

# Gamma-ray spectroscopy II

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Lectures presented at the  
IoP Nuclear Physics Summer School  
September 4 – 17, 2005  
Chester, UK

# Outline

## First lecture

- Properties of  $\gamma$ -ray transitions
- Fusion-evaporation reactions
- Germanium detector arrays
- Coincidence technique
- Nuclear deformations
- Rotation of deformed nuclei
- Pair alignment
- Superdeformed nuclei
- Hyperdeformed nuclei
- Triaxiality and wobbling

## Second lecture

- Angular distribution
- Linear polarization
- Jacobi shape transition
- Charged-particle detectors
- Neutron detectors
- Prompt proton decay
- Recoil-decay tagging
- Rotation and deformation alignment

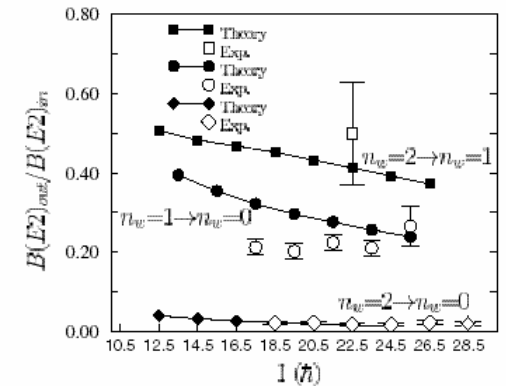
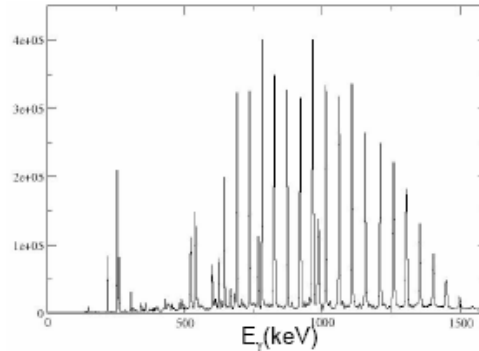
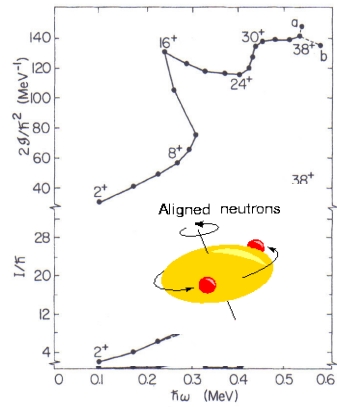
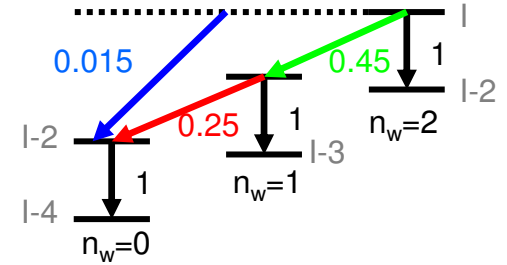
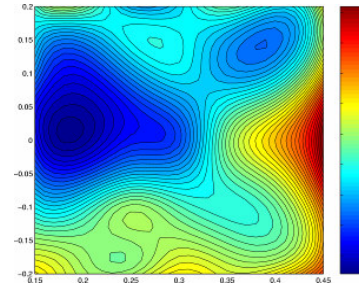
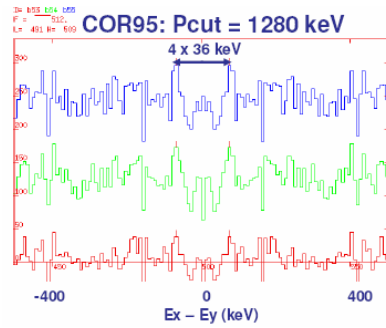
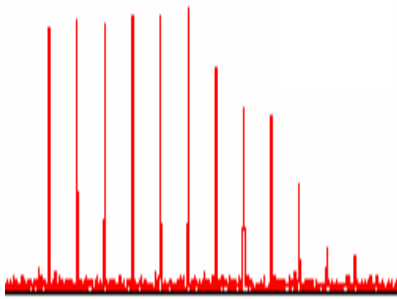
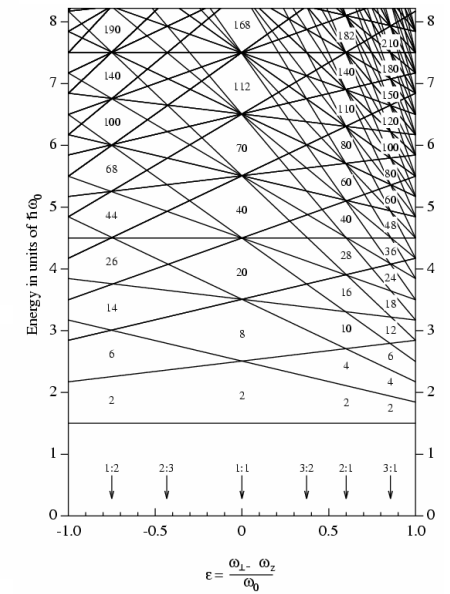
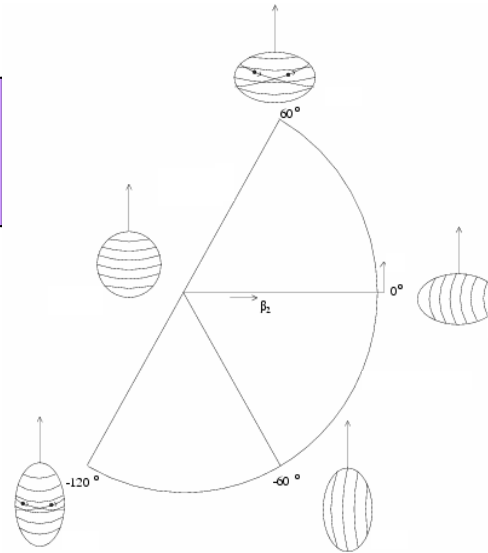
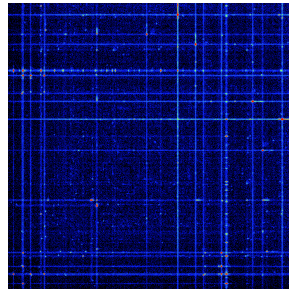
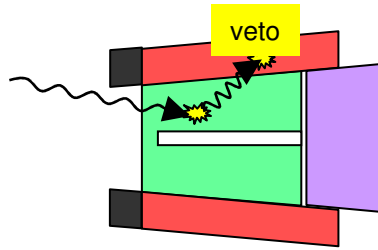
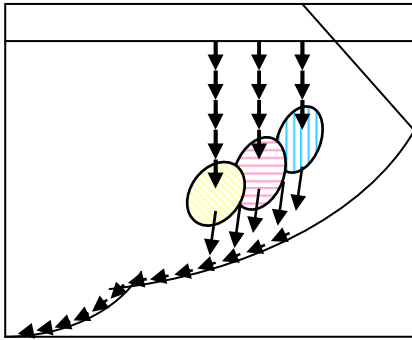
## Third lecture

- Spectroscopy of transfermium nuclei
- Conversion-electron spectroscopy
- Quadrupole moments and transition rates
- Recoil-distance method
- Doppler shift attenuation method
- Fractional Doppler shift method
- Magnetic moments
- Perturbed angular distribution
- Magnetic Rotation
- Shears Effect

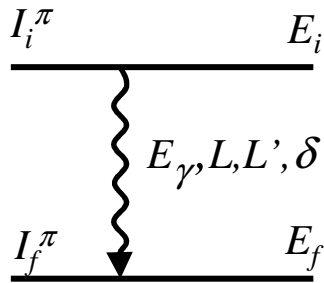
## Fourth lecture

- Fast fragmentation beams
- Isomer spectroscopy after fragmentation
- E0 transitions
- Shape coexistence
- Two-level mixing
- Coulomb excitation
- Reorientation effect
- ISOL technique
- Low-energy Coulomb excitation of  $^{74}\text{Kr}$
- Relativistic Coulomb excitation of  $^{58}\text{Cr}$
- Gamma-ray tracking
- AGATA

# Summary (I)

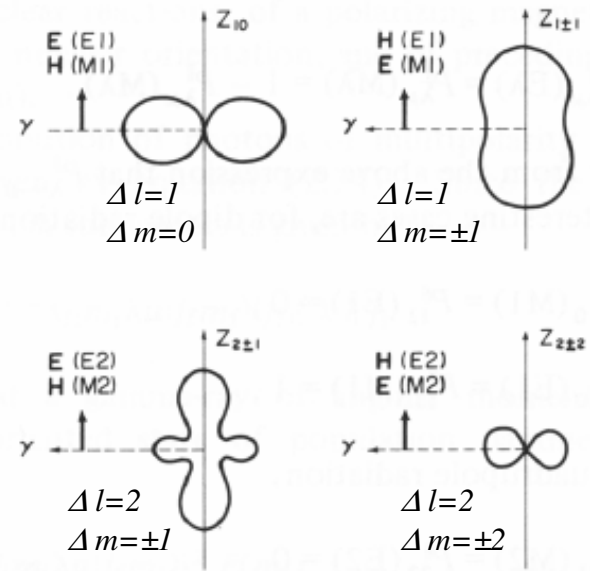
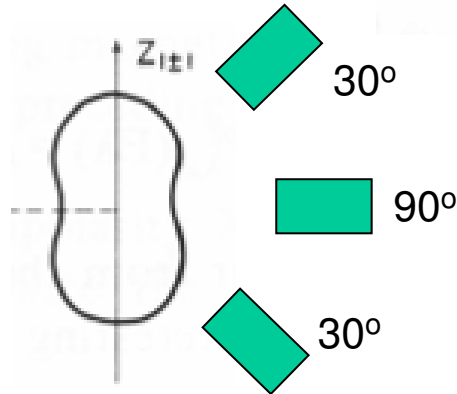
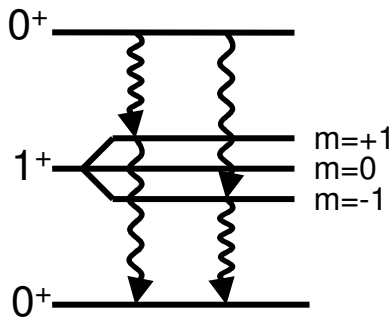


# Angular correlation



$$W(\vartheta) = 1 + \sum_k A_k P_k(\cos \vartheta)$$

simple example:

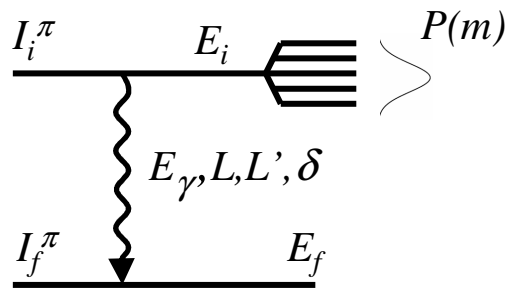


Most transitions following fusion-evaporation reactions have stretched dipole or quadrupole character.

