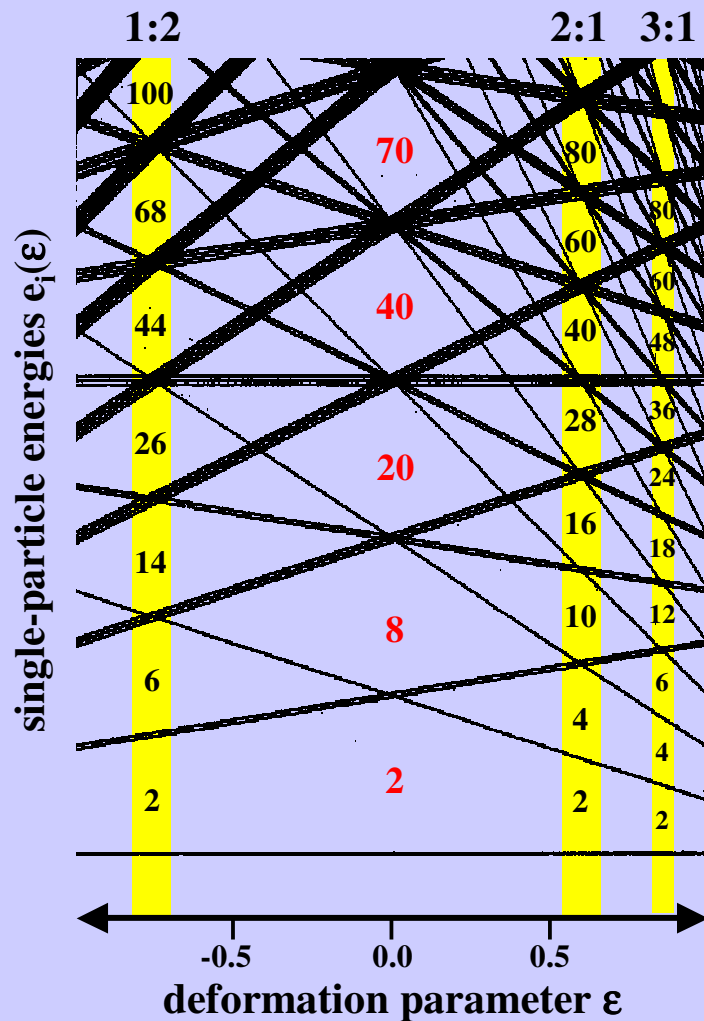


Single-particle energies in deformed harmonic oscillator potential

$$V = \frac{1}{2}M[\omega_{x,y}^2(x^2 + y^2) + \omega_z^2 z^2]$$

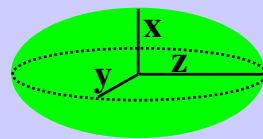
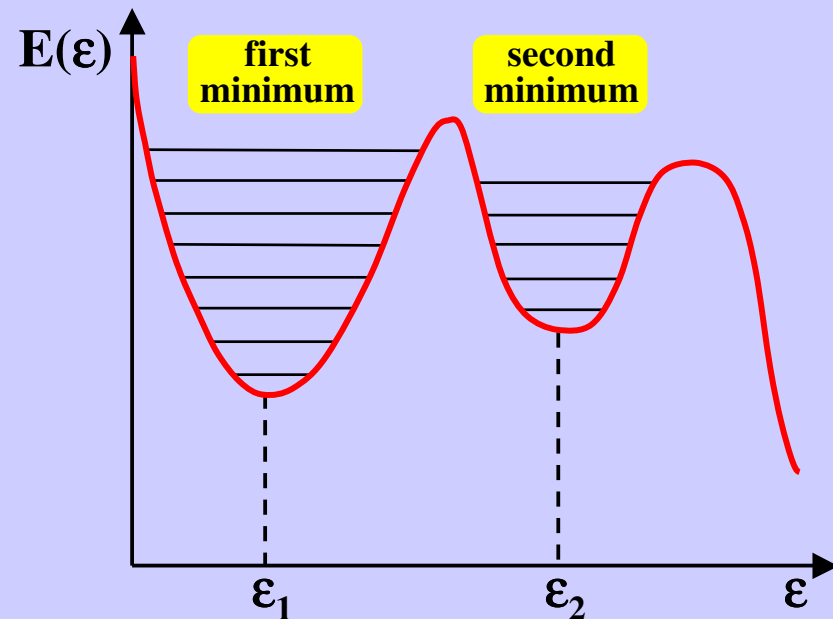


Potential energy of the nucleus as a function of the deformation:

$$E(\epsilon) = \sum e_i(\epsilon)$$

sum over the single-particle energies of all A nucleons

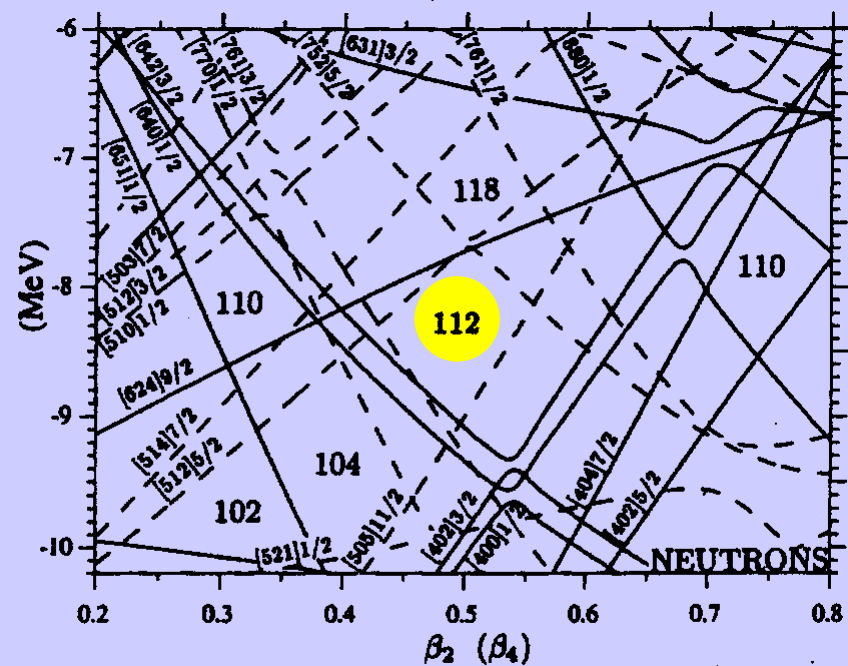
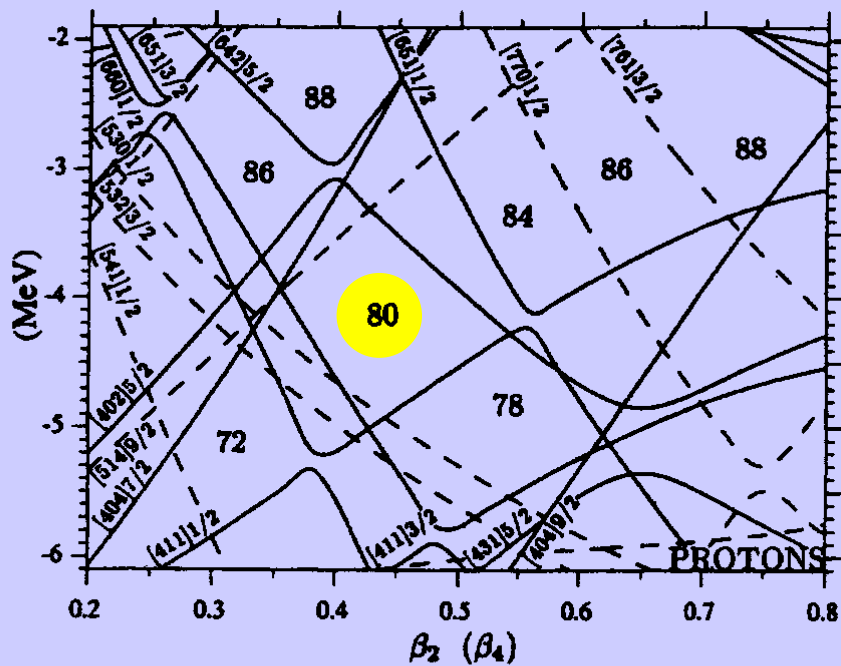
Due to the different slopes of the single-particle orbitals there might be more than one minimum for certain nucleon numbers !



$$x = y = R(1 + \epsilon)$$

$$z = \frac{R}{\sqrt{1 + \epsilon}} \approx R(1 - \frac{1}{2}\epsilon) \quad \epsilon = (\omega_{x,y} - \omega_z) / \omega_0$$

