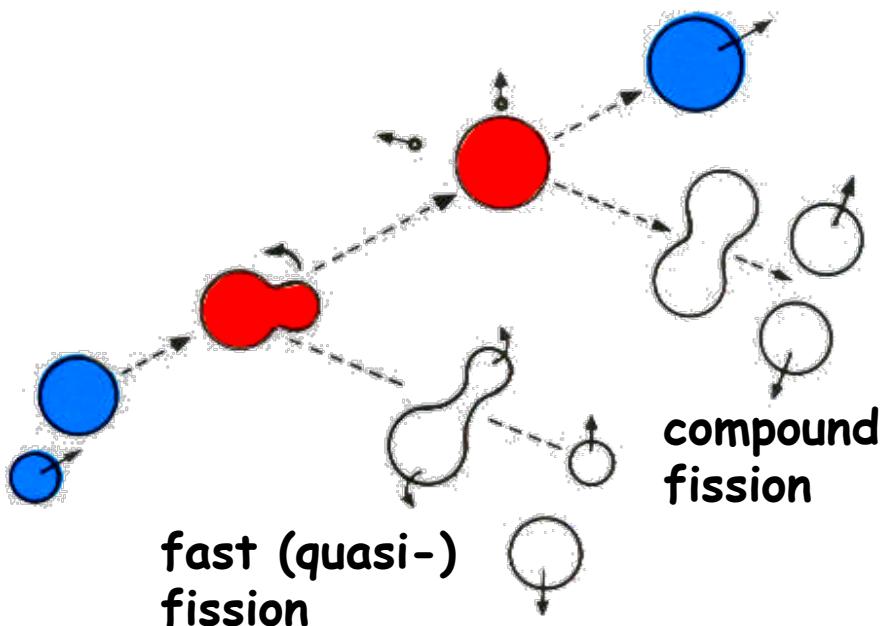
$$\frac{\Gamma_n}{\Gamma_f} = \frac{2 \cdot T \cdot A_{CN}^{2/3}}{K_0} \cdot \exp[(B_f - B_n)/T]$$

$$K_0 = \hbar^2 / 2 \cdot m \cdot r_0^2 \approx 11.4 \text{ MeV}$$

$$T = \sqrt{8 \cdot E^* / A_{CN}}$$

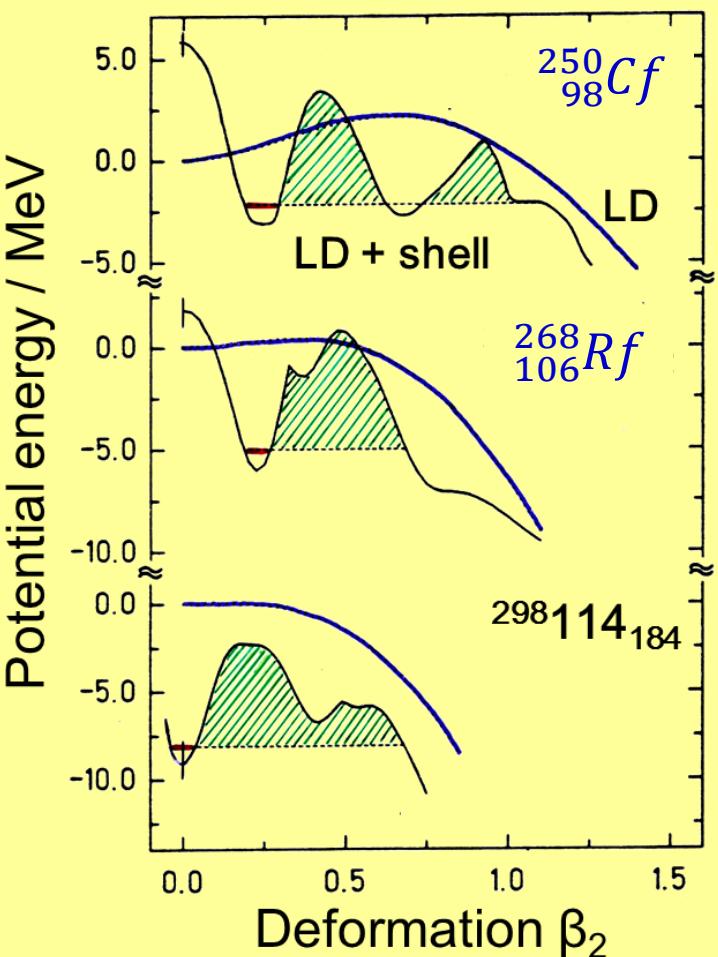
Fusion/Fission competition for SHE liquid drop + shell corrections