⇒ compare experimental  $R_{DCO}$  with values for dipole-dipole, dipole-quadrupole, quadrupole-quadrupole cascades

⇒ often sufficient for spin assignments

# Angular distribution



Alignment of angular momentum after fusion-evaporation reaction:

$$P(m) = \frac{\exp\left(\frac{-m^2}{2\sigma^2}\right)}{\sum_{m'=-I}^{+I} \exp\left(\frac{-m'^2}{2\sigma^2}\right)}$$

$$W(\vartheta) = 1 + \sum_k A_k P_k(\cos \vartheta)$$

- The coefficients  $A_k$  depend on
- the multipolarity  $L$
  - the mixing parameter  $\delta$
  - the population width  $\sigma$

$$A_k(L, L', I_f, I_i) = \rho_k(I_i) \frac{1}{1 + \delta^2} \left[ F_k(L, L, I_f, I_i) + 2\delta F_k(L, L', I_f, I_i) + 2\delta^2 F_k(L', L', I_f, I_i) \right]$$

$$\rho_k(I_i) = \sqrt{2I_i + 1} \sum_{m=-I}^{+I} (-1)^{I_i - m} \langle I_i, m, I_i - m | k, 0 \rangle P(m)$$

Ferentz-Rosenzweig coefficients

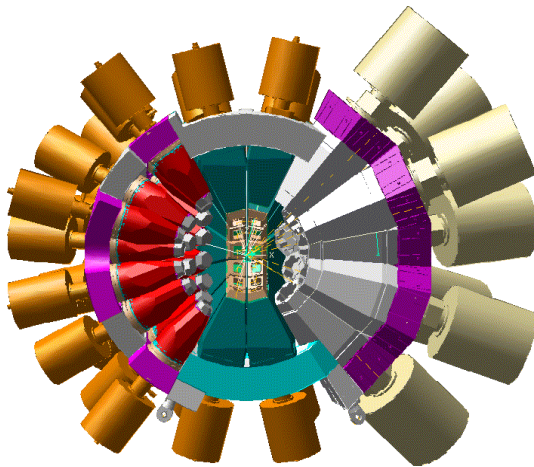
$$F_k(L, L', I_f, I_i) = (-1)^{I_f + I_i - 1} \sqrt{(2L + 1)(2L' + 1)(2I_i + 1)(2k + 1)} \begin{pmatrix} L & L' & k \\ 1 & -1 & 0 \end{pmatrix} \left\{ \begin{matrix} L & L' & k \\ I_i & I_i & I_f \end{matrix} \right\}$$

Clebsch-Gordan Racah

$\sigma/I$  is approximately constant (for a given reaction).

Normalize to transition with known multipolarity, e.g.  $2^+ \rightarrow 0^+$

# Example: Angular distribution with EUROBALL

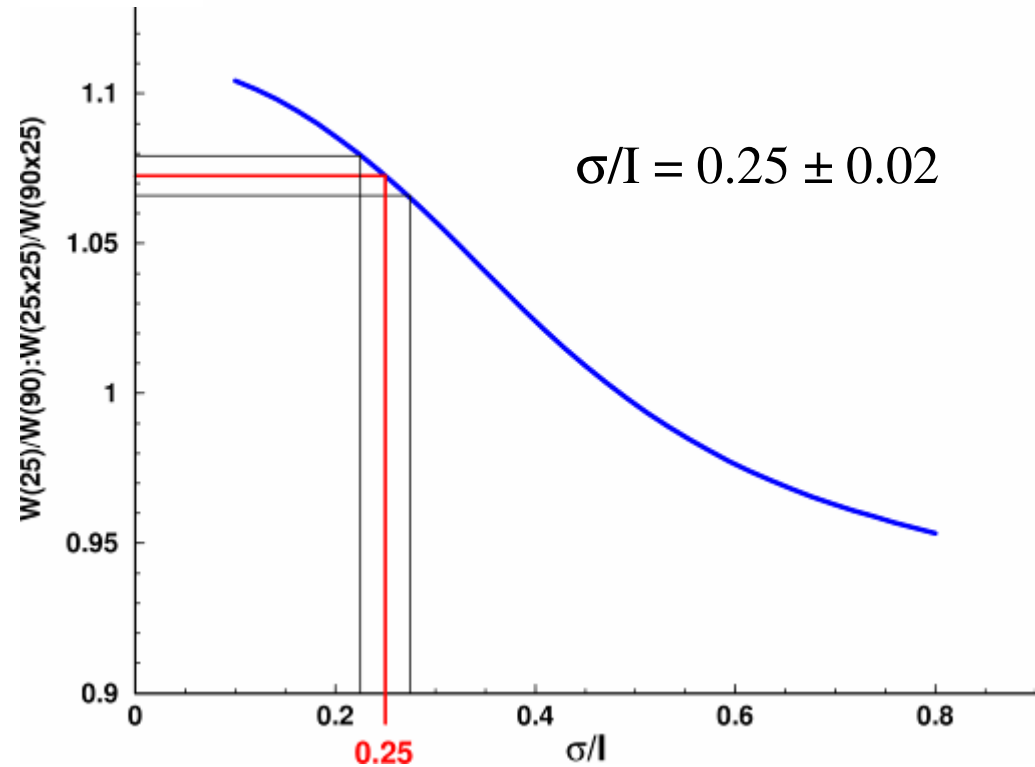
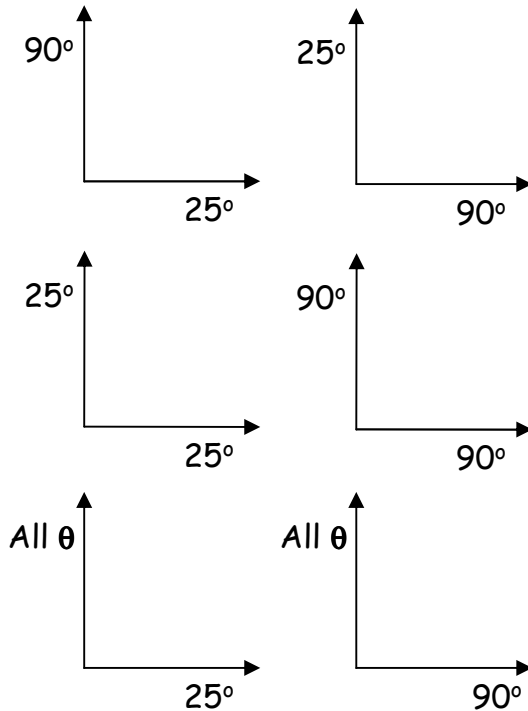


25° angle: ring1 + ring2 (tapered)  
 ring 7 (clusters)  
 90° angle: ring4 + ring5 (clovers)

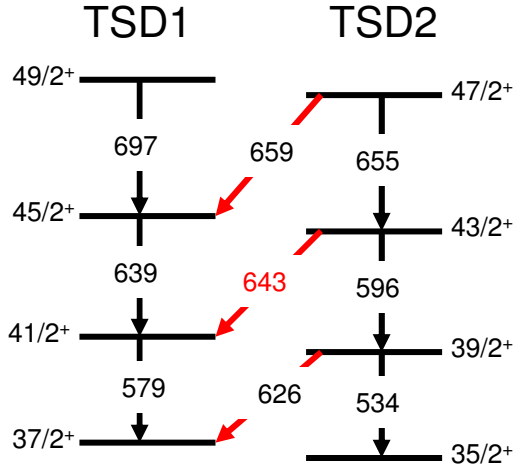
$$\frac{W(25)}{W(90)} \cdot \frac{W(25 \times 25)}{W(90 \times 25)}$$

$^{139}\text{La}(^{29}\text{Si},5n)^{163}\text{Lu}$  with  $E_{\text{beam}} = 153 \text{ MeV}$