Single-particle energies in an anisotropic Woods-Saxon potential

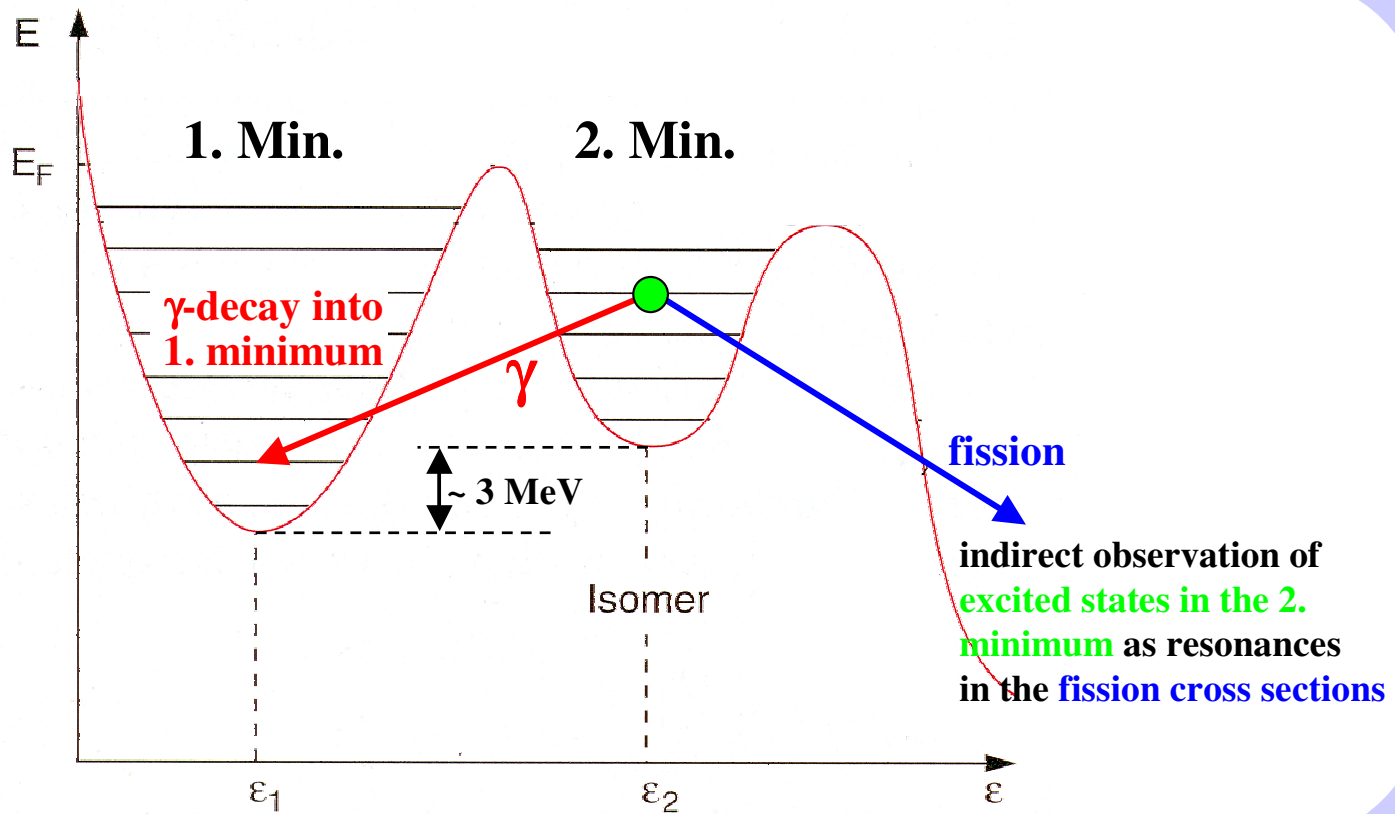


e.g.: $Z = 80, N = 112$



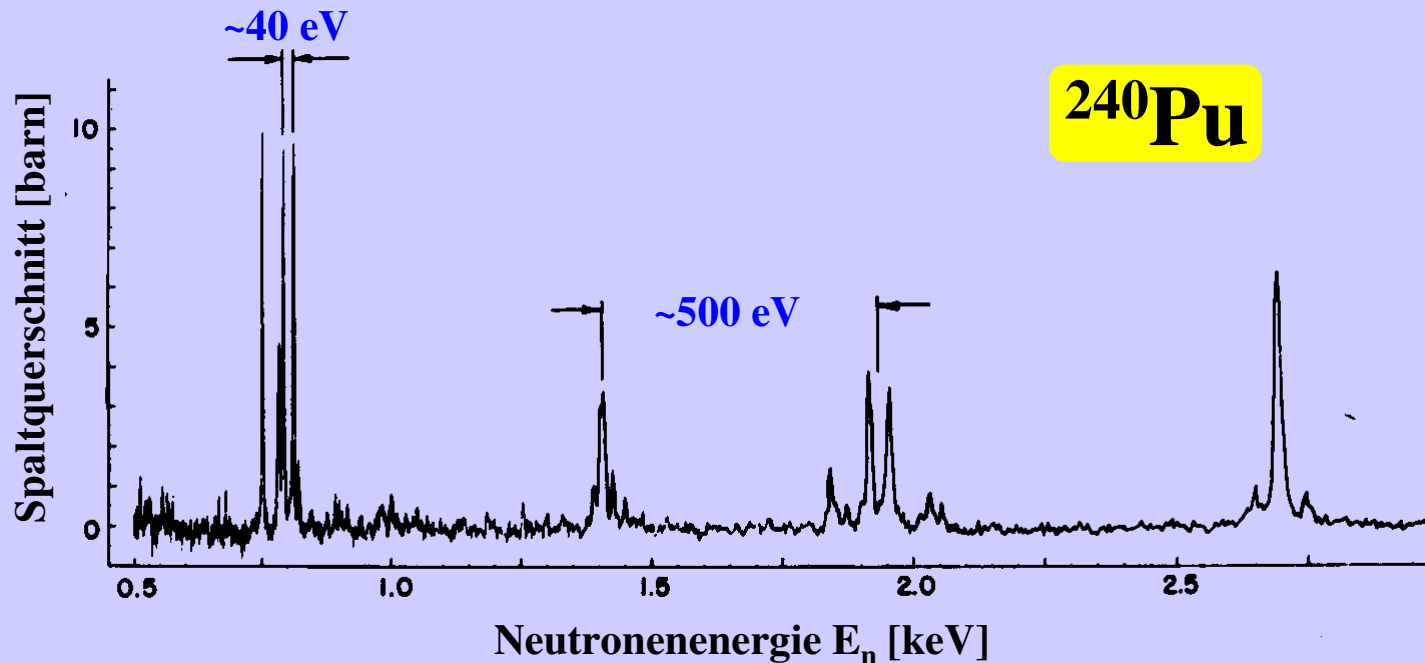
^{192}Hg

Possible decays out of the second minimum



An example: ^{240}Pu

Subbarrier fission cross section as function of the neutron energy:

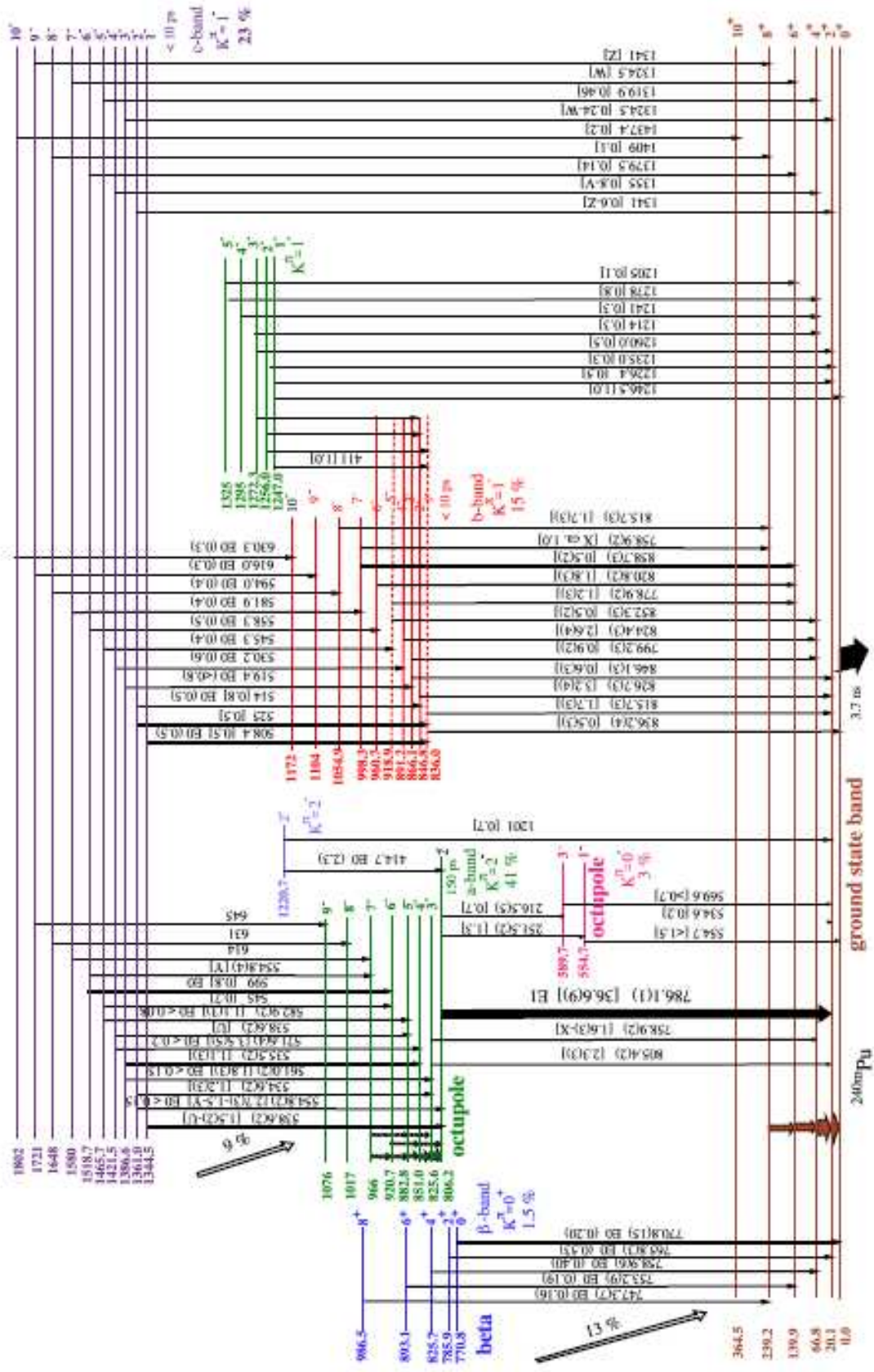


Double structure at energies far below the barrier !

No explanation in 1-barrier-picture, but ...

Spectroscopy in the Superdeformed Minimum of ^{240}Pu

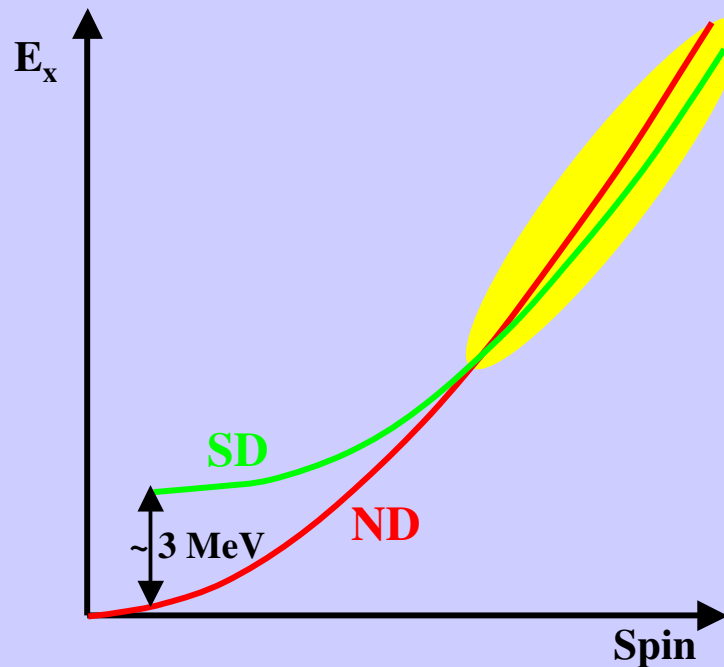
Pansegrau et al., Phys. Lett. 484B (2000) 1



γ -ray spectroscopy

moment of inertia larger for larger deformation

→ superdeformed rotational states energetically favoured
and therefore observable at **high spin** !