evaporation residue survival



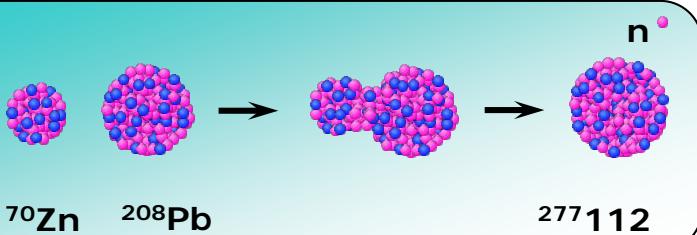
compound
fission

fast (quasi-)
fission



A. Sobiczewski et al.

Synthesis of heavy elements

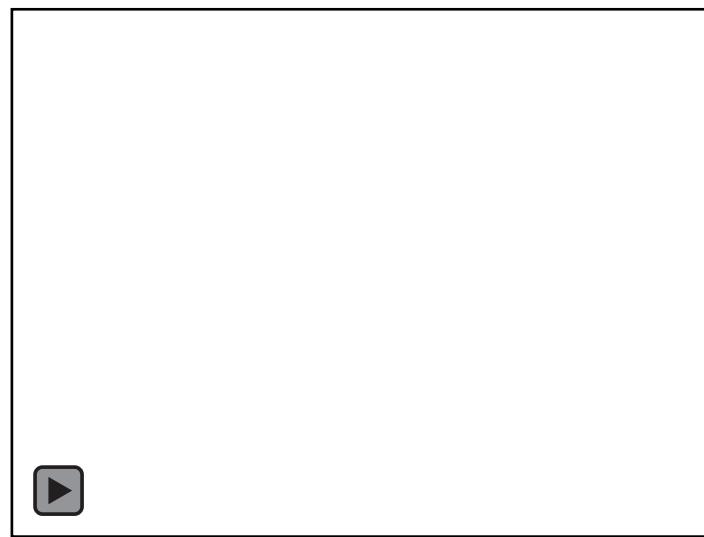


Fusion