Average of 8 stretched E2 transitions in TSD1 and TSD2

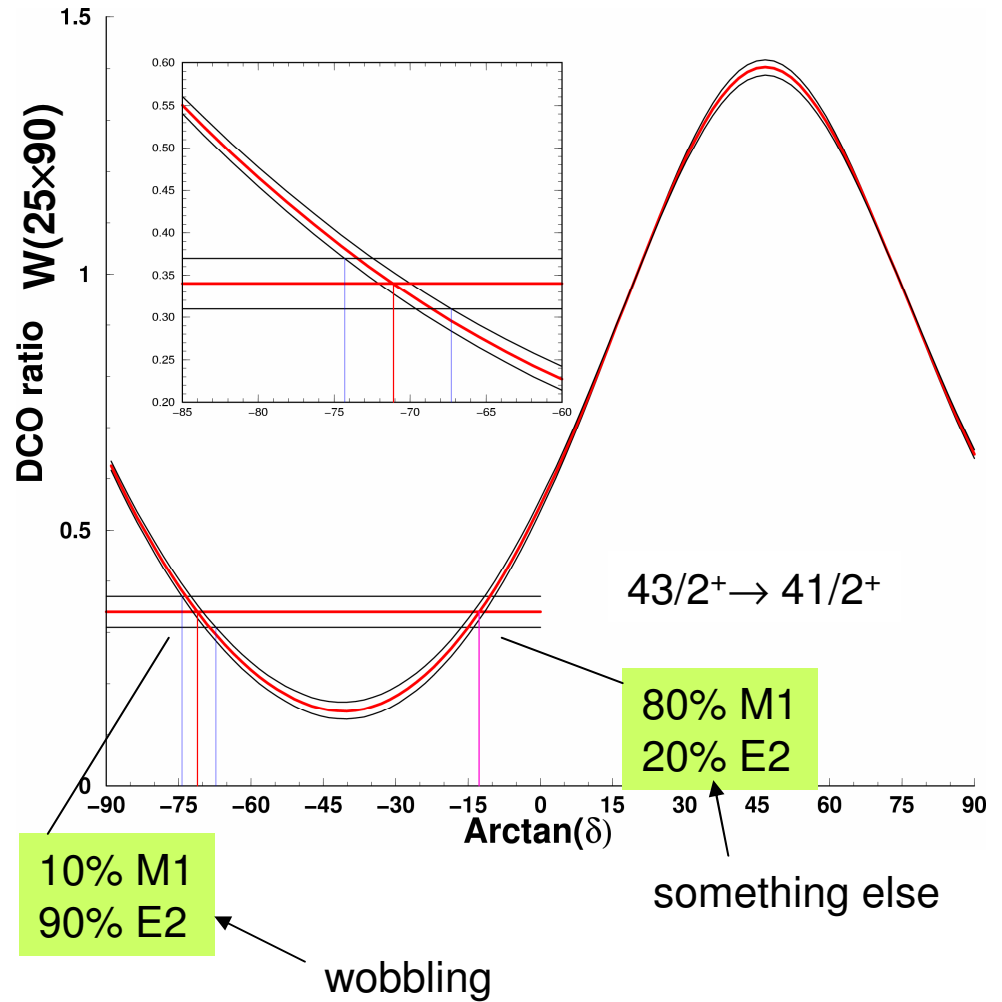


# Measuring the mixing parameter $\delta$



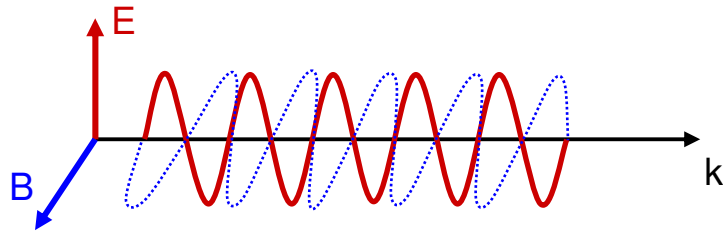
We know  $\sigma/I$  and have assigned  $I^\pi$   
 For wobbling bands, we expect  $\Delta I=1$  E2 inter-band transitions.  
 $\Rightarrow L=1, L'=2$ , large  $\delta$

Two possible solutions

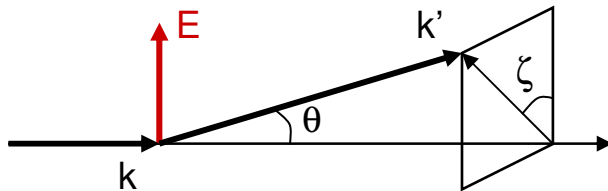


Angular distribution cannot distinguish between the two.  
 $\Rightarrow$  measure the linear polarization to establish electric or magnetic character.

# Linear polarization



linear polarization: fixed direction of electric field vector E



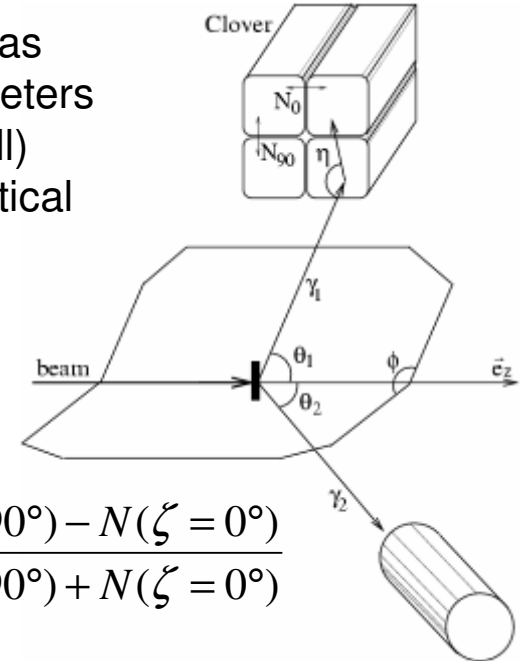
Compton scattering is sensitive to linear polarization: Klein-Nishina formula

$$\frac{d\sigma}{d\Omega} = \frac{r_0^2}{2} \frac{\omega'}{\omega^2} \left( \frac{\omega'}{\omega} + \frac{\omega}{\omega'} - 2 \sin^2 \theta \cos^2 \zeta \right)$$

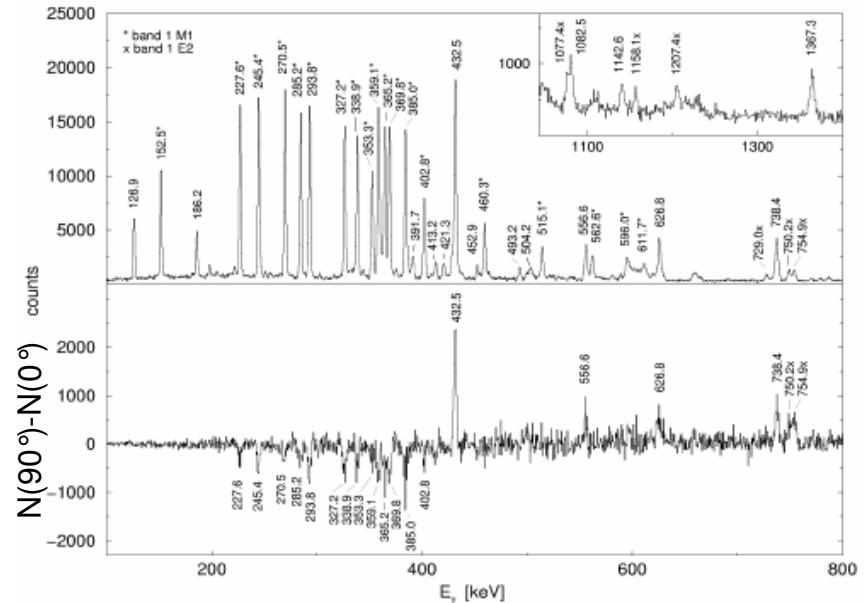
Effect is largest at  $\theta=90^\circ$

electric transitions appear positive,  
magnetic transitions negative

Clover detectors as Compton polarimeters (at  $90^\circ$  in Euroball) horizontal vs. vertical scattering

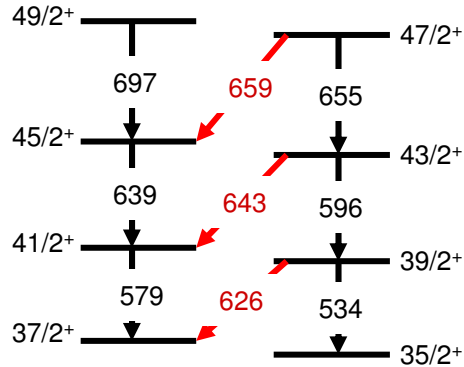


$$P = \frac{A}{Q} = \frac{1}{Q} \frac{N(\zeta = 90^\circ) - N(\zeta = 0^\circ)}{N(\zeta = 90^\circ) + N(\zeta = 0^\circ)}$$

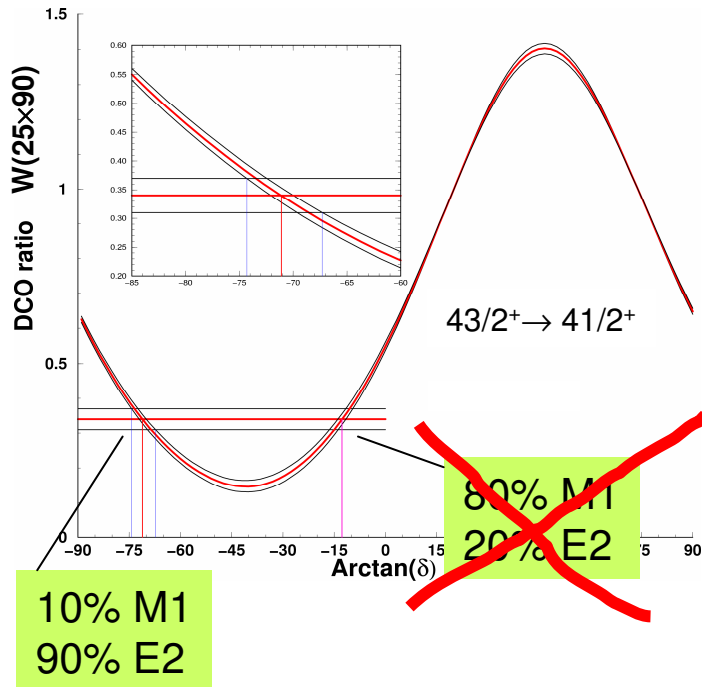




# Polarization measurement in $^{163}\text{Lu}$



	$E_\gamma$	$A = \frac{N(90^\circ) - N(0^\circ)}{N(90^\circ) + N(0^\circ)}$	
E2	579	$0.10 \pm 0.03$	positive
	697	$0.13 \pm 0.03$	
	386	$0.06 \pm 0.05$	
	534	$0.05 \pm 0.04$	
M1	349	$-0.11 \pm 0.05$	negative
inter-band	607	$0.05 \pm 0.05$	positive ⇒ electric
	626	$0.12 \pm 0.05$	
	643	$0.11 \pm 0.05$	
	659	$0.17 \pm 0.09$	
	673	$0.18 \pm 0.09$	



Confirmation of the wobbling mode in  $^{163}\text{Lu}$  through combined angular distribution and linear polarization measurement.

S.W. Ødegård et al., Phys. Rev. Lett. 86, 5866 (2001)

## MacLaurin shapes

What happens if we spin a liquid drop ?

It becomes oblate !

Jupiter:

- $T = 9 \text{ h } 50 \text{ min}$
- polar / equatorial axis  $\sim 15/16$

MacLaurin shape after C. MacLaurin (1698-1746)

But what if we spin really fast ?



## Jacobi shapes

The equilibrium shape changes abruptly to a very elongated triaxial shape rotating about its shortest axis.



piece of moon rock from Apollo mission

# The Jacobi shape transition in nuclei



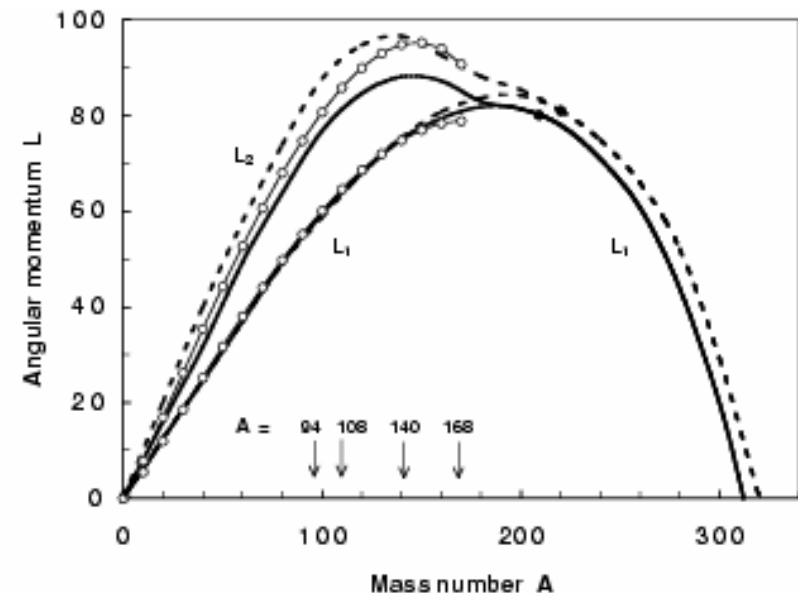
Carl Gustav Jacob Jacobi (1804 - 1851) discovered transition from oblate to triaxial shapes in the context of rotating, idealized, incompressible gravitating masses in 1834.