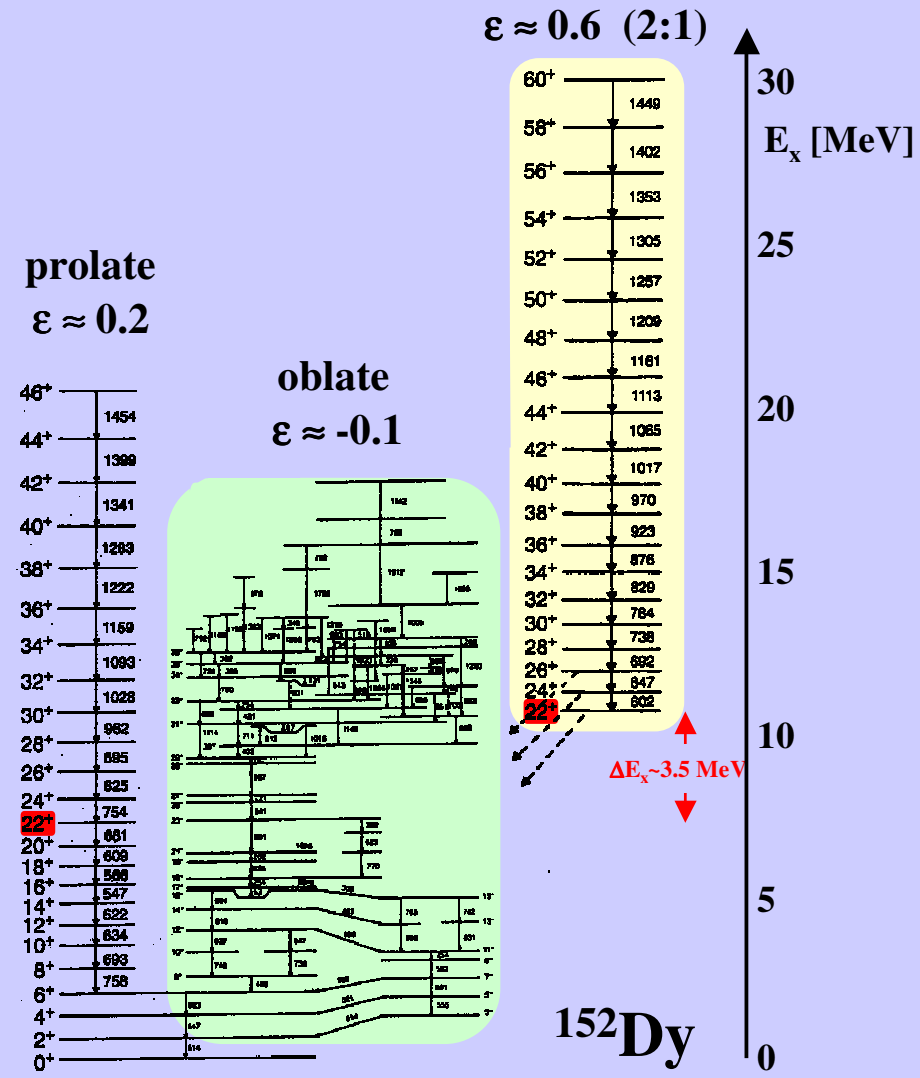
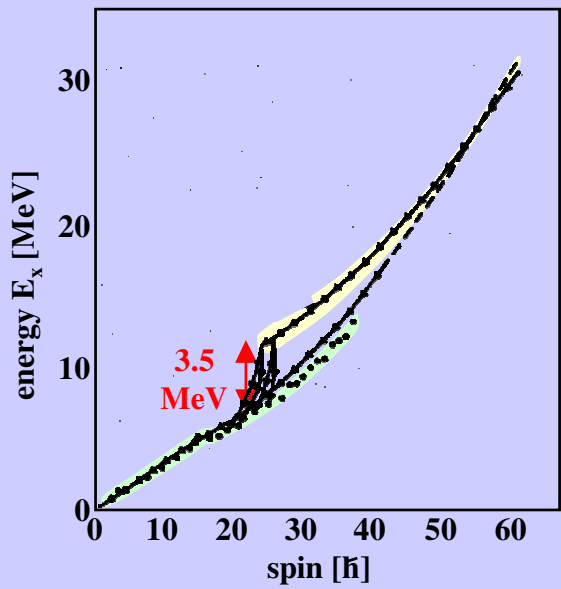
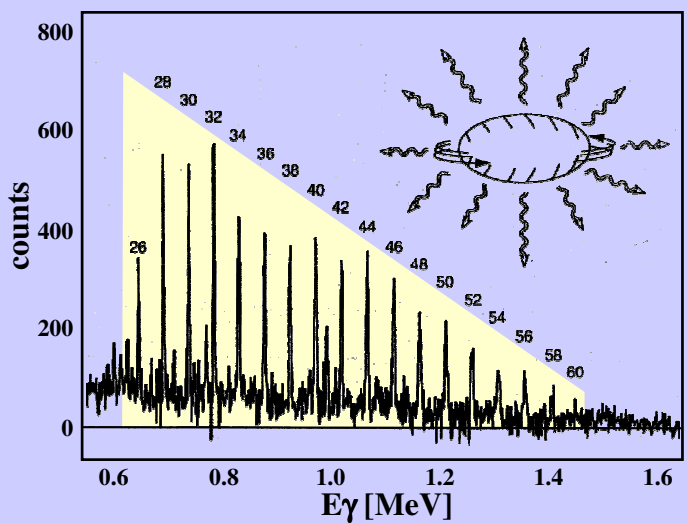
Rotational band:

$$E_x = \frac{\hbar^2}{2J} I(I+1)$$

J : moment of inertia

$$J_{SD} > J_{ND}$$

Discovery of superdeformation in $^{152}\text{Dy}_{86}$

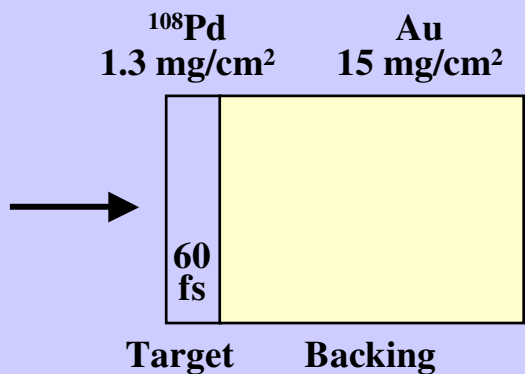


1986

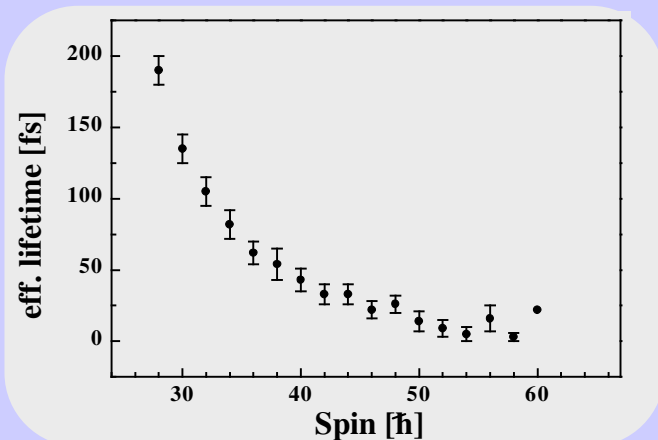
How to determine the deformation ?

Example: superdef. band in ^{152}Dy

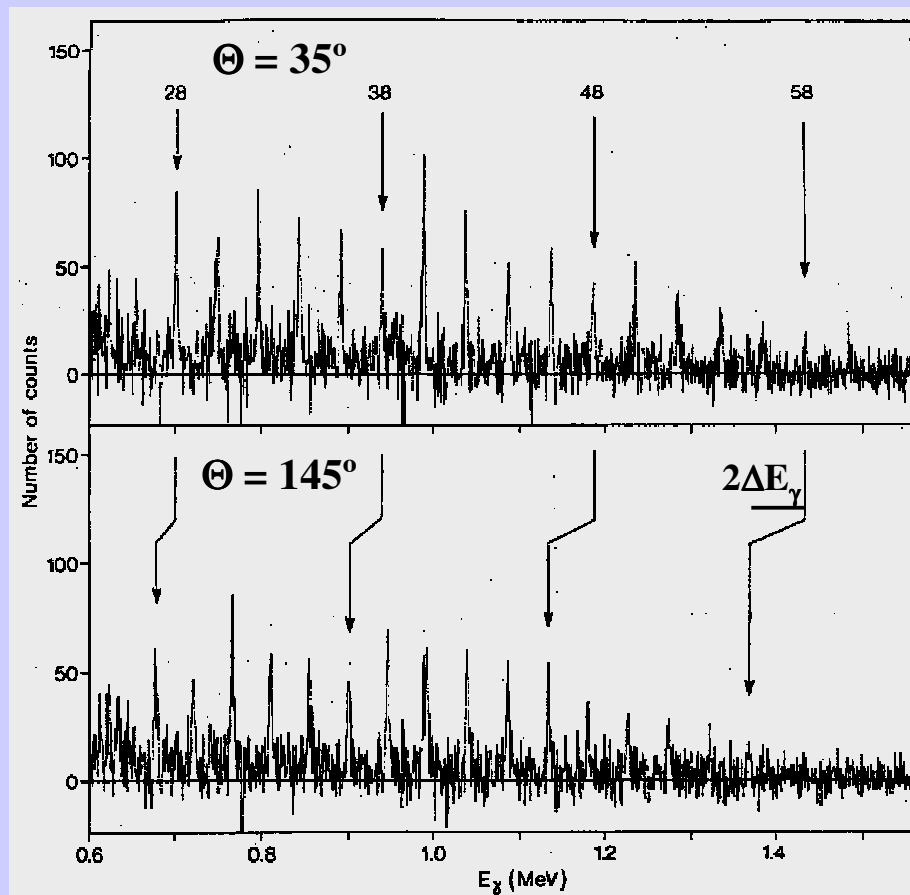
^{108}Pd ($^{48}\text{Ca}, 4n$) ^{152}Dy @ 205 MeV



$v_{\text{rec}}/c = 2.84 - 2.97 \%$ (target thickness !)



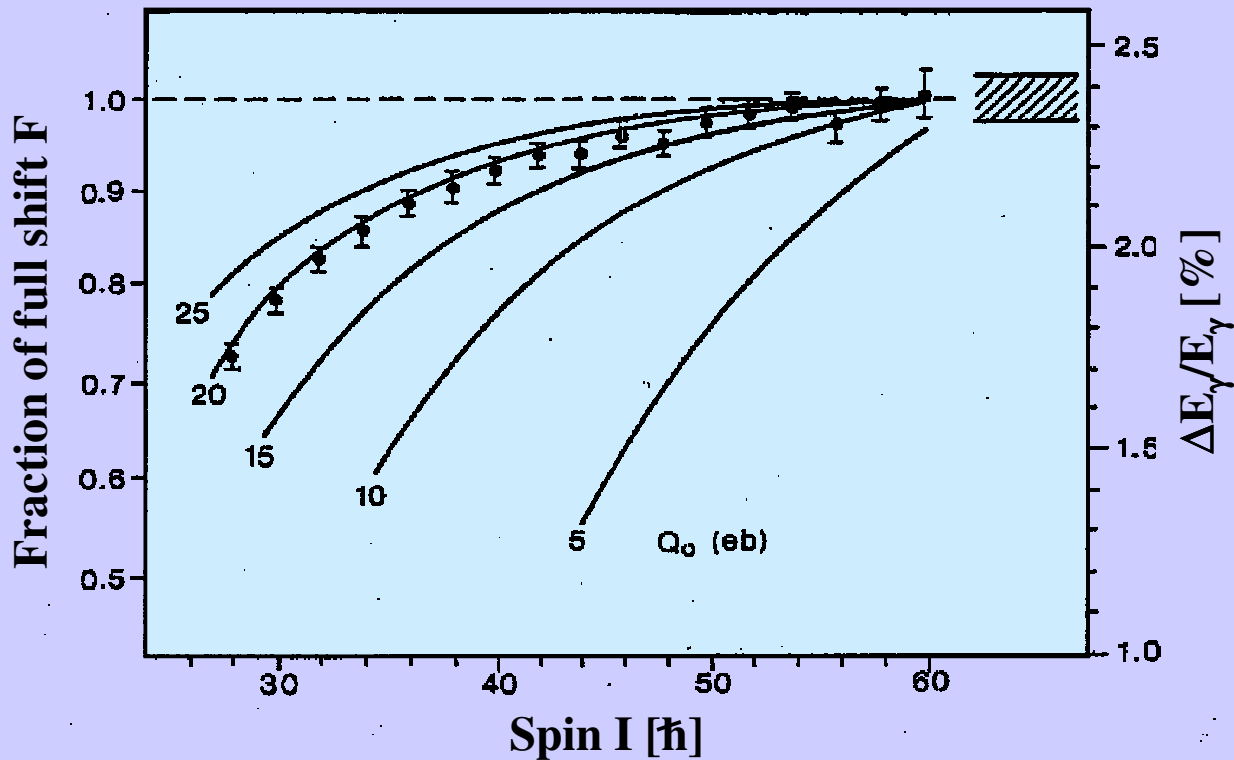
Spectrum of superdef. band



$$\Delta E_{\gamma}(v) = E_{\gamma} \cdot v/c \cos \theta = E_{\gamma} \cdot F(\tau) v_{\text{max}}/c \cos \theta$$