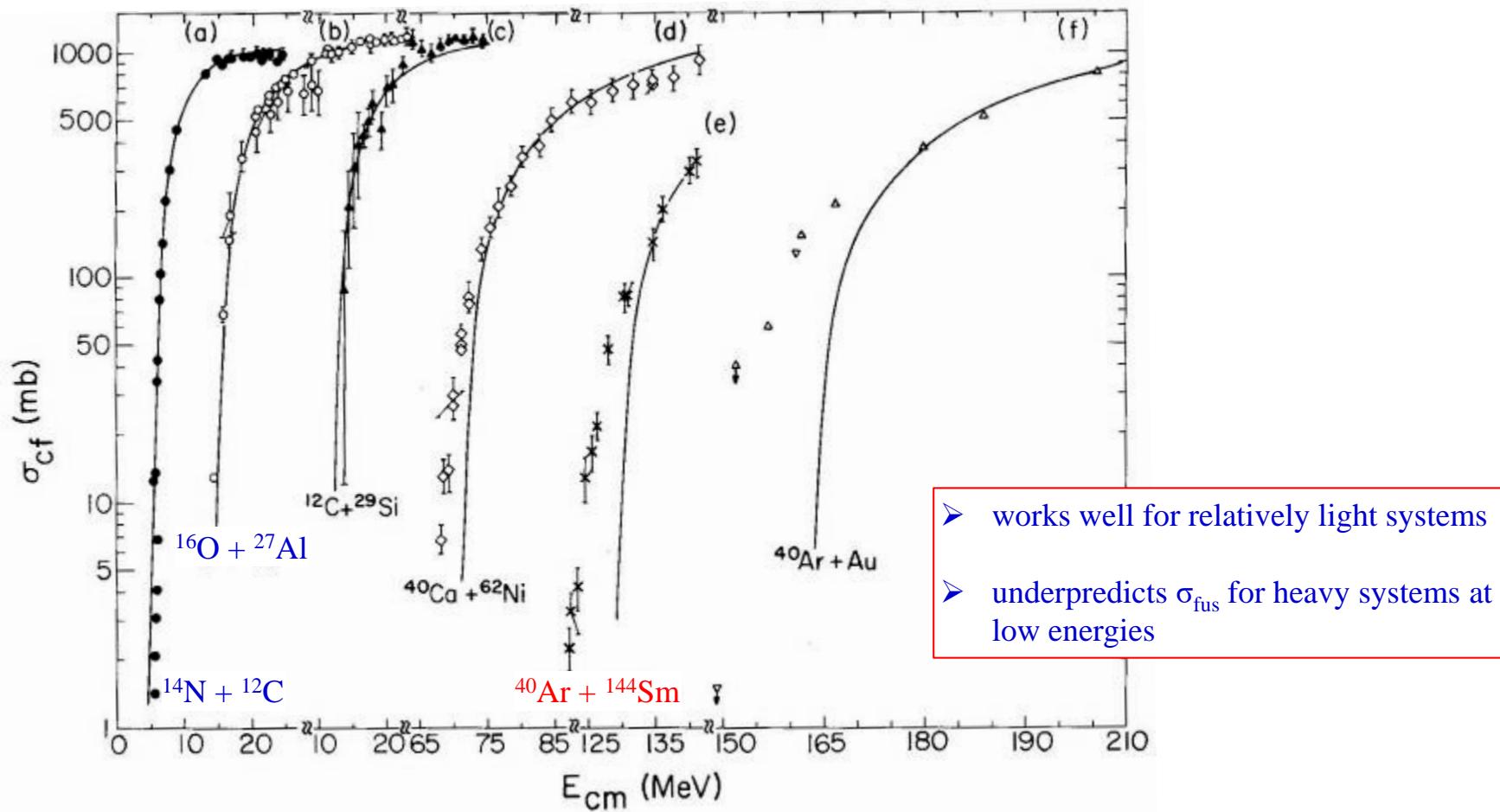
$\frac{1}{10^{12}}$

Compound Fission



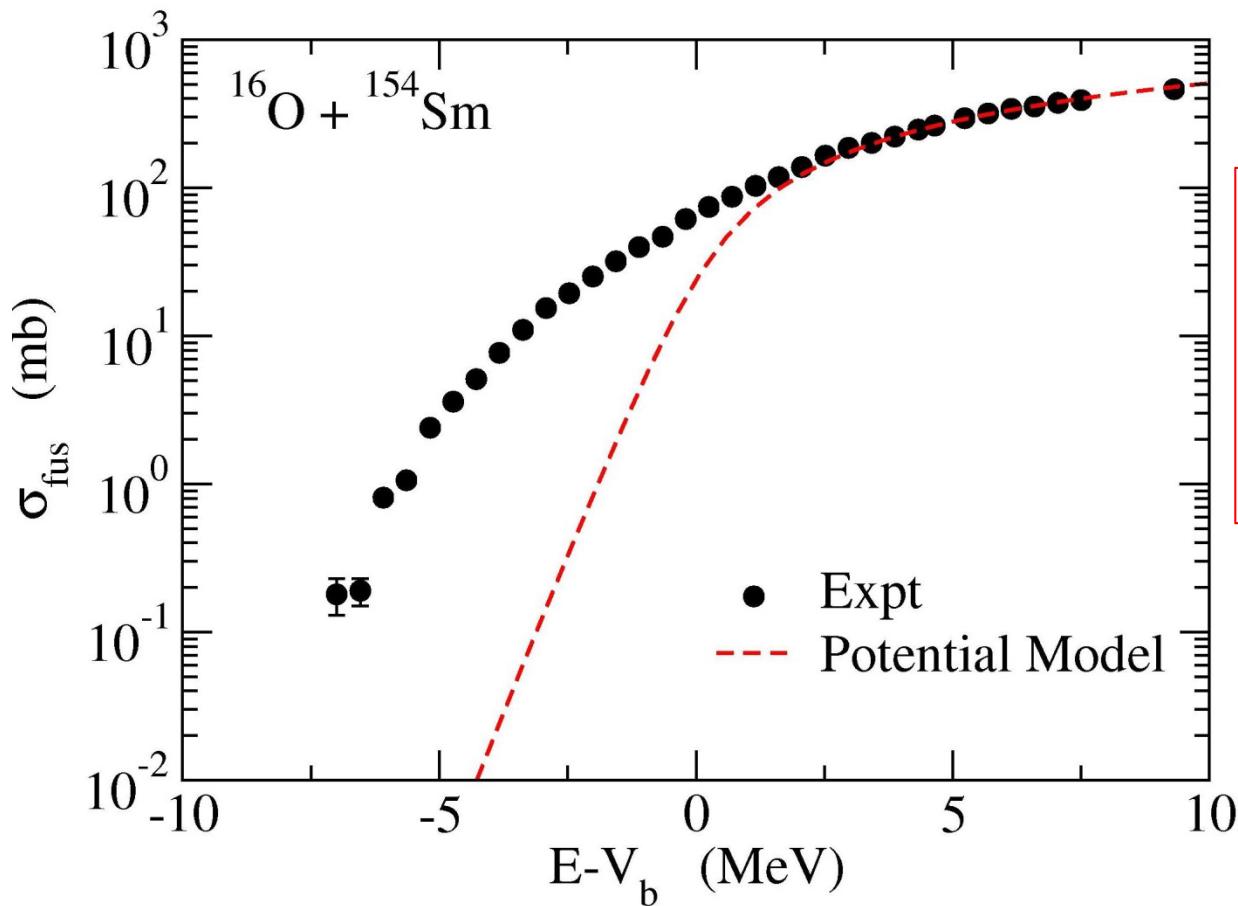
Comparison between potential model and experimental data

Fusion cross sections calculated with a