In 1961 Beringer and Knox suggested a similar transition in the case of atomic nuclei, idealized as incompressible, uniformly charged, liquid drops endowed with surface tension.

Liquid drop calculation

Jacobi transition for  $L > L_1$   
Fission barrier vanishes for  $L > L_2$

W.D. Myers and W.J. Swiatecki  
Acta Phys. Pol. B 32, 1033 (2001)



## What is the signature of a Jacobi transition in nuclei ?

- sharp decrease of frequency with increasing angular momentum (giant backbend of the moment of inertia)
- frequency of collective rotation is related to the E2  $\gamma$ -ray energy:  $\hbar\omega = \frac{1}{2} E_\gamma$
- many rotational bands at high spin quasi-continuous transitions
- measure the energy of the quasi-continuous 'E2 bump' as a function of angular momentum
- series of experiments with Gammasphere

$^{48}\text{Ca} + ^{50}\text{Ti}$  @ 200 MeV

$^{48}\text{Ca} + ^{64}\text{Ni}$  @ 207 MeV

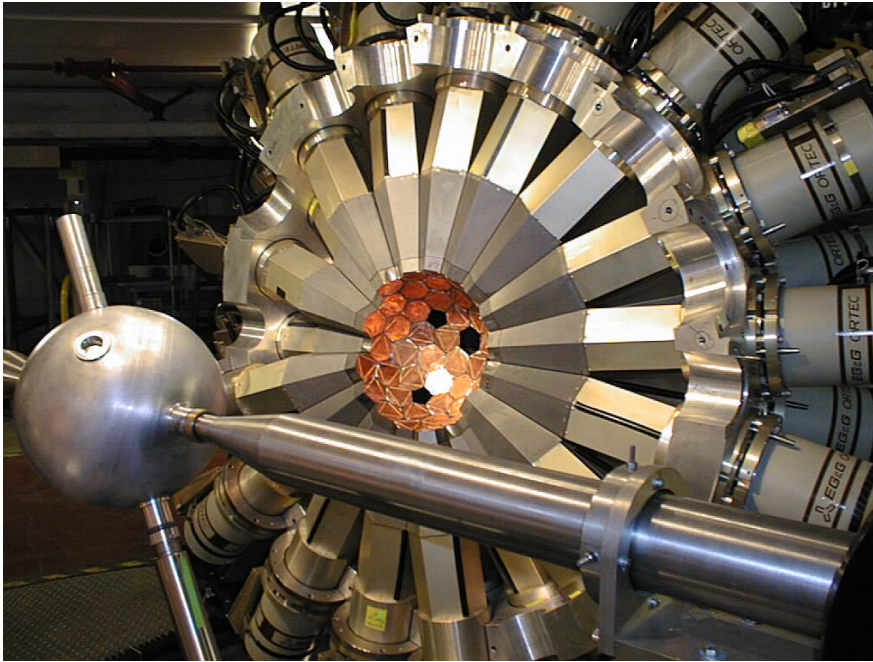
$^{48}\text{Ca} + ^{96}\text{Zr}$  @ 207 MeV

$^{48}\text{Ca} + ^{124}\text{Sn}$  @ 215 MeV

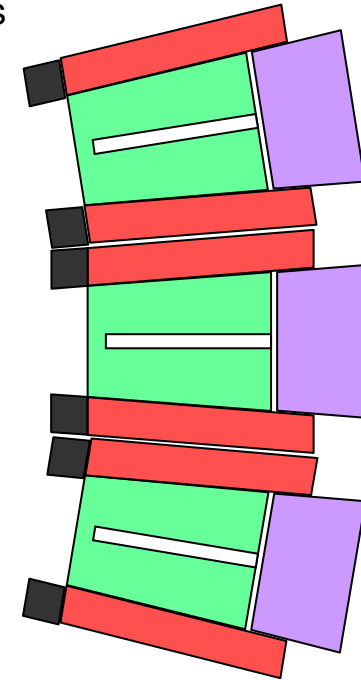
as neutron rich as possible:  
 $\Rightarrow$  higher fission barrier



# Measuring angular momentum with Gammasphere

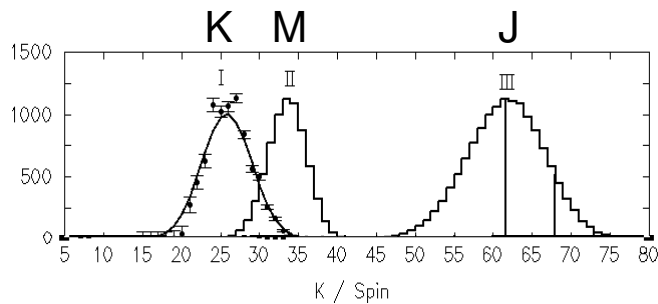


108 Compton-suppressed  
HPGe detectors



108 Ge detectors  
6 x 108 = 648 BGO detectors

increase in false veto signals  
reduced Ge efficiency  
but very high granularity

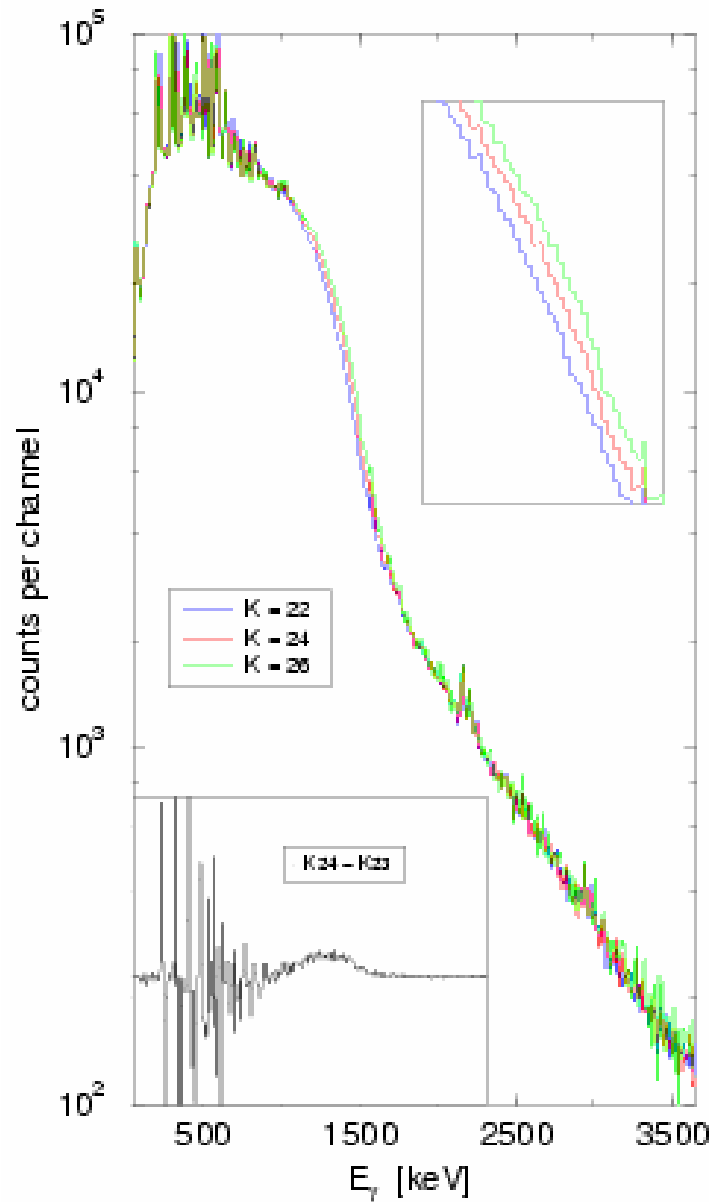


K = number of hits = fold

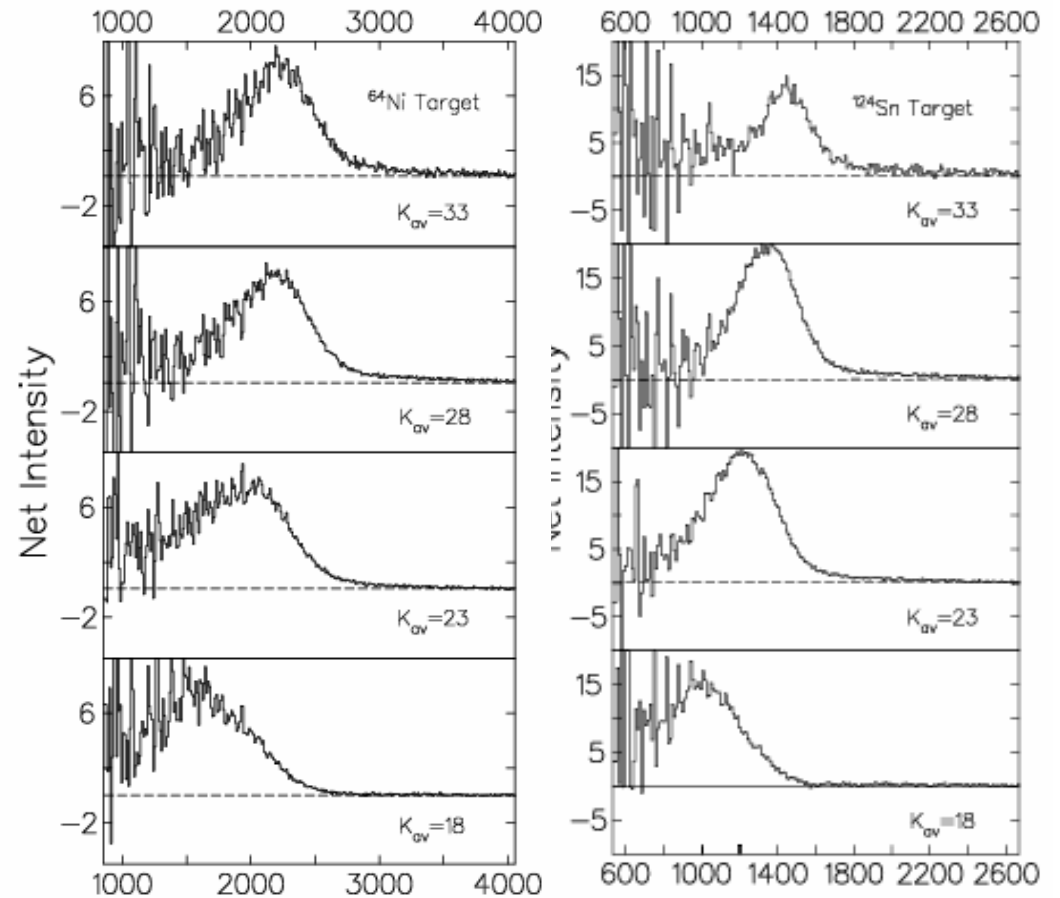
M =  $\gamma$  rays emitted = multiplicity (from response function)