$$\text{mit } F(\tau) = \int_0^{\infty} \frac{1}{\tau} \frac{v(t)}{v_{\text{max}}} \cdot e^{-t/\tau} dt$$

Energy shifts as function of spin



Assuming constant deformation \longleftrightarrow constant Q_0

- calculation of $F(I)$ for different Q_0
- comparison with exp. values

→ $Q_0 = 19(3) \text{ e}\cdot\text{b}$
 $B(E2) \sim 2660 \text{ W.u.}$

$\beta \sim 0.82$

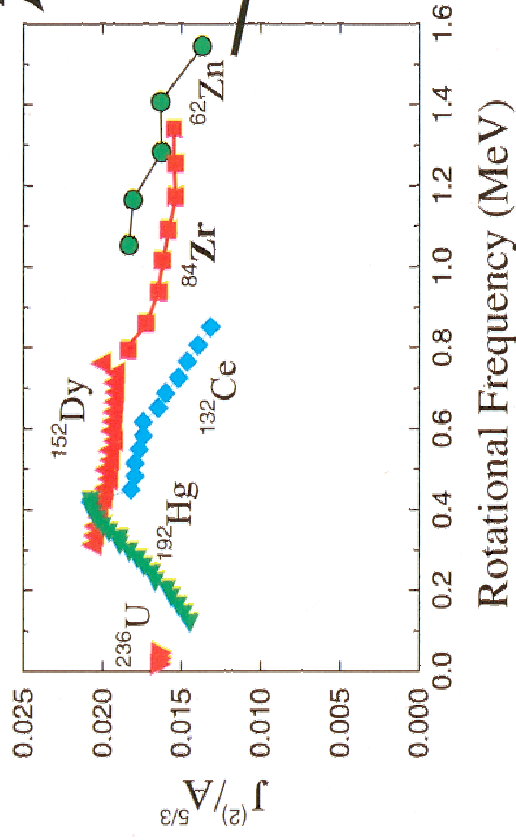
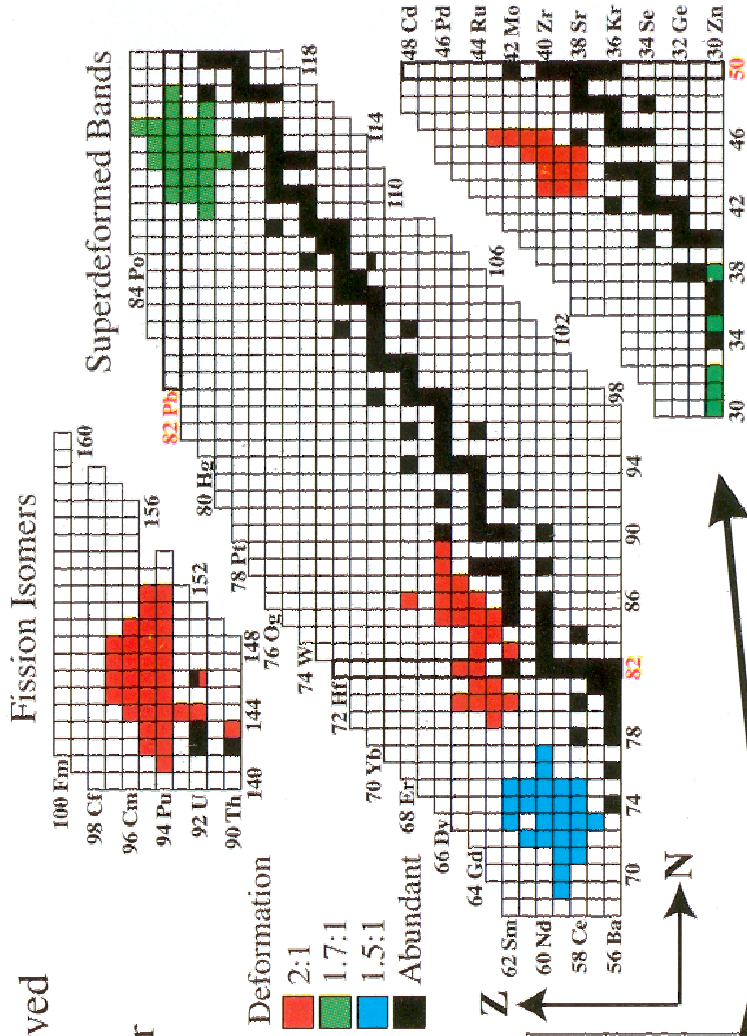
Proof of superdef. !

$$Q_0 = \frac{3}{\sqrt{5\pi}} ZR_0^2 \beta (1 + 0.16\beta)$$

mit $R_0 = 1.2 \cdot A^{1/3} \text{ fm}$

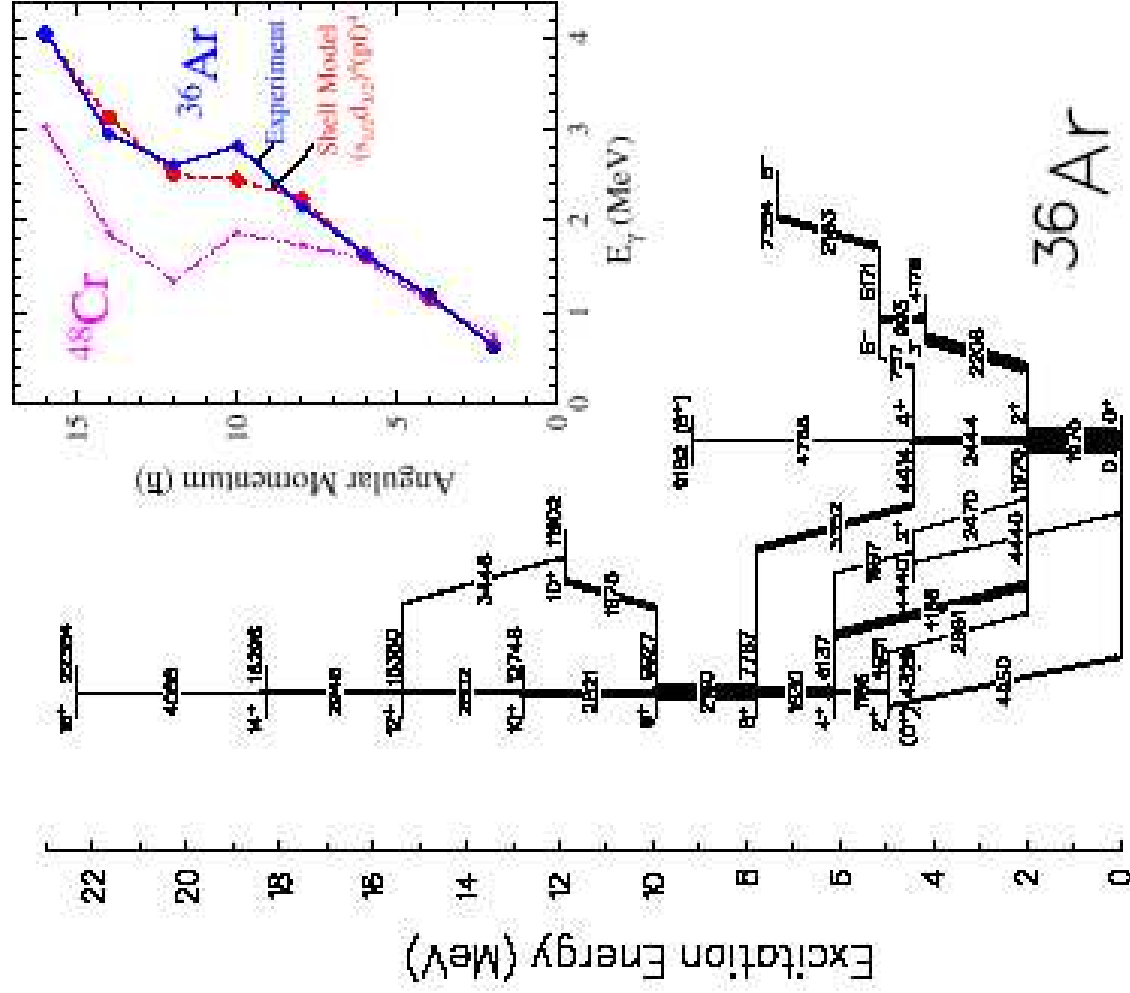
The Regions of Superdeformation

Superdeformed (SD) bands have been observed throughout the nuclear chart. The moments of inertia, after scaling to $A^{5/3}$, are all similar due to the fact that the shape (deformation) of a SD band is largely independent of the mass. The shell gaps, which define the SD shape, depend on the underlying symmetry of the nucleus, and occur at near-integer axis ratios.