static, energy independent potential



Comparison between potential model and exp. data

Fusion cross sections calculated with a static, energy independent potential



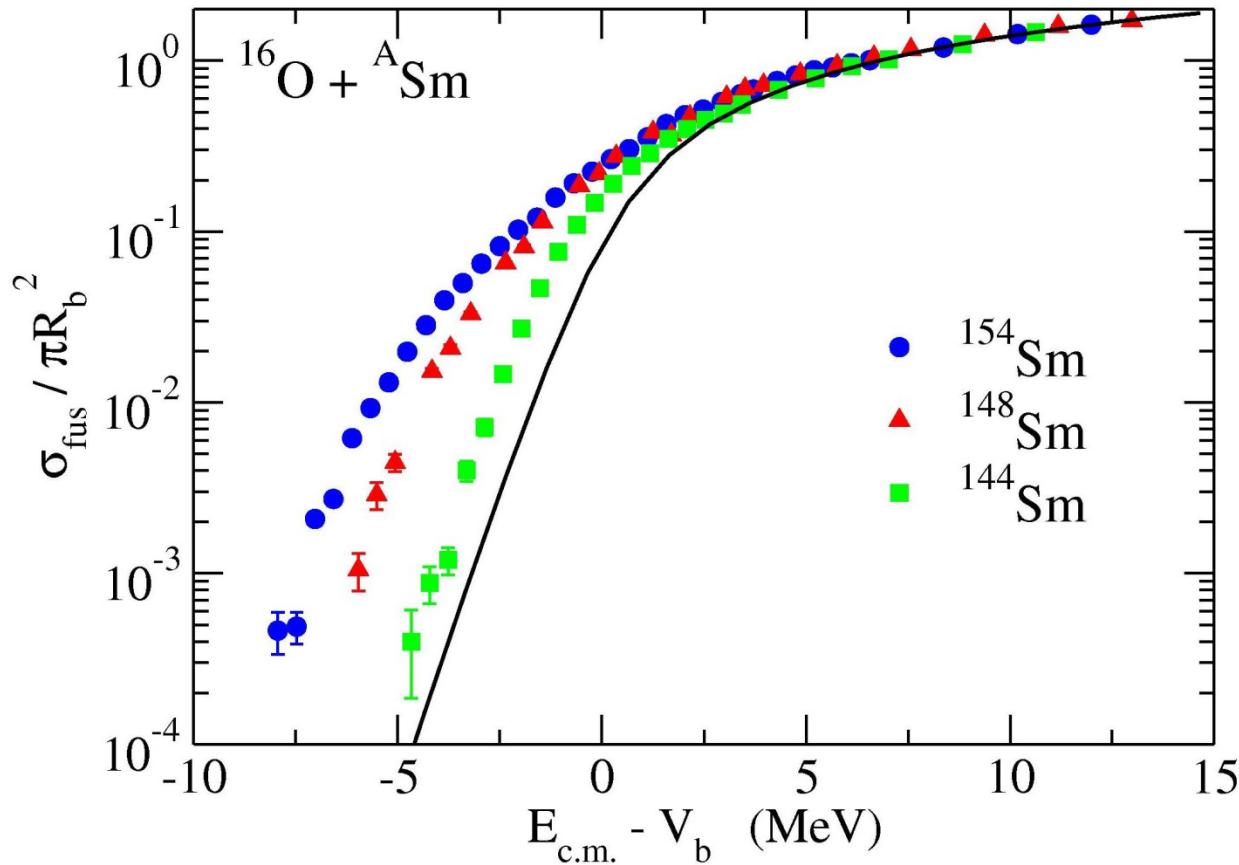
Potential model:

reproduces the data reasonably well
for $E > V_B$

underpredicts σ_{fus}
for $E < V_B$

Comparison between potential model and exp. data

Fusion cross sections calculated with a static, energy independent potential



Effect of collective excitation

$$\beta_2 \cong \frac{4\pi}{3 \cdot Z_t \cdot R_t^2} \cdot \sqrt{\frac{B(E2 \uparrow)}{e^2}}$$

$$\beta_4 \cong \frac{4\pi}{3 \cdot Z_t \cdot R_t^4} \cdot \sqrt{\frac{B(E4 \uparrow)}{e^2}}$$

(MeV)
1.81 — 3⁻
1.66 — 2⁺

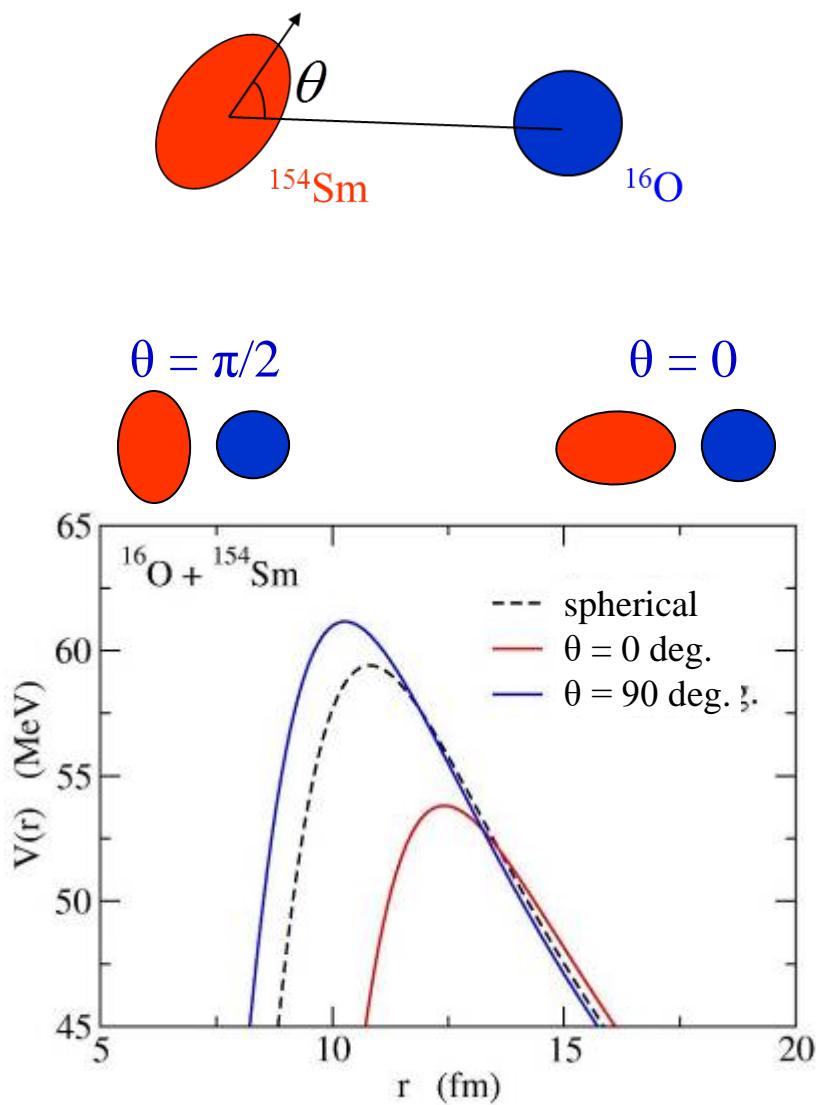
(MeV)
1.18 — 4⁺
1.16 — 3⁻

(MeV)
0.55 — 2⁺
0 — 0⁺

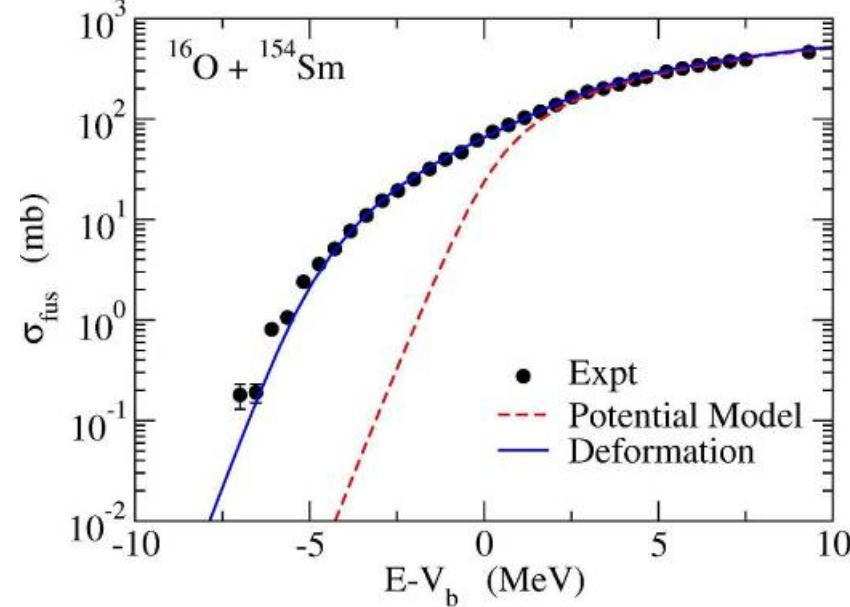
(MeV)
0.90 — 8⁺
0.54 — 6⁺

(MeV)
0.27 — 4⁺
0.082 — 2⁺
0 — 0⁺

Fusion for a deformed nucleus



$$\sigma_{fus}(E) = \int_0^1 \sigma_{fus}(E, \theta) d(\cos\theta)$$



- ❖ The barrier is lower for $\theta = 0$
- ❖ The barrier is higher for $\theta = \pi/2$

Deformation enhances σ_{fus} by a factor of 10 - 100

Fusion below and above the barrier inconsistent

■ $^{16}\text{O} + ^{208}\text{Pb}$

● $^{16}\text{O} + ^{204}\text{Pb}$

— a = 0.66 fm

— a = 1.18 fm

— a = 1.65 fm

$\sigma(\text{mb})$

— a = 0.66 fm

— a = 1.18 fm

— a = 1.65 fm

$E_{\text{c.m.}} - V_B$ (MeV)

$E_{\text{c.m.}} - V_B$ (MeV)