J = initial angular momentum (from angular distribution)

# The E2 bump

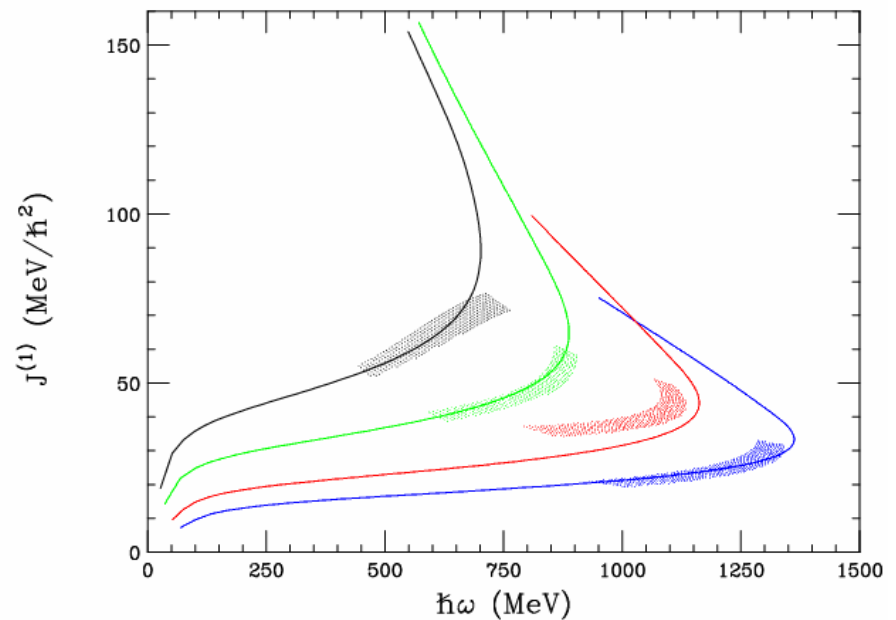
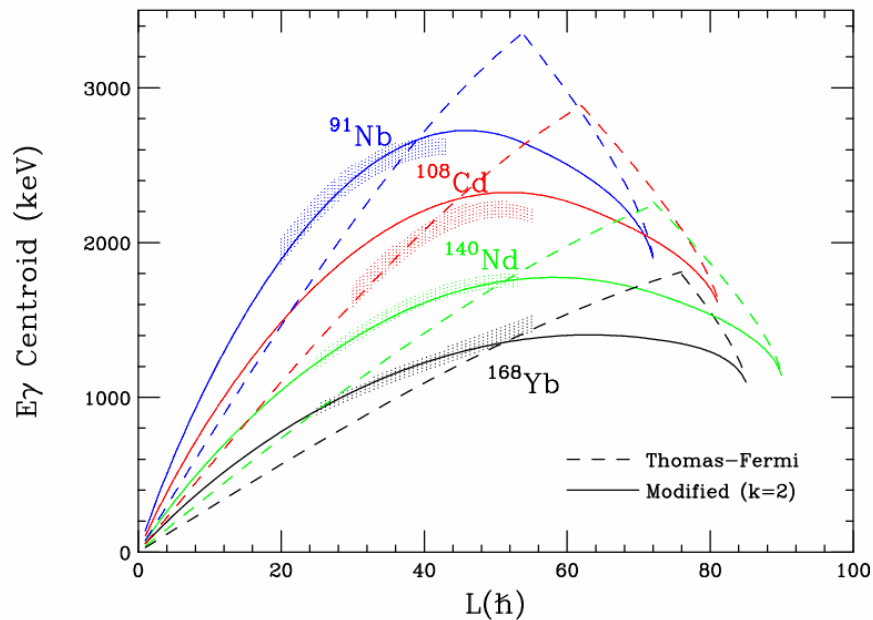


Incremental spectra:  
 Multiplicity ( $K_{av}-1$ ) gated spectrum  
 subtracted from ( $K_{av}+1$ ) spectrum



**K measures the angular momentum**  
**E2 bump measures rotational frequency**

## Comparison to liquid drop calculations



D. Ward et al., Phys. Rev. C 66, 024317 (2002)

two modifications:

- lower effective moment of inertia at low spin due to pairing
- no collective rotation about axially symmetric (MacLaurin) shapes in nuclei, instead, collective rotations are associated with (mostly) prolate shapes  
→ no sharp transition caused by breaking of axial symmetry, but smooth transition



# Charged-particle evaporation

$^{40}\text{Ca}+^{40}\text{Ca}$  @ 167 MeV

The nucleus of interest is often only weakly populated compared to a large background of other nuclei.

neutrons are deeply bound, charged-particle evaporation favored despite Coulomb barrier

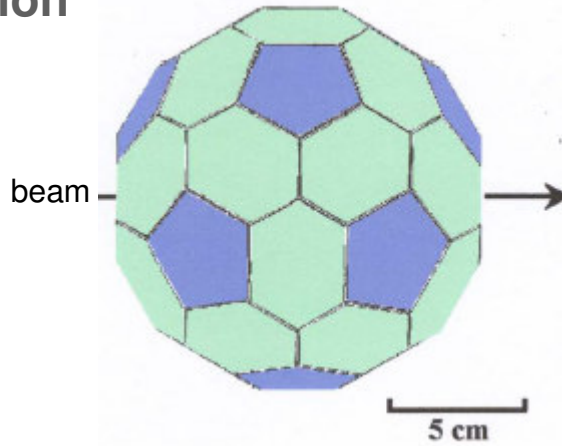
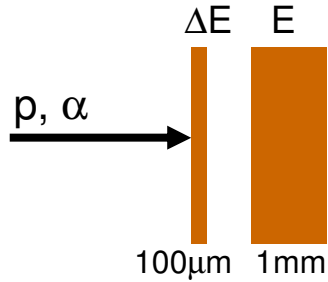
- Additional sensitivity from:
- charged-particle detectors
  - neutron detectors
  - recoil detectors
  - tagging techniques

						$^{78}\text{Zr}$ 2n	$^{79}\text{Zr}$ 1n	$^{80}\text{Zr}$
					$^{76}\text{Y}$ p3n	$^{77}\text{Y}$ p2n 0.18	$^{78}\text{Y}$ pn	$^{79}\text{Y}$ 1p
					$^{75}\text{Sr}$ $\alpha$ n	$^{76}\text{Sr}$ 2p2n 4.06	$^{77}\text{Sr}$ 2pn 2.95	$^{78}\text{Sr}$ 2p
			$^{72}\text{Rb}$ $\alpha$ p3n	$^{73}\text{Rb}$ $\alpha$ p2n 0.01	$^{74}\text{Rb}$ $\alpha$ pn 2.49	$^{75}\text{Rb}$ $\alpha$ p 5.35	$^{76}\text{Rb}$ 3pn 101	$^{77}\text{Rb}$ 3p 2.31
	$^{70}\text{Kr}$ 2 $\alpha$ 2n	$^{71}\text{Kr}$ 2 $\alpha$ n 0.18	$^{72}\text{Kr}$ 2 $\alpha$ 0.37	$^{73}\text{Kr}$ $\alpha$ 2pn 52	$^{74}\text{Kr}$ $\alpha$ 2p 20.3	$^{75}\text{Kr}$ 4pn 132	$^{76}\text{Kr}$ 4p 74.2	
	$^{69}\text{Br}$ 2 $\alpha$ p2n	$^{70}\text{Br}$ 2 $\alpha$ pn 6.82	$^{71}\text{Br}$ 2 $\alpha$ p 11.4	$^{72}\text{Br}$ $\alpha$ 3pn 38.2	$^{73}\text{Br}$ $\alpha$ 3p 128	$^{74}\text{Br}$ 5pn 3.23	$^{75}\text{Br}$ 5p 68.2	
	$^{67}\text{Se}$ 3 $\alpha$ n	$^{68}\text{Se}$ 3 $\alpha$ 5.07	$^{69}\text{Se}$ 2 $\alpha$ 2pn 1.57	$^{70}\text{Se}$ 2 $\alpha$ 2p 94	$^{71}\text{Se}$ $\alpha$ 4pn 0.46	$^{72}\text{Se}$ $\alpha$ 4p 102	$^{73}\text{Se}$ 6pn	$^{74}\text{Se}$ 6p 1.57
	$^{66}\text{As}$ 3 $\alpha$ pn	$^{67}\text{As}$ 3 $\alpha$ p 15	$^{68}\text{As}$ 2 $\alpha$ 3pn	$^{69}\text{As}$ 2 $\alpha$ 3p 35.6	$^{70}\text{As}$ $\alpha$ 5pn	$^{71}\text{As}$ $\alpha$ 5p 2.03	$^{72}\text{As}$ 7pn	$^{73}\text{As}$ 7p
$^{64}\text{Ge}$ 4 $\alpha$ 1.48	$^{65}\text{Ge}$ 3 $\alpha$ 2pn	$^{66}\text{Ge}$ 3 $\alpha$ 2p 8.3	$^{67}\text{Ge}$ 2 $\alpha$ 4pn	$^{68}\text{Ge}$ 2 $\alpha$ 4p 0.37	$^{69}\text{Ge}$ $\alpha$ 6pn			