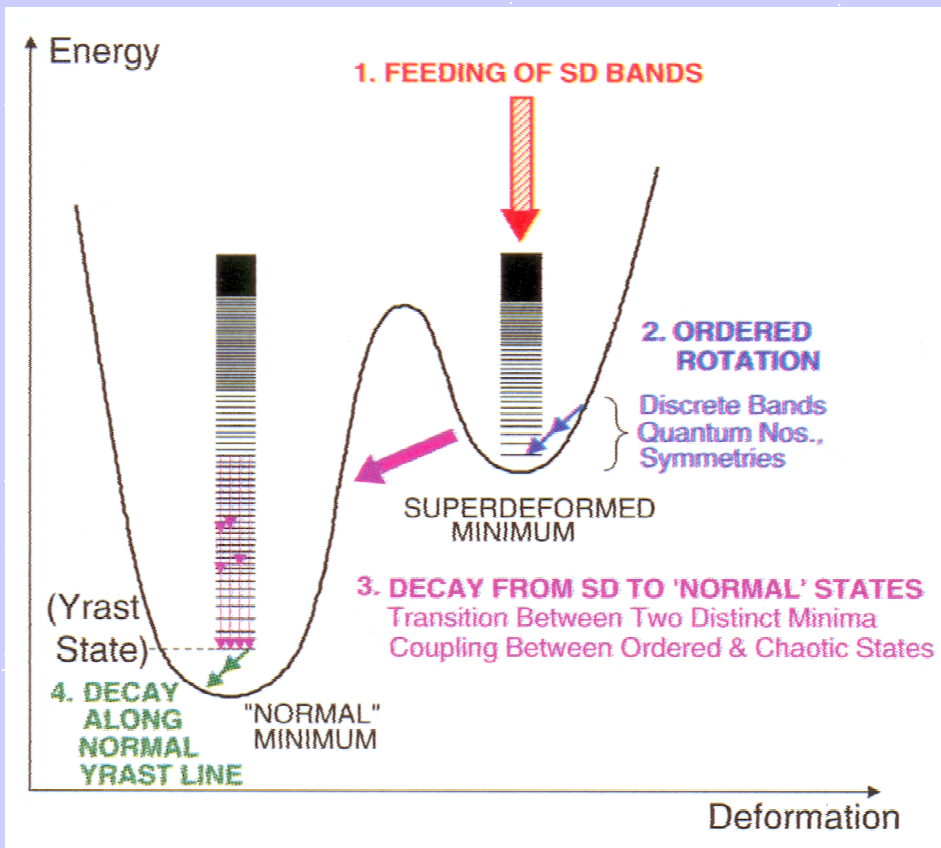
The recent observation with GAMMASPHERE of a new region of superdeformation ($N \sim Z \sim 30$) opens up many exciting opportunities, including the possibility to study the most collective states known with a "Shell Model" description.

C. Svensson et al., Phys. Rev. Lett. 85 (2000) 2693



The γ -decay of superdeformed bands

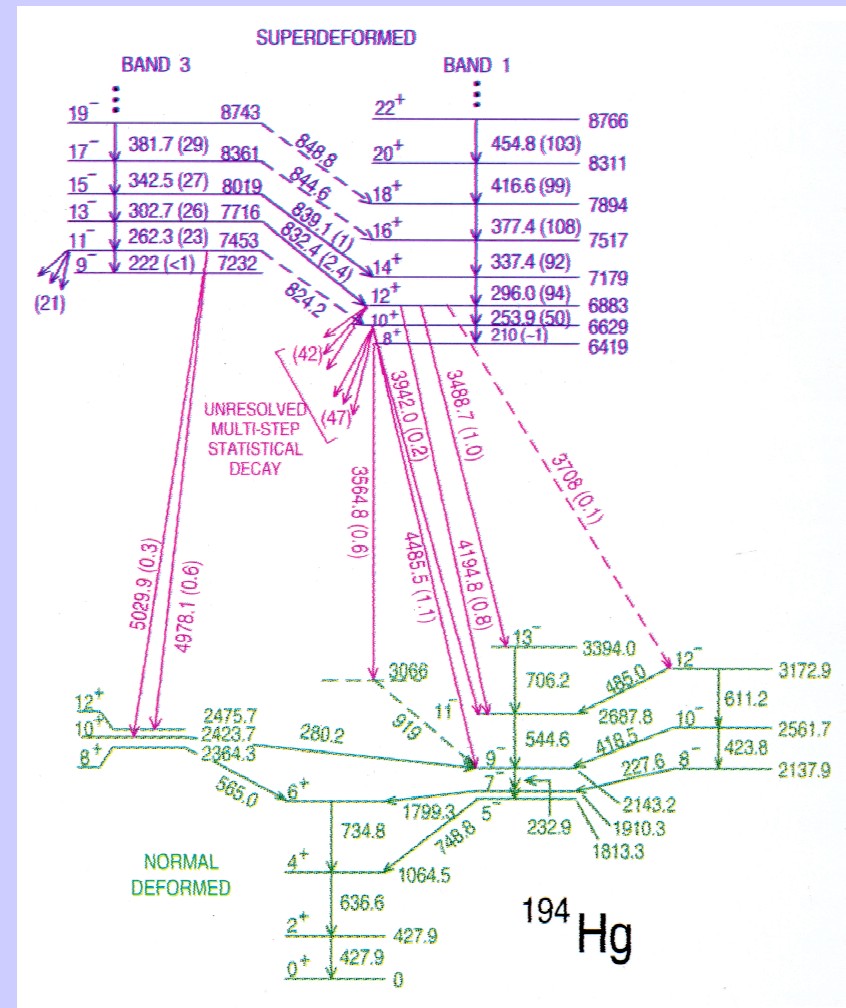
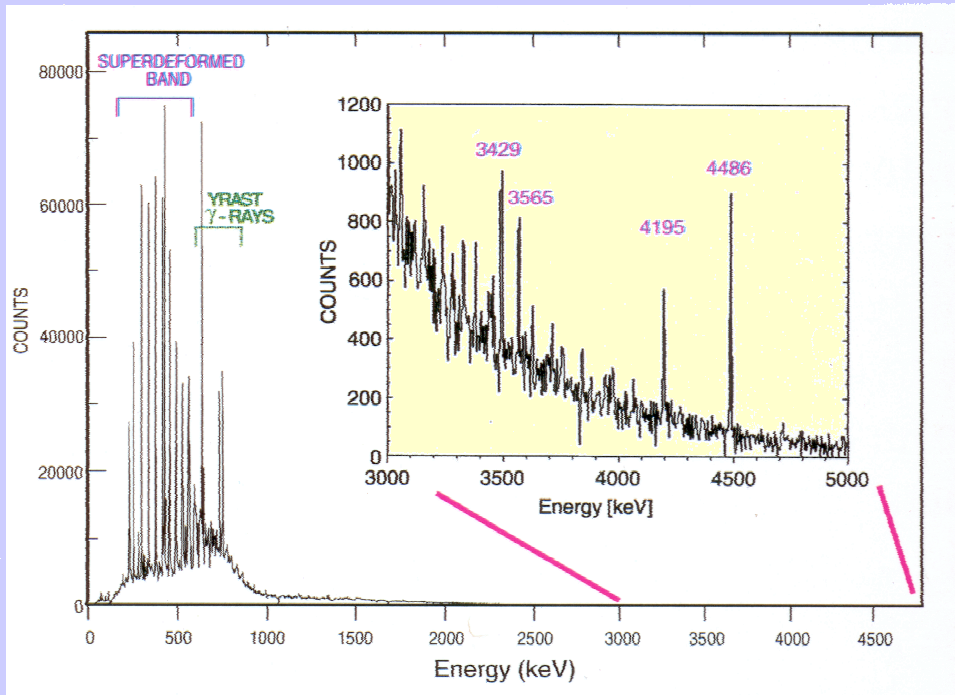
Why is it so difficult to observe ?



- high level density
→ fragmentation, low γ -intensities
- high γ -ray energies
→ low detection efficiency
- very different structure of the states in the two minima
→ low transition strengths

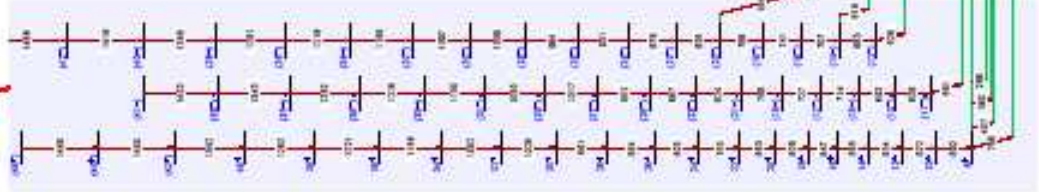
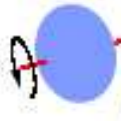
First observation of discrete linking γ -decays into the first minimum: The nucleus ^{194}Hg

Gammasphere-Experiment

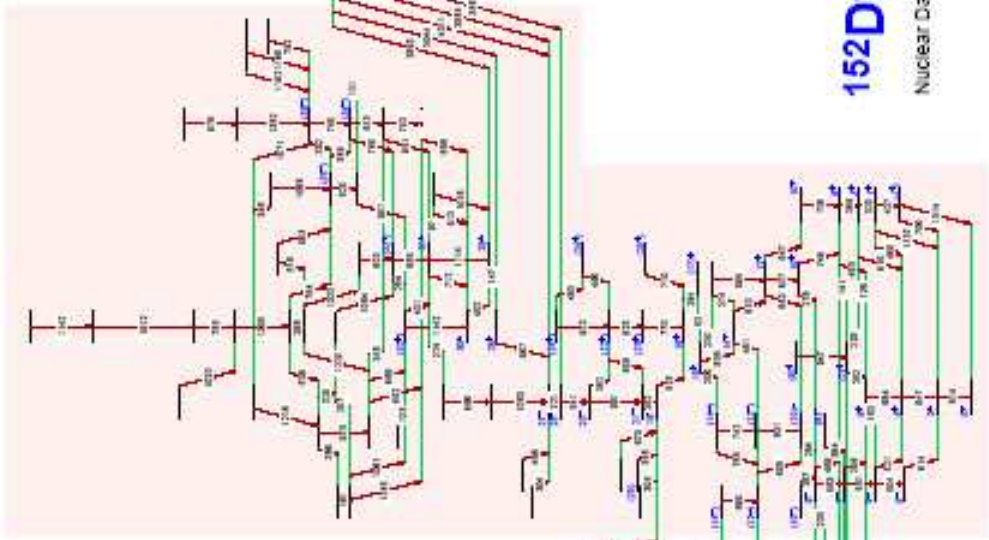


Coexistence of collective and noncollective motion

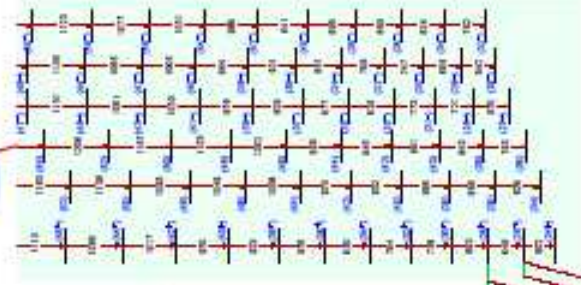
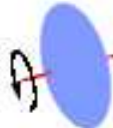
Triaxial bands