cross sections in mb

# Charged particle detection

Si telescope



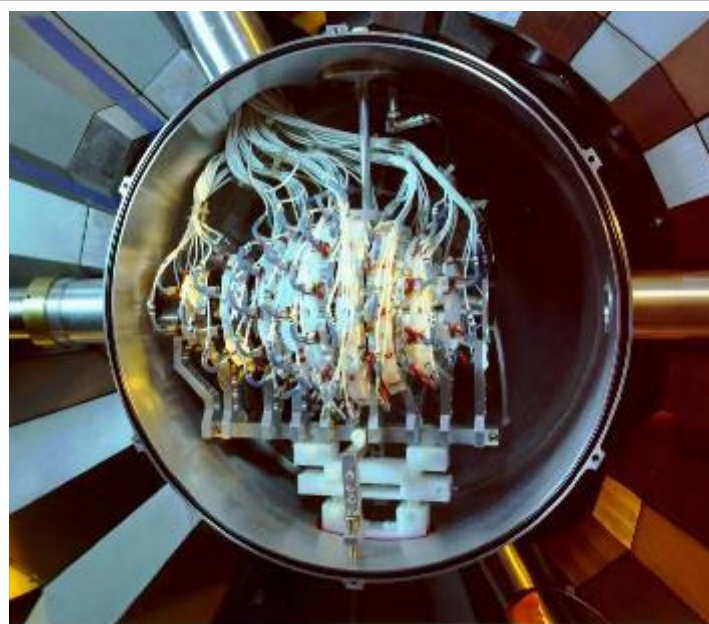
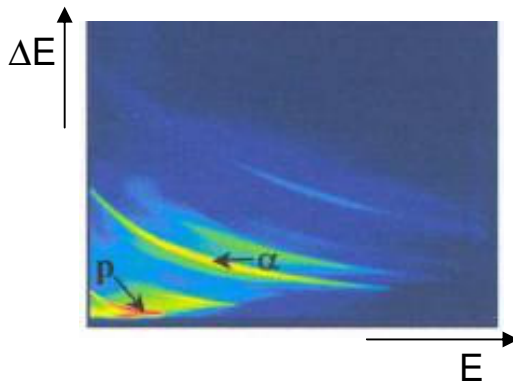
Italian Silicon Sphere ISIS  
Laboratori Nazionali di Legnaro

30 hexagons  
12 pentagons

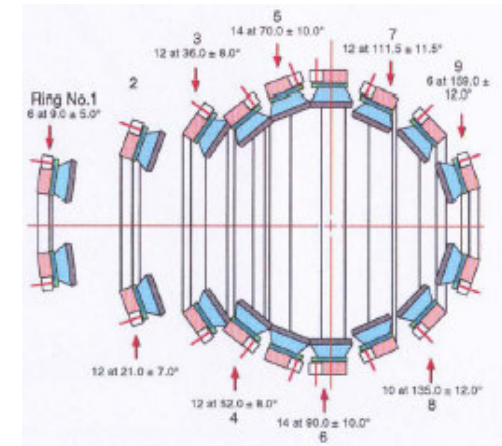
used with GASP and Euroball

stopping power

$$\frac{dE}{dx} \propto \frac{mZ^2}{E}$$



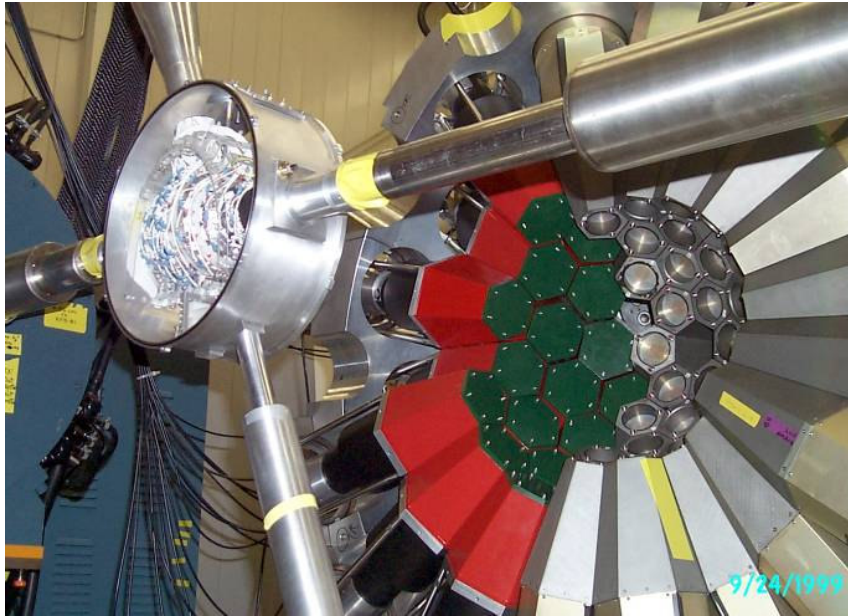
Microball,  
Washington University St. Louis



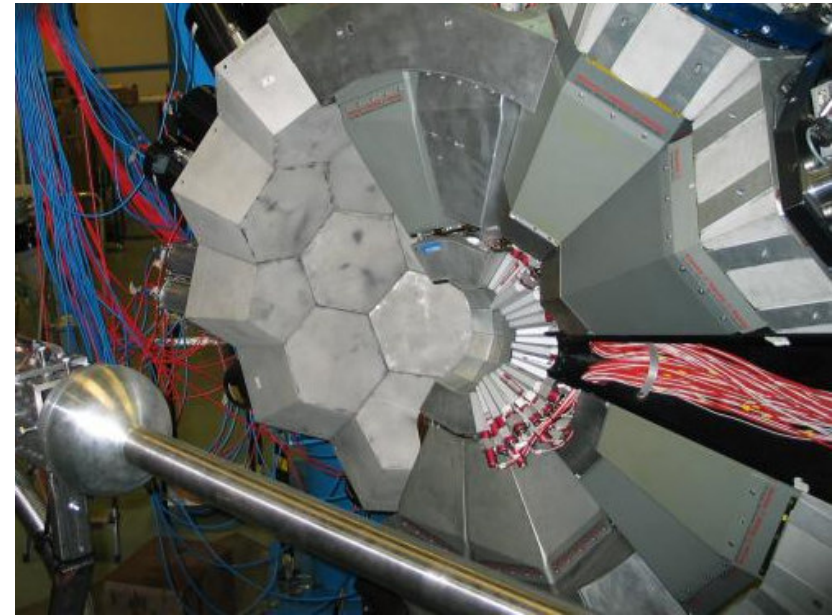
95 CsI(Tl) Scintillators  
in 9 rings

used with  
Gammasphere

# Neutron detection

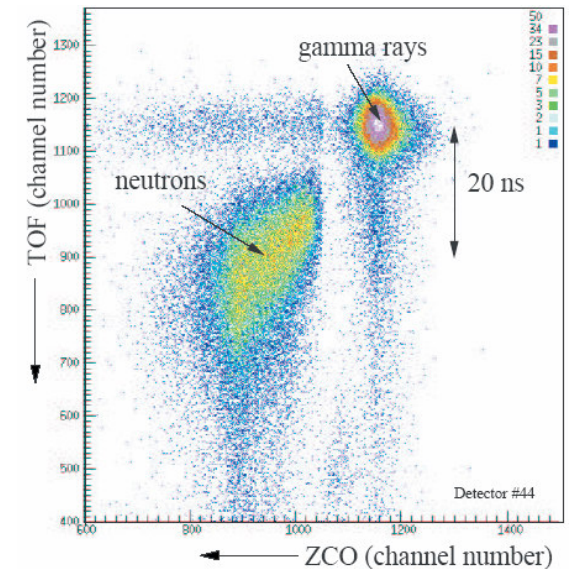


Gammasphere with Microball and Neutron shell (Washington University, St. Louis)



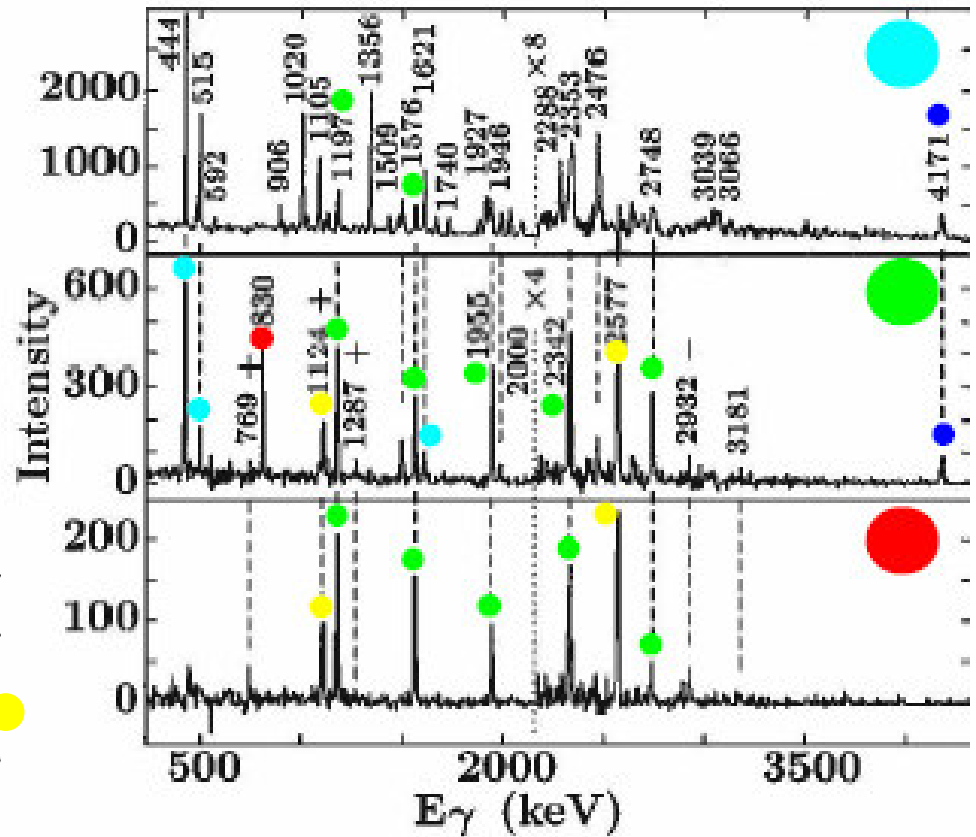
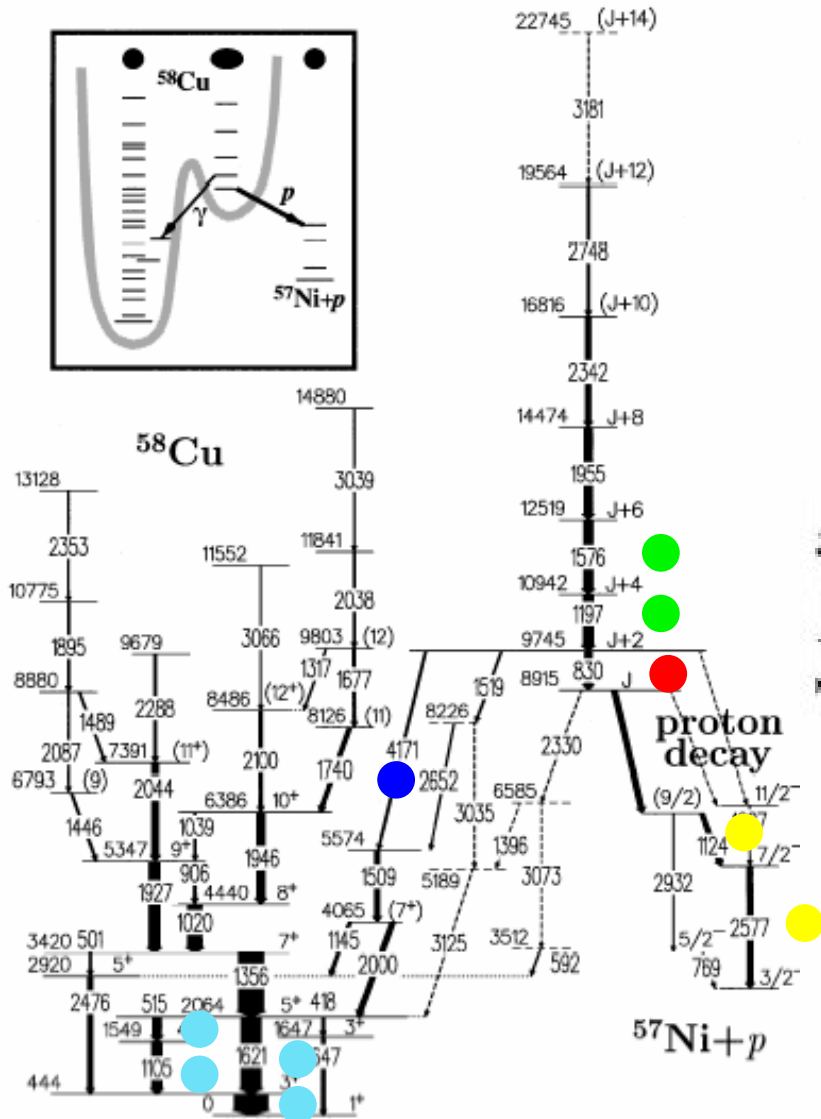
Euroball with Neutron wall (Uppsala University)