Noncollective states



Superdeformed bands



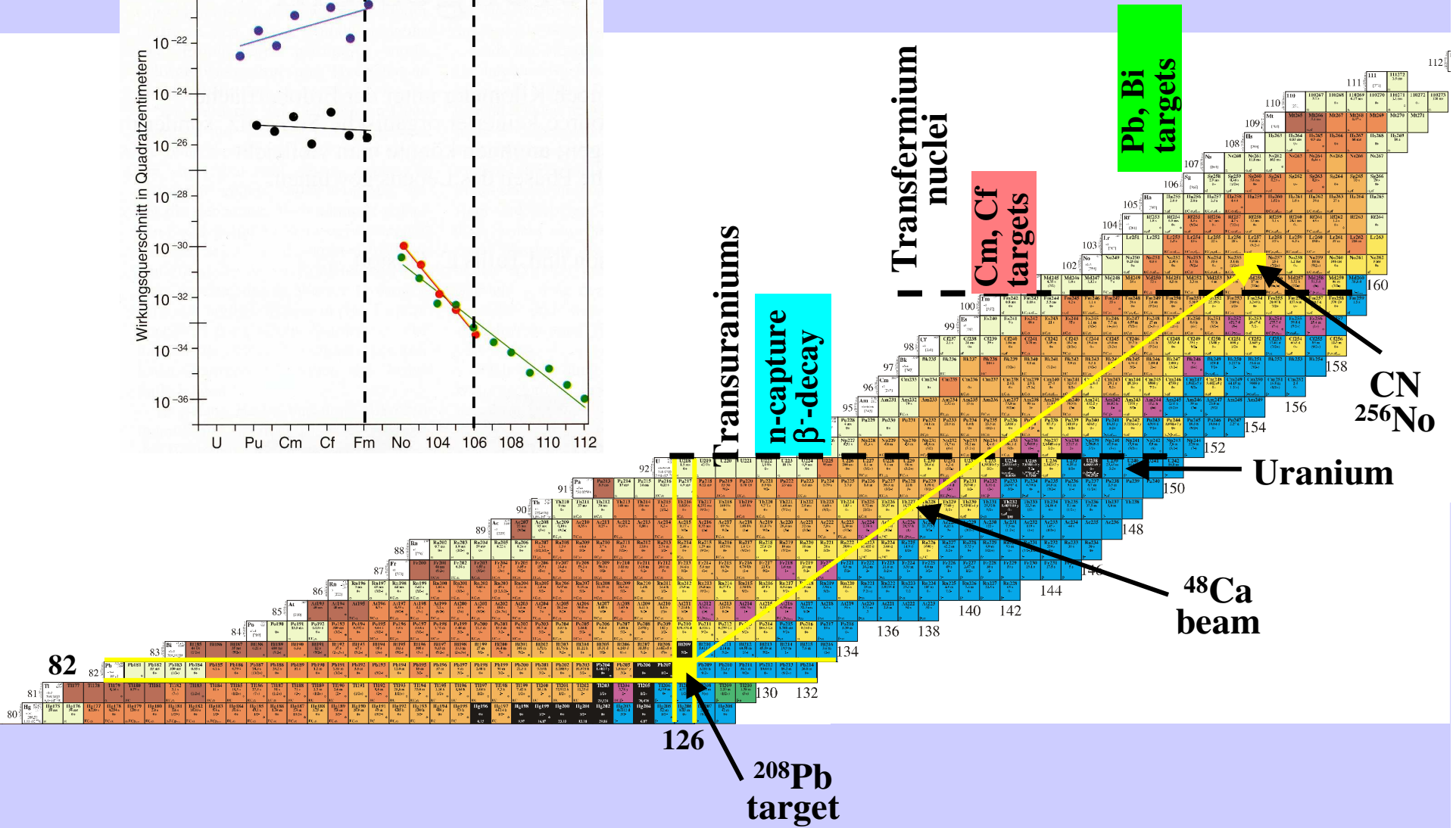
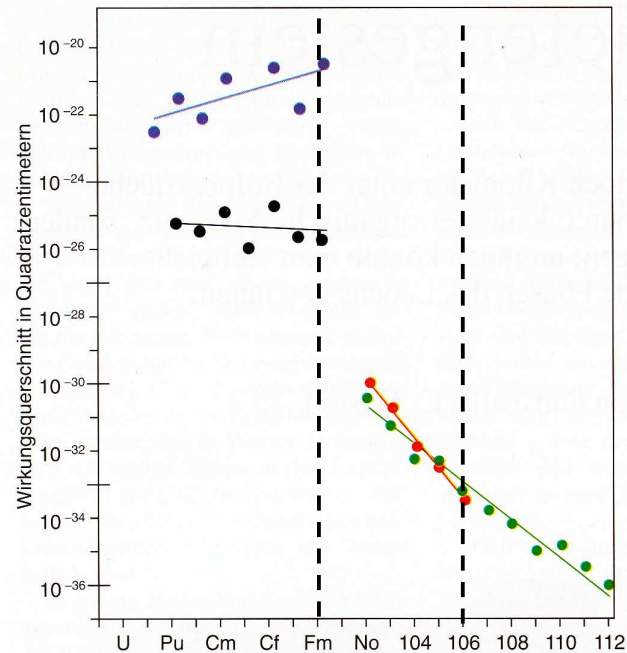
152Dy

Nuclear Data Sheets 95 (2002) 995

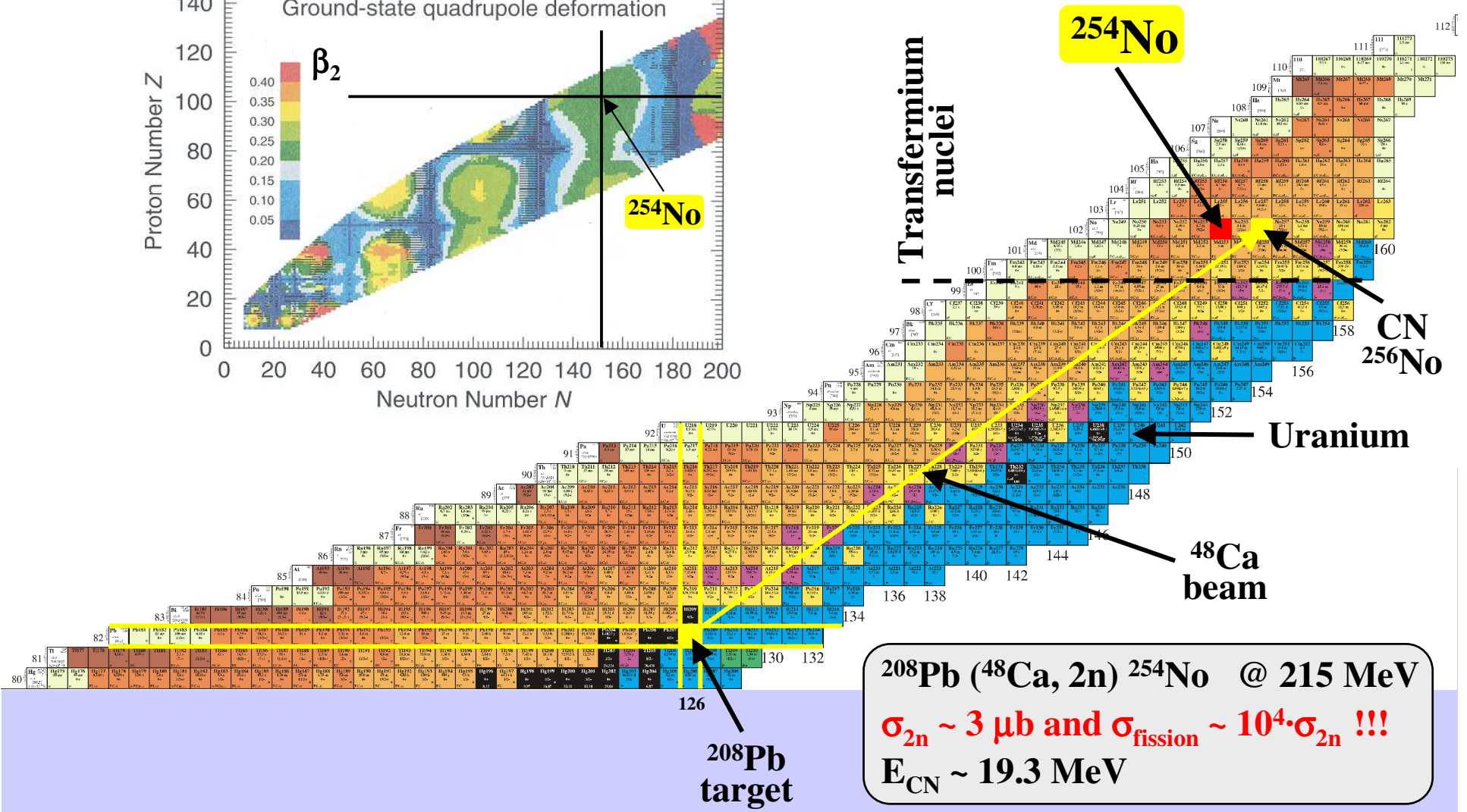
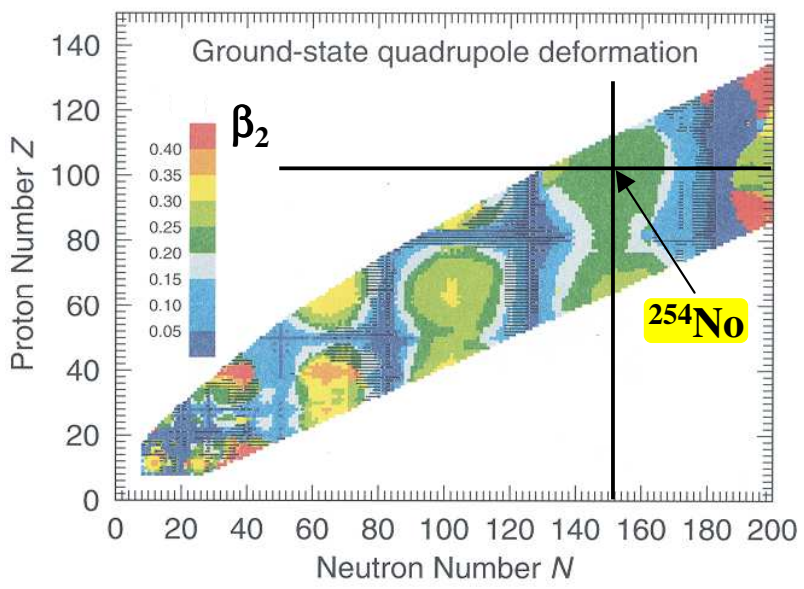
Some recent examples ...

- Prompt particle decay from deformed excited states
- Superdeformed bands all over the chart of nuclides
- **Spectroscopy of transfermium nuclei: Towards the SHE's ...**
- Ground state proton decay: spectroscopy beyond the dripline

Possible reactions for the synthesis of heavy elements



Where is ^{254}No ?