- can be used to select or veto neutron evaporation
- most powerful together with charged-particle detection
- used to study nuclei near  $N=Z$  line
  - isospin symmetry
  - proton-neutron pairing
  - shape coexistence
  - astrophysical rapid-proton capture process
- neutrons are separated from  $\gamma$  rays by time of flight and pulse shapes (zero-crossing time)
- difficult to distinguish two-neutron hit from scattering



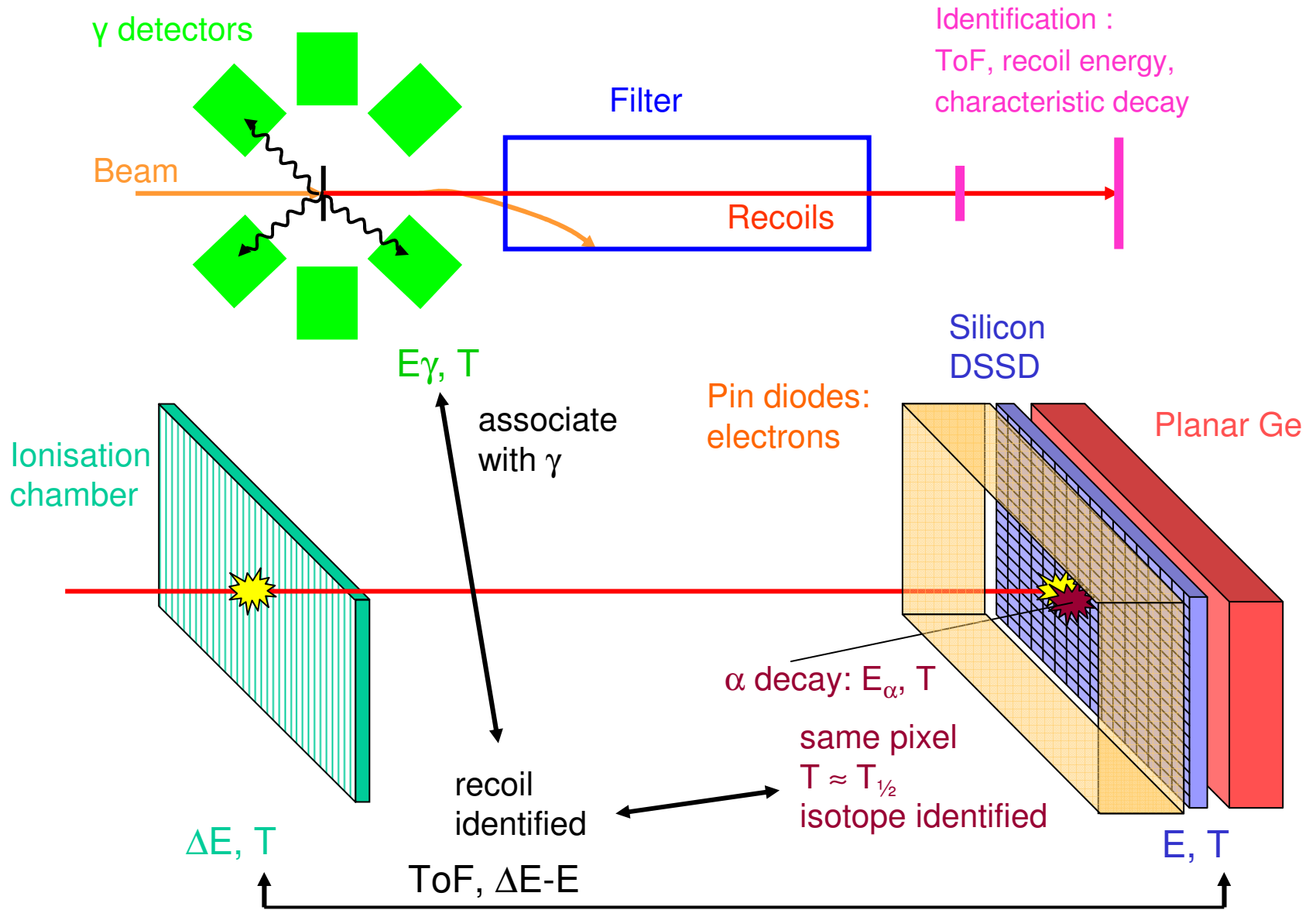
# Prompt proton decay in $^{58}\text{Cu}$

- $^{28}\text{Si}(^{36}\text{Ar}, \alpha pn)^{58}\text{Cu}$
- gate on  $1\alpha + 1p + 1n$
- $\sigma_{\text{rel}} = 0.3\%$
- Gammasphere, Microball, Neutron Shell

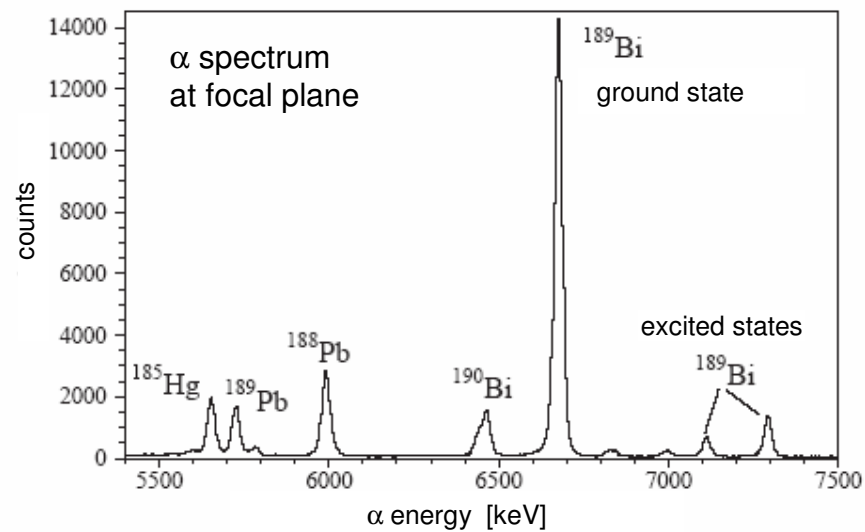
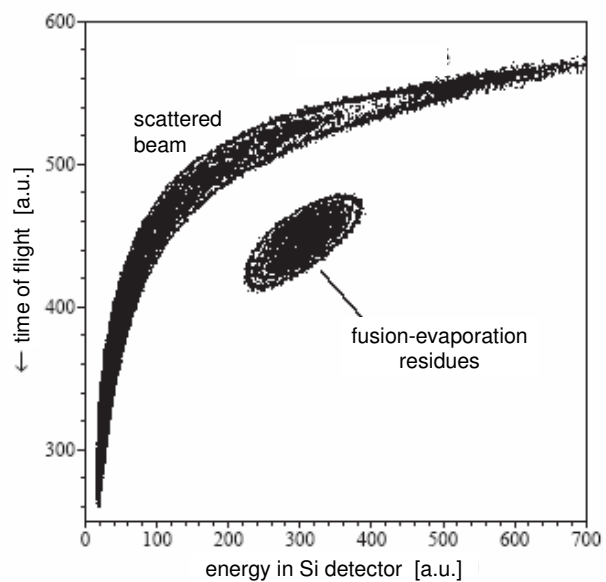
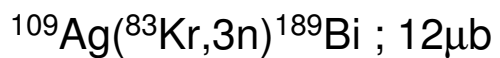
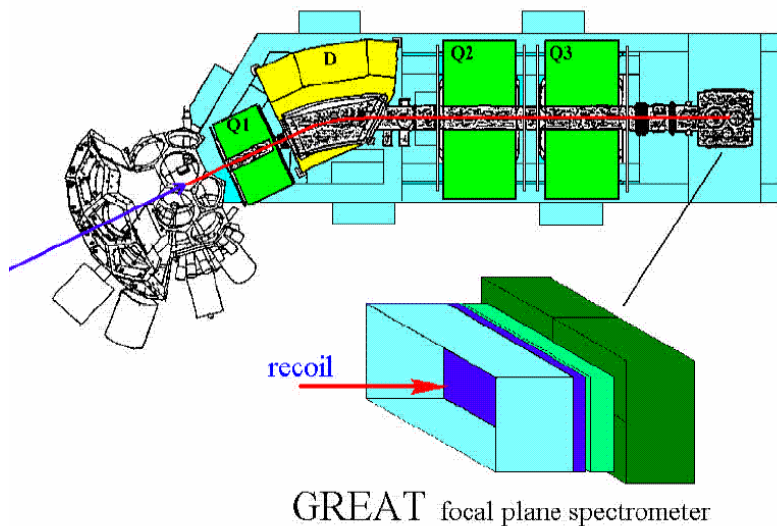