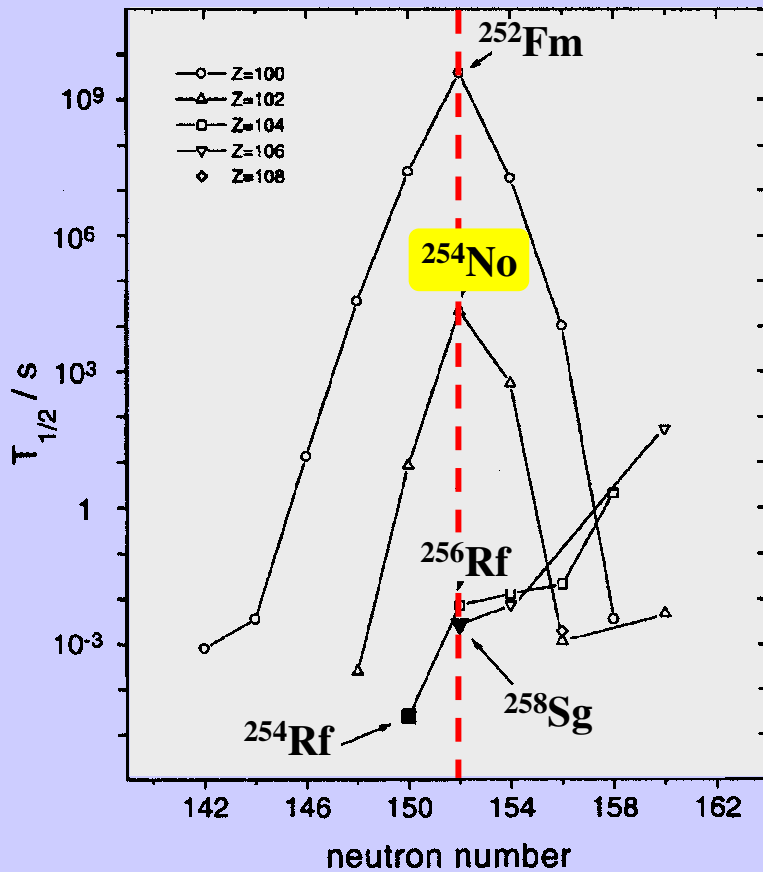


$^{208}\text{Pb} (^{48}\text{Ca}, 2n) ^{254}\text{No}$ @ 215 MeV

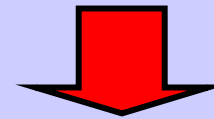
$\sigma_{2n} \sim 3 \mu\text{b}$ and $\sigma_{\text{fission}} \sim 10^4 \cdot \sigma_{2n}$!!!

$E_{\text{CN}} \sim 19.3 \text{ MeV}$

Fission halflives of transfermium nuclei



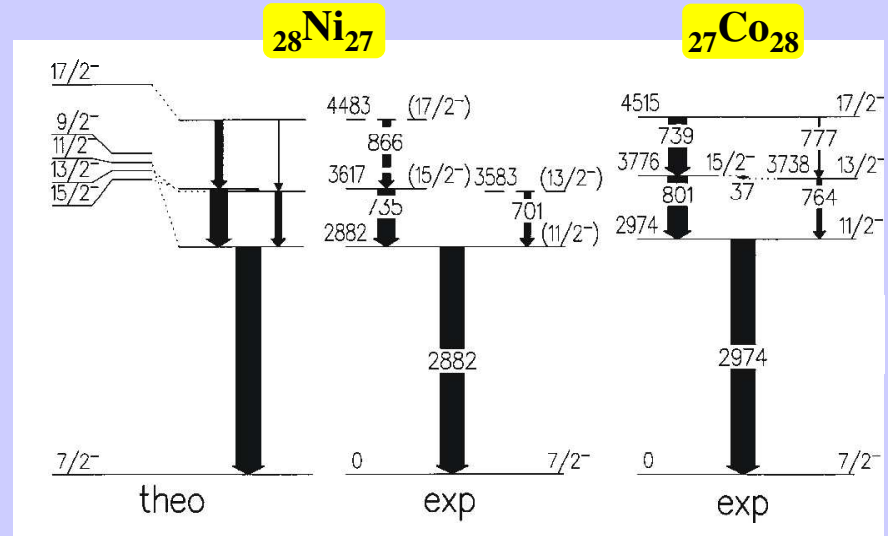
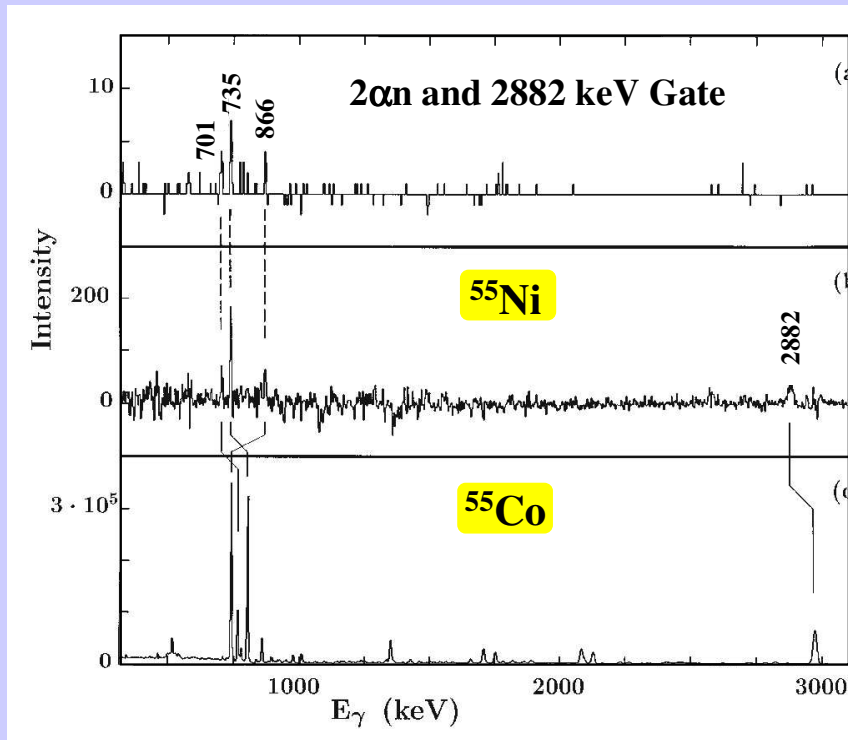
relatively long fission half-life



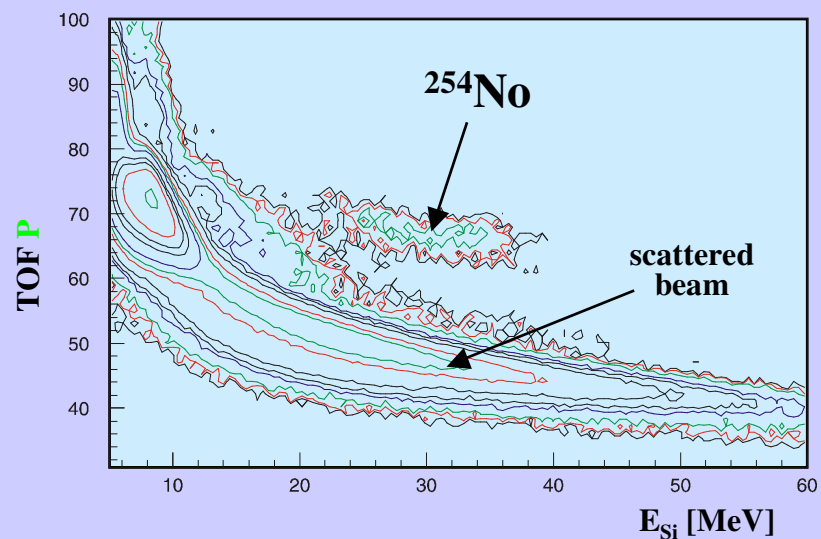
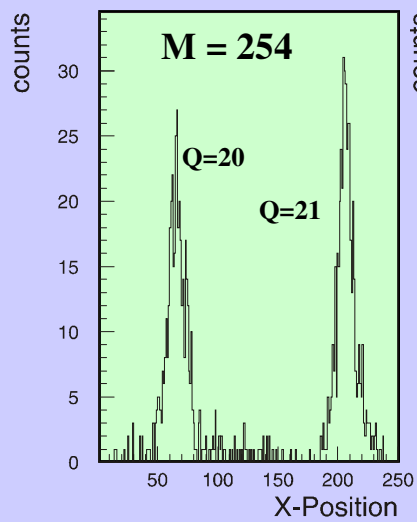
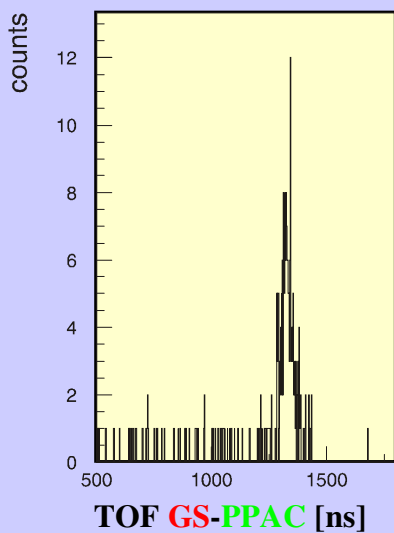
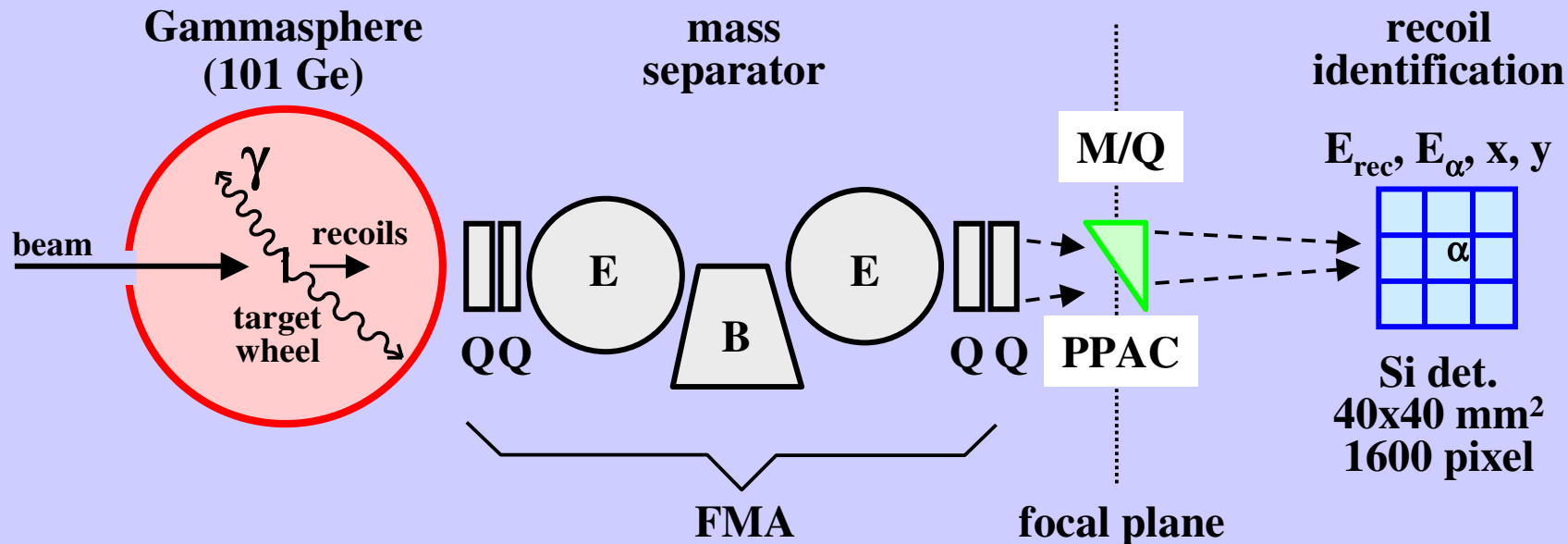
better chances to observe α -decay

Reminder ...