D. Rudolph et al., Phys. Rev. Lett. 80, 3018 (1998)  
Eur. Phys. J. A 14, 137 (2002)

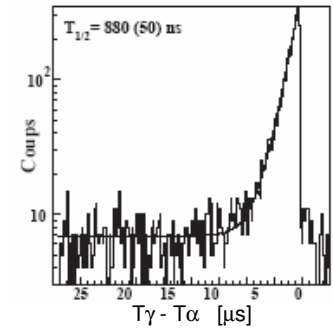
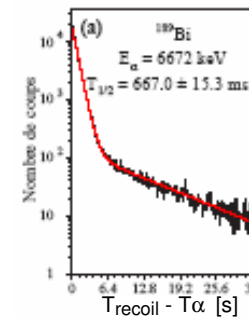
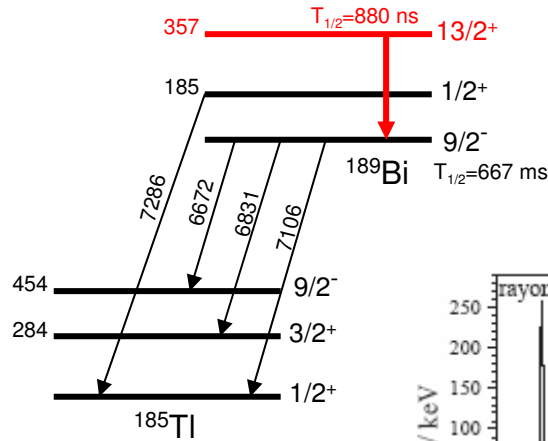
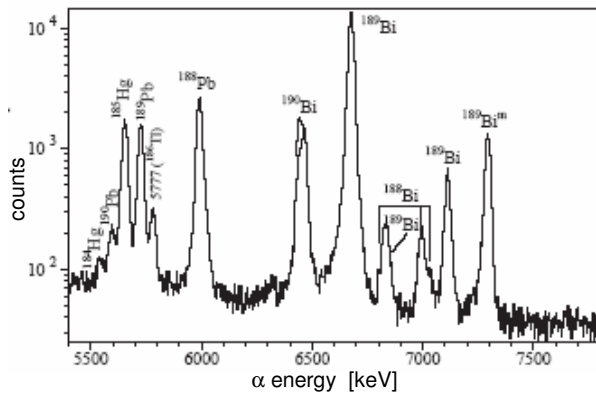
# Recoil decay tagging



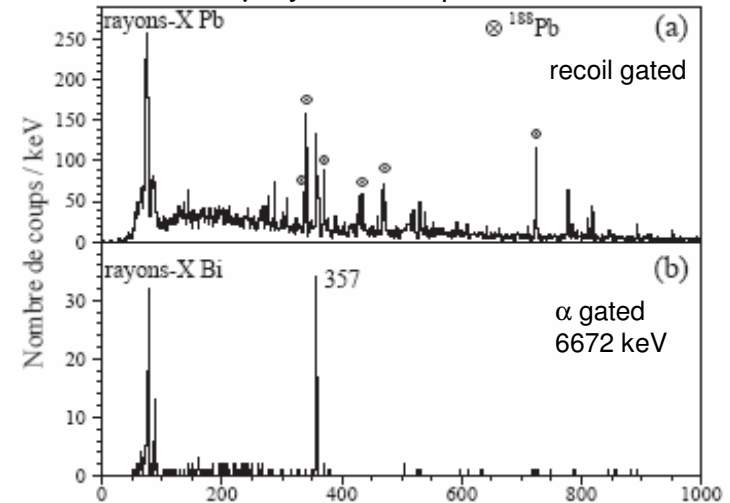
# JUROGAM – RITU – GREAT at Jyväskylä



# Recoil-decay tagging: $^{189}\text{Bi}$ spectra

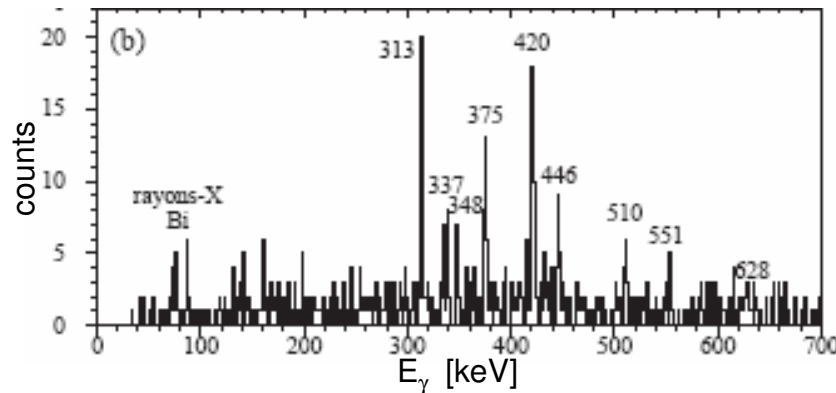
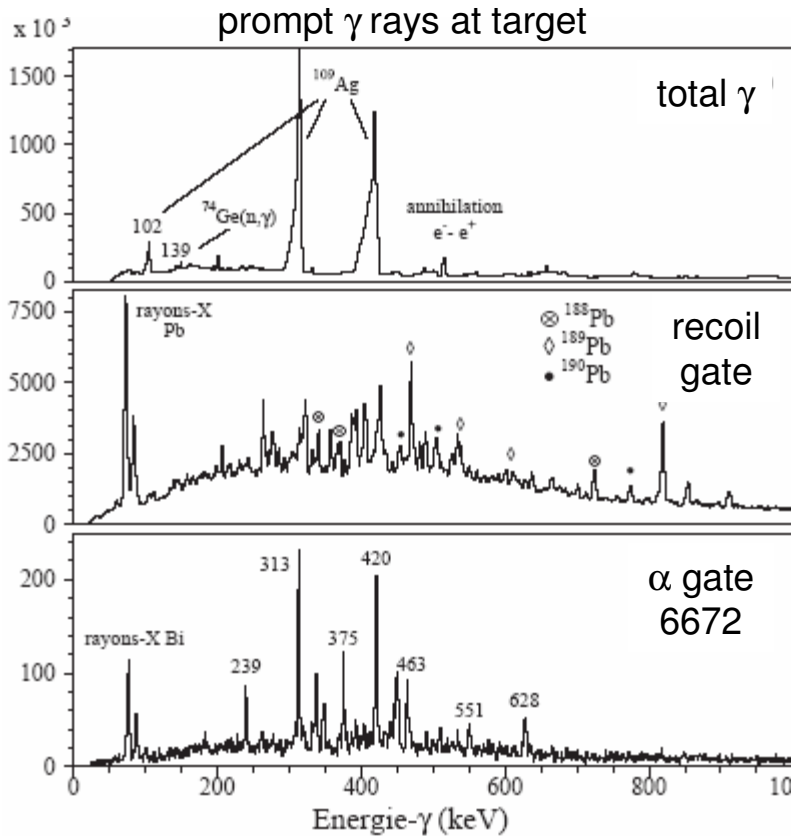


$\gamma$  rays at focal plane



$\alpha$  and  $\gamma$  tagging:  
 $\alpha$  gate: 6672 keV;  $\gamma$  gate: 357 keV

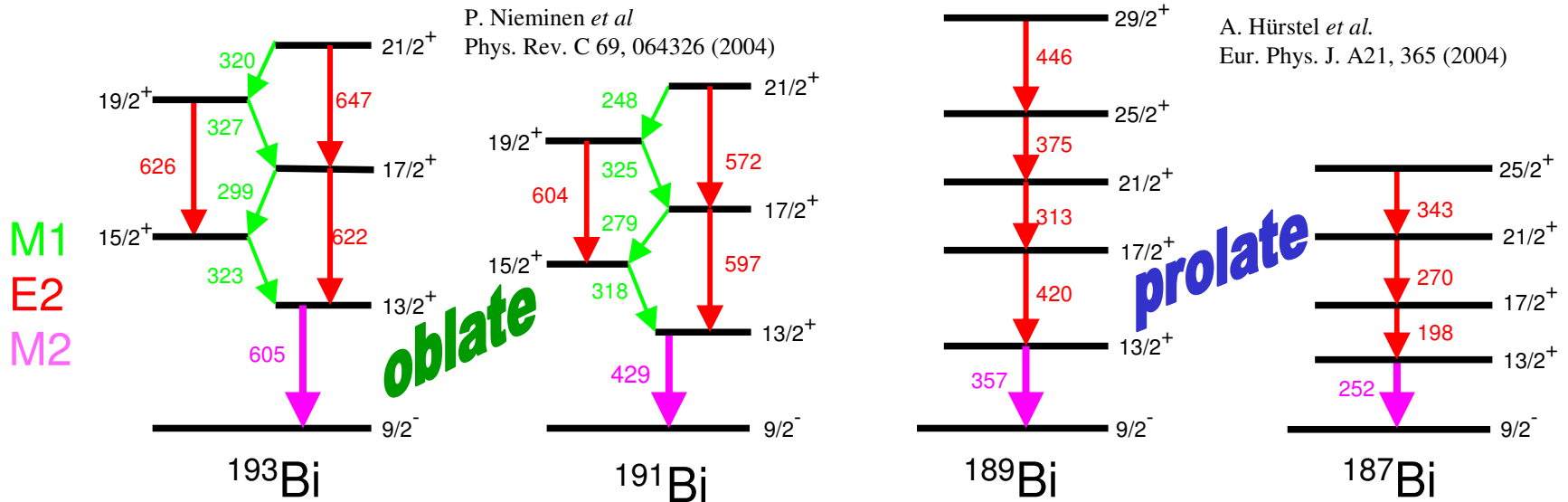
A. Hürstel et al.,  
 Eur. Phys. J. A 15, 329 (2002)







# Systematics of the neutron-deficient Bi isotopes



strongly coupled bands  
 → deformation aligned

decoupled bands  
 → rotation aligned