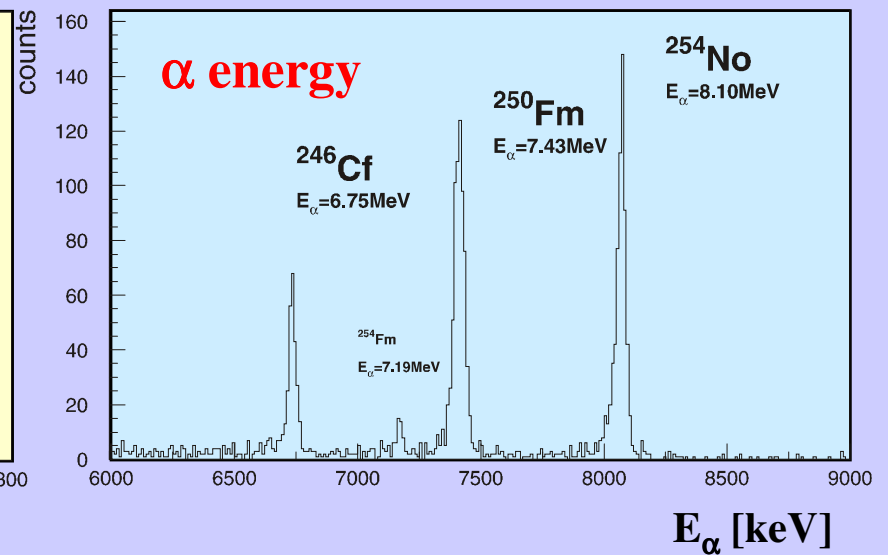
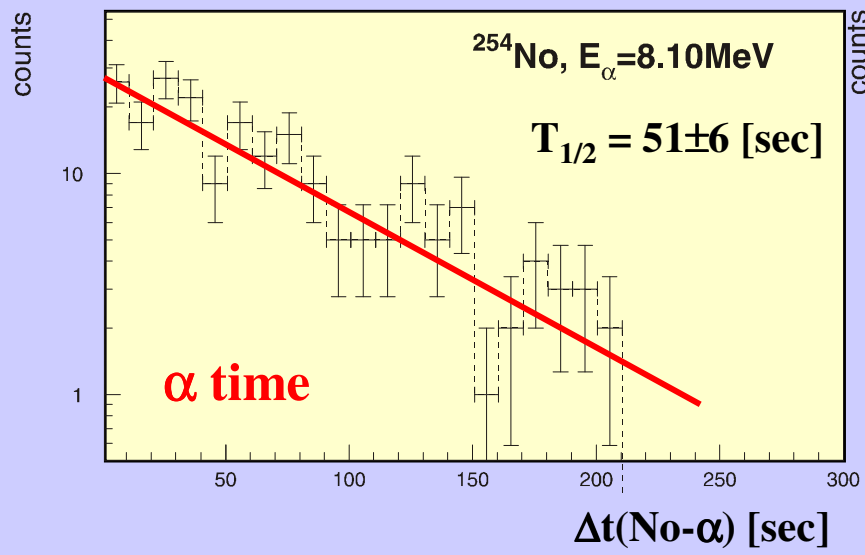
At the limit: ^{55}Ni with $\sigma \sim 0.004\% \cong 40 \mu\text{b}$



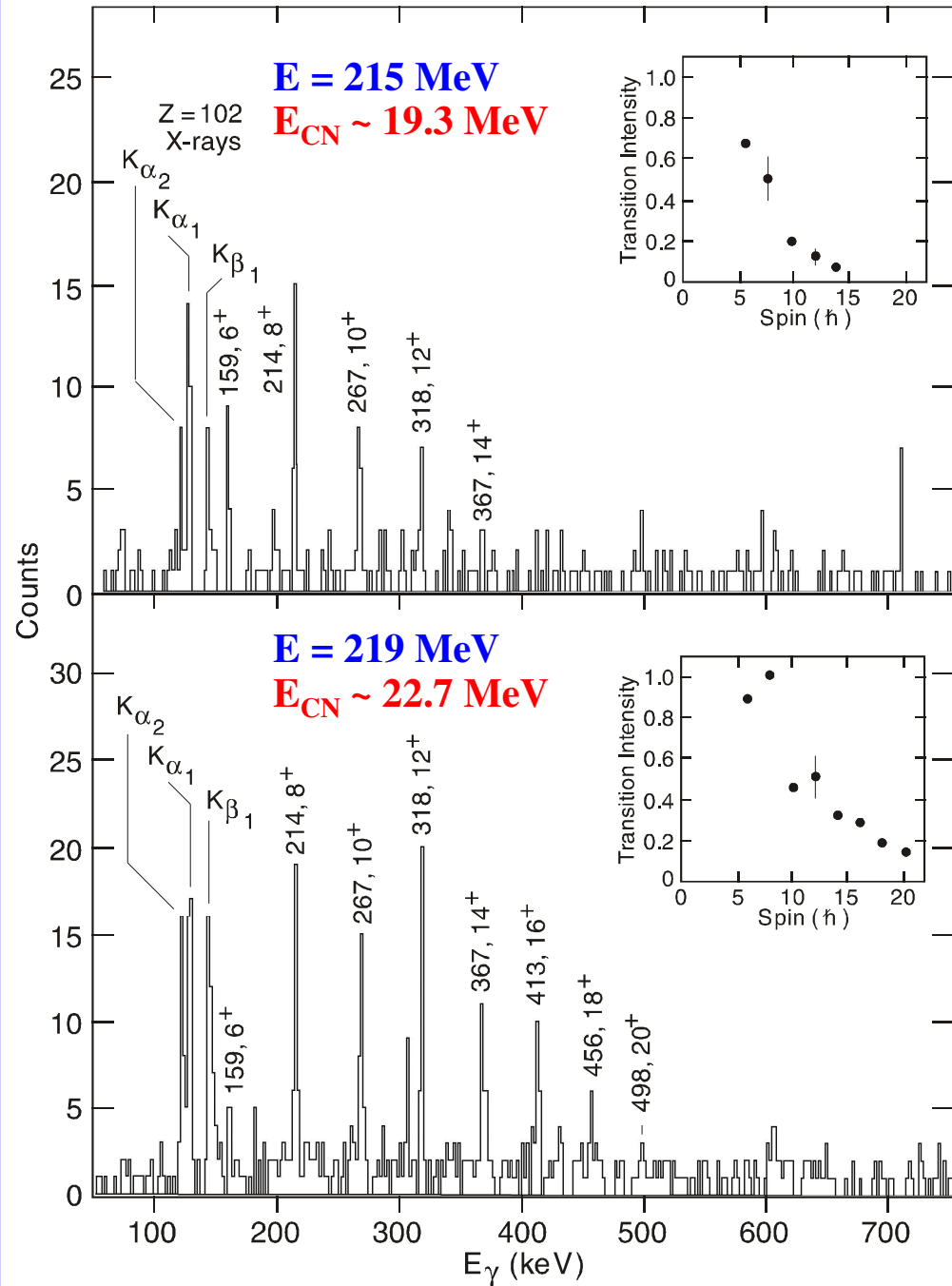
The experimental setup to study ^{254}No



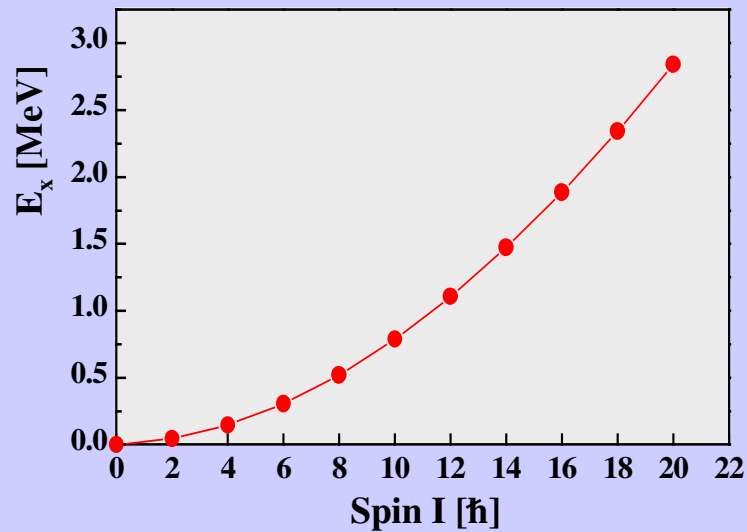
**Unequivocal identification of ^{254}No :
 α -decays in the same Si pixel**



^{254}No γ spectra from Gammasphere

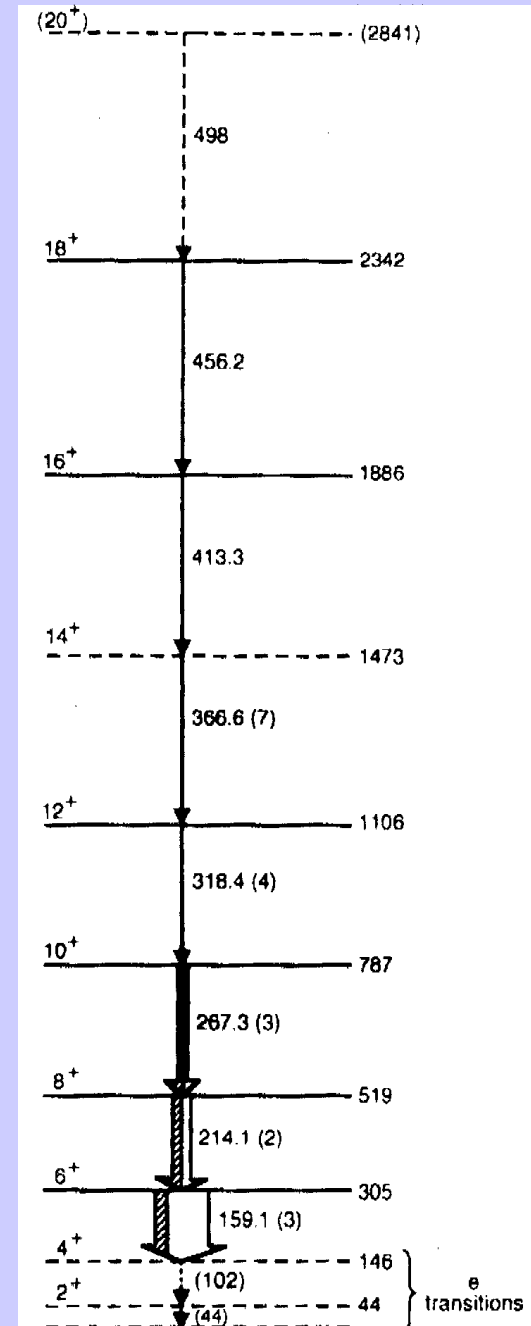


Rotational band and deformation of ^{254}No



Using an empirical formula to deduce the deformation from the energy of the first 2^+ state:

$$\beta = 0.27(2)$$



More transfermium spectroscopy ...

^{252}No (220 nb)

^{250}Fm

